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INVESTIGATION OF THE EXTERNAL FLOW ANALYSIS
FOR DENSITY MEASUREMENTS AT HIGH ALTITUDE

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NASA RESEARCH GRANT NSG 1630



Department of
Mechanical and
Aerospace Engineering

INVESTIGATION OF THE EXTERNAL FLOW ANALYSIS FOR
DENSITY MEASUREMENTS AT HIGH ALTITUDE

FINAL REPORT

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I. INTRODUCTION

Accurate experimental determination or verification of aerodynamic force coefficients (C_D and/or C_L) requires accurate simultaneous measurements of the forces (or accelerations) and the dynamic pressure ($q = 1/2 \rho U^2$). Comparison with theoretical predictions requires independent knowledge of the density (ρ) and velocity (U) to establish the proper values of the non-dimensional parameters such as Reynolds number (Re) and Mach number (M). These parameters in principle require independent measurement of temperature and measurement or inference of viscosity. Under hypersonic conditions during the early phases of re-entry the Mach number becomes a secondary parameter while the relevant Reynolds number is based on viscosity within the gas layer near the vehicle and is only weakly dependent on free stream temperature. During the earliest part of re-entry, independent knowledge of density is necessary to establish the degree of rarefaction generally measured by the Knudsen number $Kn = \lambda/L$ where λ is a relevant mean free path and L the characteristic physical dimension (either vehicle size for the overall flowfield or entrance dimensions for the local behavior at the instrument).

The Shuttle Upper Atmosphere Mass Spectrometer (SUMS) Experiment⁽¹⁾ is designed to provide independent measurement of $q = 1/2 \rho U^2$ within the high altitude range. When combined with information of vehicle velocity, it will provide independent determination of upper atmosphere density and

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coupled to accelerometer data will give the aerodynamic force coefficients within a regime difficult to simulate on the ground. The experiment is primarily intended to provide information between about 80 Km and 140 Km where rarefaction effects on the force coefficients are most important for a vehicle of the size of the Space Shuttle. It is also a regime where information on the atmosphere is relatively sparse as it lies below the altitude traversed by satellites and above that regularly assessed by ground launched meteorological vehicles. The interpretation of the measurements, however, requires an adequate understanding of the flowfield around the Space Shuttle within the vicinity of the SUMS experiment in order to provide the proper data reduction procedure and an assessment of the accuracy of the results.

At sufficiently low altitudes (below about 80 Km for the Space Shuttle), conventional pitot probe measurements can provide the dynamic pressure with straightforward data reduction and relatively minor corrections. At sufficiently high altitude (above about 150 Km), free molecular theory can be used to infer free stream conditions from surface measurements. The forces, however, are small and of little interest while the measurements require instruments of high sensitivity and are therefore difficult. In the intermediate range of altitudes where SUMS is designed to provide data, the typical molecular mean free path is of the same order as the characteristic vehicle dimensions. Figure 1 shows the

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variation with altitude of the free stream Knudsen number ($Kn_{\infty} = \lambda_{\infty}/D$) based on the free stream mean free path λ_{∞} and the Shuttle diameter D at the location of the SUMS orifice just ahead of the wheel well. Note that $Kn_{\infty} = .01$ at about 87 Km and $Kn_{\infty} = 10.0$ at about 136 Km. Intermolecular collisions can, therefore, neither be neglected (free molecular theory) nor represented by the resultant transport properties (continuum theory) over the major portion of the SUMS measurement regime. The gas properties at the entrance to the instrument are, therefore, dependent on a flowfield that can only be determined on the basis of a "molecular" theory.

In addition to the above "external flow problem", needed to establish properties near the surface of the vehicle, the entrance region of the instrument is typically either smaller or comparable to a local mean free path. In such circumstances, the connection between the gas properties at some distance into the internal plumbing and those at the vehicle surface, can be very sensitive to the velocity and angular distribution of the incoming molecules. This requires both a high degree of detail from the "external flow" and a local analysis that must assess the molecular behavior at the instrument entrance. We shall refer to this as the "entrance problem". Figure 1 also shows an approximate band of Knudsen numbers for the entrance region of the SUMS experiment. Kn_s is based on the mean free path λ_s at the vehicle surface and the orifice diameter ($d_o = .235$ cm) with surface properties fitted between free molecular results at high altitude and continuum Newtonian values at low altitude. Kn_c is a similar Knudsen number based on

conditions behind the "entrance" tube with λ_c estimated on free molecular results using Hughes and deLeeuw theory (2) at high altitude, with continuum constant pressure results applied at low altitude. Note that conditions within the entrance tube range from clearly free molecular behavior above about 110 Km to transitional behavior near 80 Km with fully continuum results only approached at the lowest altitude of interest.

The subsequent connection between the properties immediately behind the entrance tube and the measurements at the mass spectrometer shall be referred to as the "internal problem". The analytical procedures for calculating pressure profiles through the internal plumbing are well established and will be further verified by instrument calibration (1) .

During the preliminary phase of NASA Grant NSG 1630 (July 1979 to November 1979), the feasibility of examining the "external flow problem" for the Space Shuttle nose region within the relevant altitude range was established. A previously developed Direct Simulation Monte Carlo Computer Code (3,4,5) was found to be suitable as the starting point for this geometry and altitude range. Preliminary results were obtained at 87, 95, 105, 115 Km altitudes.

During the subsequent grant periods (November 1979 to September 1982) improvements in the modelling of the geometry and the molecular interactions have been incorporated in the external flow computer code. A number of runs at altitudes of 87, 95, 105, and 115 Km have been made to obtain a range of the relevant parameters and to provide input information

at the SUMS entrance location. The "entrance problem" has been examined both by using published information (2,6,7), and a previously developed Monte Carlo code for internal geometries (8). Because of the combination of entrance geometry (very long tube) and the range of local Knudsen number over the altitudes considered, a totally new "entrance" computer code had to developed. This code provides the connection between the flux information at the orifice entrance obtained from the external code and the local gas properties behind the entrance tube where the gas is in equilibrium with the "cold" walls of the internal plumbing. This new code has only been exercised to a limited extent, but preliminary results relating the pressure within the tube behind the tile to the free stream dynamic pressure have been obtained. This information coupled with an appropriate calibration of the mass spectrometer provides the basis for a viable data reduction procedure of the SUMS experiment.

Section IIA contains a brief description of the operation of the EXTERNAL computer code (the detailed code is attached in Appendix A). Section IIB describes the issues associated with geometric modelling of the shuttle nose region and the modelling of intermolecular collisions including rotational energy exchange and a preliminary analysis of the vibrational excitation and dissociation effects. Section IIC discusses the selection of the trial runs and presents the major results.

Section IIIA contains a brief description of both the

original version and the modified present code (INTERNAL) for the entrance problem (Appendix C contains the code listing). Section IIIB contains a disucssion justifying the selection of geometric, collisional and surface modelling parameters used for the trial runs. Section IIIC presents the preliminary results and discusses the major effects.

Section IV presents the conclusions that can be drawn from the present study, provides a preliminary estimate of the data reduction procedure and suggests future work.

III. EXTERNAL FLOWFIELD

The physical properties of the gas monitored by the SUMS instrument are not those of the free stream but are altered both by the intermolecular interactions in the external flowfield and by the combination of intermolecular and surface interactions within the entrance orifice and tubing leading to the instrument. External flowfield effects can be summarized in terms of the relation between local "stagnation" pressure and the free stream dynamic pressure ($q = 1/2 \rho U^2$) at sufficiently low altitudes. Within the "transition" regime of interest (80-140 Km) the gas entering the orifice is neither in equilibrium with the surface nor simply related to the free stream. The only currently available technique for describing the flowfield within this regime and providing sufficiently detailed data on the physical state of the gas at the surface is the Direct Simulation Monte Carlo Code.

A. DIRECT SIMULATION MONTE CARLO COMPUTER CODE

"Monte Carlo" is the technique of using a simulated situation and random numbers to generate solutions from which information for the real case is then deduced statistically. The Monte Carlo approach ranges from being a strictly mathematical technique for evaluating the complicated multi-dimensional Boltzmann collision integral to a complete simulation of a number of molecules, with randomness

introduced only in the initial conditions. A modification of this latter approach is the one used in the present development. It consists of simultaneously following a large number of particles which yields, to some degree, a "direct simulation" of the processes taking place. Because there are finite limits on computer storage space, a modification to the direct simulation technique was developed by G.A. Bird (Ref. 8) wherein the real gas is simulated by several thousand "sample" particles populated into cells of the sample space considered. For collision calculations, all the particles in one cell are used as a representation of the local distribution function from which collision pairs are chosen at random, but in proportion to their collision probability based on the real gas. This preceding discussion applies to a general program incorporating the direct simulation procedure. A specific computer program for the generalized three-dimensional program for axisymmetric bodies in a hypersonic multi-fluid flow is described below.

The program (EXTERNAL) conducts numerical experiments with a model multi-component gas. The real gas is simulated by several thousand molecules which may be thought of as a representative sample of the many billions of molecules in the corresponding real gas. The positions and velocity components of the simulated molecules are stored in the computer and typical collisions are computed among them as a time parameter is advanced. Since the flow is at an angle of attack to the body, three position coordinates, three velocity components and appropriate internal energy levels

must be stored for each simulated molecule.

The computation of collisions starts at zero time with the molecules moving along the flow axis at the required freestream Mach number. The body is inserted into this flow at the zero time and the desired steady flow is obtained as the large-time solution of the resulting unsteady flow.

The free-stream flow vector lies in the x-y plane. The simulated region is bounded by the x-y plane as a plane of symmetry, an outer cylindrical boundary (the axis of the cylinder is the x axis), and two planes parallel to the y-z plane. These boundaries must be set sufficiently far from the body to eliminate interference. The simulated region is divided into a number of cells which are sufficiently small for the expected change in flow properties across the cells to be small.

The first step is to generate the initial, or zero time, configuration of molecules. The molecules are distributed over the simulated region and the velocity components assigned are appropriate to a gas in Maxwellian equilibrium and moving at the required Mach number. The body is then inserted into the flow and the molecules are allowed to move and collide among themselves. The move and collide processes are uncoupled by computing a number of collisions appropriate to a time interval Δt_m equivalent to a small fraction of the mean time between collisions, and then moving the molecules through distances appropriate to Δt_m and their instantaneous velocities. The distortion produced in the molecular paths by this approximation is small as long as Δt_m is small.

compared with the mean time between collisions, and smaller than the typical transit time of a molecule through a cell.

Since the change in flow properties over the width of one cell is assumed small, the molecules in a cell at any instant may be regarded as a sample of the molecules at the location of the cell. The relative location of the various molecules within the cell can then be disregarded. A pair of molecules is chosen at random from those within the cell under consideration and is retained or rejected in such a way that the probability of retention is proportional to the relative collision probability for the appropriate interaction law. When a pair is retained, a typical collision is computed between the two molecules and the new velocity components and internal energies are stored in place of the old ones.

In general, the relative number ratio of the species of molecules in the multi-component gas will differ from unity, requiring the computation of different types of collisions. There is, therefore, one time counter for each type of collision in each cell. For each collision, the correct time counter is advanced for the cell by an amount appropriate to the collision parameters. The probability of collision, and therefore the time advancement per collision, is made proportional to the number of molecules in the cell, and the relative velocity and cross-sections of the colliding molecules. Collisions are computed in each cell until all the time counters have advanced through at least a time Δt_m . When this procedure has been carried out for all cells, the

overall experiment time is advanced by Δt_m and the molecules are moved through appropriate distances.

The set of molecules in each cell changes as the molecules are moved and appropriate conditions must be applied at the boundaries of the region begin simulated. Every boundary is treated as a source of molecules with velocity components representative of molecules moving in thermal equilibrium at the appropriate fraction of the free-stream Mach number. (The fractional Mach number is determined by the cosine of the angle between the local boundary normal and the flow direction.) Any molecule which moves outward across any boundary is regarded as being lost and is removed from the sample. The plane of symmetry (the x-y plane) is regarded as a specularly reflecting surface. Interactions with the body are also computed. The body consists of a number of conic sections rotated about the axis of symmetry. Each section must be separately specified according to the coefficients of the defining equation, a procedure to be described later in the report. For the purpose of computing the momentum and energy transfers to the surface, each region of the body can be subdivided into smaller sections. Within these smaller sections, the following three parameters must be specified: wall temperature/gas temperature, energy accommodation coefficient for each species, and tangential (momentum) accommodation coefficient for each species.

After the flow has settled down to a steady state, the number flux, momentum and energy transfers to the surface are

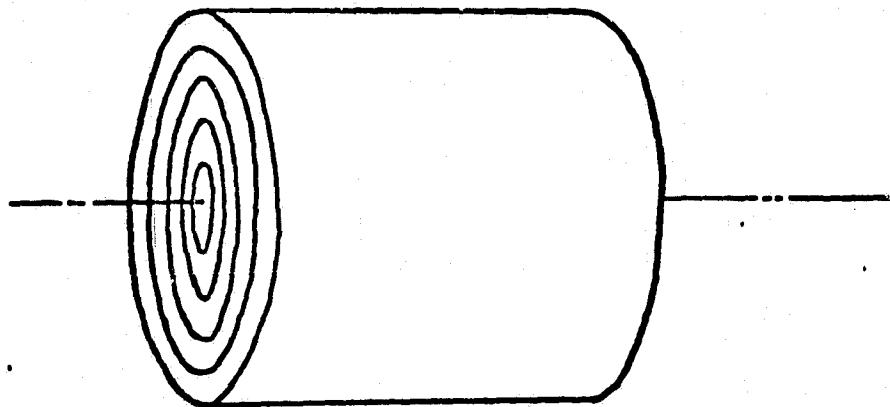
accumulated and used to compute the aerodynamic data. The time required to establish steady flow is usually assumed to be close to the time required for the free-stream flow to traverse several body lengths. The overall number flux, drag, and heat transfer coefficients are determined, along with their distribution along the surface.

In addition, it is possible to generate data on the body surface which can be used as input to the companion program, INTERNAL, described separately. (INTERNAL) computes the flow field regime inside an axially symmetric cavity, such as might be used for a spacecraft-borne sensor. The input data needed for this computation includes the molecular distribution functions present at the orifice to the cavity, the orifice being on the surface of the spacecraft.) This data consists of velocity and internal energy samples in three coordinate directions for all species of the mixture.

Flow field properties are also computed. Instantaneous values are sampled at appropriate time intervals and these are time-averaged for greater accuracy. Number densities, velocities and temperatures are printed for each cell.

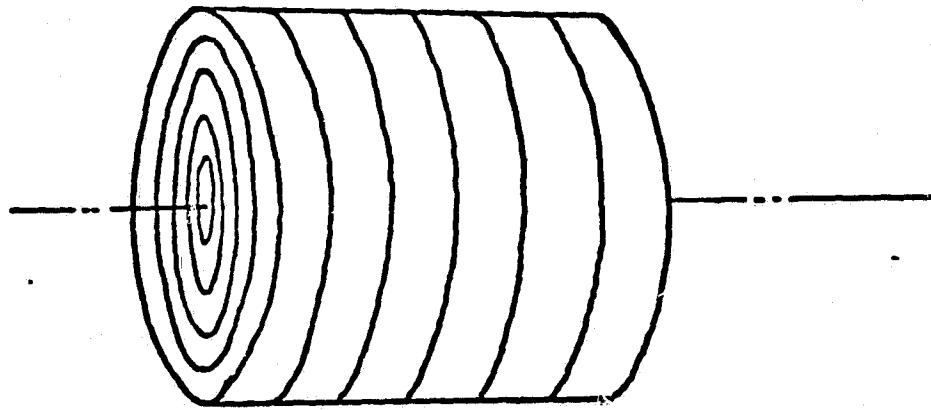
The numerical experiment takes place in a cylindrical block of space whose axis is coincident with the axis of revolution of the conic surfaces comprising the test body. This space is subdivided into cells in which the flow field properties of the experiment can be monitored. Cylindrical surfaces concentric with the axis partition the space into nested cylindrical volumes, as shown:

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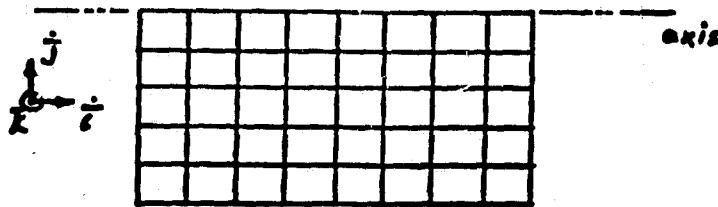


Planes parallel to the end-faces of the cylinders divide the cylinders into a stack of nested rings.

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Finally, planes perpendicular to the preceding planes and passing through the axis, called radial planes, divide up the rings in the azimuthal direction, producing cells which are wedge-shaped pieces of rings. The geometry is more easily visualized if one considers a trace of the cell configuration in a radial plane. A typical planar trace is shown:



The axis shown is the axis of symmetry of the sample space (and of the test body). This axis is considered to be the x axis. When this direction must be specified in vector algebra computations, a unit vector \hat{i} is assumed in the x direction.

Assume that the planar trace shown above is bounded by the x-axis as described and by the -y axis. This plane is at 0° azimuth angle and is called the zero plane. It is the plane normally depicted when describing the sample space.

A unit vector \hat{j} points in the +y direction. The z direction points out of the paper, and in this direction is the unit vector \hat{k} , which is given by $\hat{i} \times \hat{j}$.

The flow velocity is in the direction $\hat{i} \cos\alpha + \hat{j} \sin\alpha$, where α is the defined angle of attack between 0° and 90° . Since the flow is thus in the xy plane, there is symmetry in the z direction. That is, any condition in effect at $+z$ is also in effect at $-z$. Thus azimuth angles need be specified only from 0° to 180° where 0° is in the -y direction, 90° is in the +z direction, and 180° is in the +y direction.

Now the fact is that the gas density in the vicinity of the stagnation point of the body can become many tens of times higher than the density far from the body. It is thus desirable to use small cells in this region while the cells are larger in the regions of relatively low density. This partitioning of cell sizes is accomplished in two ways.

First, the rings can be divided azimuthally into two different sizes. This is done by specifying an angle, called THETAZ, and the number of wedge divisions both below and above this theta plane, called NWEDGE₁ and NWEDGE₂. ("Below the theta plane" means azimuth angles between 0° and THETAZ, and "above the theta plane" means azimuth angles between THETAZ and 180°.)

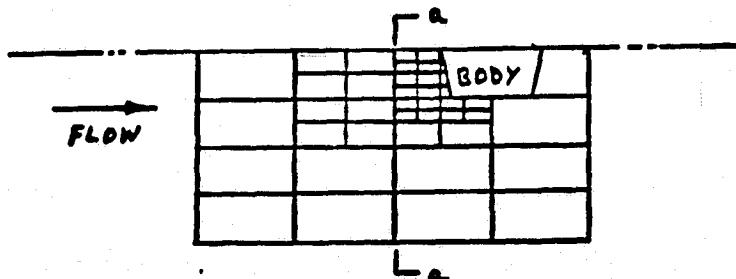
Second, the cells below the theta plane can be subdivided in the axial and radial directions down to second and then third level cells. In the 0° radial plane representation of the sample space, this would appear as large rectangles being sub-divided into small rectangles.

In this way, the sample space geometry can be tailored to the configuration of the test body angle of attack to the flow. The following examples are presented to clarify the above statement. Assume for all cases that the test body in question is a short cylinder. As explained in the section TEST BODY, the cylinder cannot have flat end faces, so the ends are cones with apex angles of about 175°.

- a) For 0° angle of attack, the flow impinges directly on the left face of the cylinder.

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It is thus desirable to have, if possible, constant azimuthal angles since there is no angles since there is no azimuthal assymetry. One possible configuration is therefore: THETAZ = 180° NWEDGE₁ = 6, NWEDGE₂ = 0 (producing 30° wedges), and the axial and radial directions can be subdivided any convenient way, producing a zero radial plane that looks like:



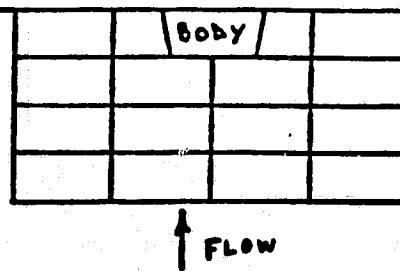
In this type of geometry, all radial planes are the same as the zero plane.

b) For 90° angle of attack, the flow impinges on

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the curved cylindrical surface of the cylinder at the bottom. It is thus necessary to have small azimuthal wedges on the lower portion of the cylinder at and near the stagnation point, while larger wedges will suffice on the upper portions of the body (in the wake of the flow).

For instance, an acceptable set of parameters is: $\text{THETAZ} = 60^\circ$, $\text{NWEDGE}_1 = 3$, $\text{NWEDGE}_2 = 3$ (producing 20° and 40° wedges), and again the axial and radial directions can be subdivided in any convenient way. Any radial plane up to 60° looks like the radial plane in example (a), with the body and smaller cells centered about a-a, while any radial plane between 60° and 180° looks like:



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- c) For an angle of attack between 0° and 90° , the configuration looks generally like that of example (b)). In this case, however, the theta plane generally should be at an azimuth angle which is relatively low (near 60°) for a high angle of attack ($45^\circ - 90^\circ$) and relatively high (near 120°) for a low angle of attack.

Because the test body is located on the axis of the cylindrical sample space, for each particle that interacts directly with the body, many more do not. In the interest of minimizing the program running time necessary to permit a statistically sufficient number of particles to strike the body, the computation makes use of zonal weighting factors. That is to say, each particle in the sample-space in reality represents one or more particles, the actual number depending upon the weighting factor of the zone in which the particle currently exists.

Up to five cylindrical boundaries are selected across which the zonal weighting factors change. These boundaries

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are specified in terms of the number of first level cells between the axis and the boundary. The change in the zonal weighting factor across each boundary can be given by:

$$LF_n = \frac{LD_{n+1} + LD_n}{n-1} \quad n=1,2,\dots,5$$
$$\prod_{j=0}^{n-1} LF_j$$

where the LD values are the number of first level cells between the axis and the cylindrical boundaries; and $LF_0=1$, $LD_6=LD_5=LD_K=NH$ were $K \geq K_{last}$.

This equation is the result of having the zonal weighting factors defined in such a way that they are equal to the ratio of sample space volume in the zone above the weighting-factor boundary to the sample space volume in the zone below the weighting factor boundaries. The importance of the region near the axis can be emphasized by choosing LF values larger than those given above.

The sample space is populated with a distribution of gas molecules. Each molecule is assigned a velocity, a rotational

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energy and a position, such that the sample space is uniformly (albeit in a random manner) filled with a gas in thermal equilibrium and flowing at the required Mach number. Each molecule is assigned a number corresponding to its zonal weighting factor. The molecule thus represents in actuality a number of molecules (including itself) equal to the zonal weighting factor. While the molecule moves in such a way as to stay within the given zone, its weighting factor does not change. If it crosses a weighting factor boundary while moving in toward the axis, a number of molecules is added to the distribution. The number of molecules added is equal to:

$$\frac{\text{old weighting factor}}{\text{new weighting factor}} - 1$$

For instance, assume that above a boundary, the zonal weighting factor is 6, and below the boundary it is 3. Hence $6/3 - 1$, or 1, molecule must be added to the distribution when the molecule crosses the boundary. This is clearer if one considers that above the boundary, the molecule represents a total of 6 molecules. When the molecule drops below the boundary, it can only represent 3 molecules, so another molecule must be added to the distribution to represent the other 3 molecules.

The added molecule(s) is (are) assigned the same position and velocity components as the original molecule. While this does not approximate true kinetic theory at first, in practice the positions and velocities are soon randomized

by collision processes.

On the other hand, if a molecule crosses a weighting factor boundary while moving away from the axis of the sample space, there is a probability that it must be dropped from the molecular distribution. The probability is given by

$$1 - \frac{\text{old weighting factor}}{\text{new weighting factor}}$$

The random number generator is used to generate a histogram of disappearing molecules to match the actual probability of disappearance. The whole idea behind using weighting factors is to increase test body - flow field interaction in a given running time. Thus, when body surface quantities are accumulated (like flux, energy, momentum, heat), they are accumulated in terms of the weighted number of molecules striking the body. This is particularly important if the body exists in two or more weighting zones, so that surface quantities in different zones can be correctly compared.

The program is set up to handle a test body consisting of a sequence of connected conic sections rotated about a common axis of symmetry. Some typical surfaces which can be considered without modification are sections of spheres, cones, cylinders, ellipsoids, hyperboloids, and paraboloids. A disc perpendicular to the symmetry axis cannot be handled without modification, since it represents a multi-valued function in r . Cones with very large apex angles ($\sim 180^\circ$) are

used in place of discs. The procedure used to specify body surfaces is to generate a form of the standard equation of the surface in question, substitute into this equation the actual values of the constants, and reduce the standard form to the following type of equations:

$$Ak^2 + Br^2 + Ck + D = 0$$

where A, B, C, and D are numerical values

These coefficients are used as input data to the program. The program requires additional parameters for each conic surface. Each surface can be axially divided into several segments for the purpose of accumulating the body surface parameters like flux, heat transfer, etc. The x-coordinate on the right side of the segment is required as data. (The segments are divided azimuthally by the same radial planes that partition the sample-space.) In addition, the temperature, energy accommodation coefficients and tangential accommodation coefficients for each species are required for each segment. The temperature of a surface is normalized by the free-stream temperature.

The tangential accommodation coefficient specifies the fractional part of incoming tangential momentum that is lost to the surface. For the reflection model used in this program, this also represents the fractional part of the molecules colliding with the body surface whose collisions are diffuse. The remainder collide specularly.

The energy accommodation coefficient specifies the ratio

of the net molecular energy flux absorbed by the body to that energy which would be absorbed if all re-emission were appropriate to equilibrium at the surface temperature of the body. This coefficient together with the tangential accommodation coefficient and the surface temperature determines the effective temperature of the diffuse component of the re-emitted molecules in the reflection model used in this program.

It is also possible to collect velocity samples of the colliding molecules on a restricted number of surface segments to generate distributions for the internal flow program. For this purpose, molecules that collide with the body surface are considered in two classes. The class called "UNCOLLIDED" consists of "free-stream" molecules. That is, these molecules have not previously collided with the body surface or with any other molecules other than "free-stream" molecules. The other class, called "COLLIDED", consists of molecules which have previously collided with either the body surface or with other "COLLIDED" molecules.

Appendix A contains the full listing of the computer code in Fortran IV. The main program consists of dimensioning statements coupled to a fairly detailed description of the input cards (using the NAMELIST input). The main operating program is called RUN which in turn calls the appropriate subroutines. Figure 2 shows a schematic of the flow chart for the operation of the Program EXTERNAL.

B. Modelling of the SUMS Problem

1. Geometric Modelling

The computer code described above requires an axisymmetric geometry of the body, although the flow vector can be at an arbitrary direction. The Space Shuttle nose geometry in the vicinity of the SUMS orifice has to be modelled by an equivalent body of revolution. A paraboloid of revolution around an axis inclined at about 8° with respect to the shuttle axis models the lower surface cross sections of the body both in the symmetry plane and in the transverse direction reasonably well. Figure 3 shows a sketch of the actual and modelled Space Shuttle nose geometry. This model will be called the parabola model, and the flow direction of 32° with respect to the axis will be used to represent the 40° angle of attack of the Space Shuttle.

Both for the purpose of benchmarking the test runs against continuum results and to achieve better resolution an alternative axisymmetric flow model was examined. This model consists of a hyperbola rotated about the flow velocity vector passing through the stagnation point. Figure 4 shows a typical cell geometry for this model. This is also the model used by Professor Clark Lewis for his continuum calculations. As will be seen from the results in the next section the latter (axisymmetric) model is questionable as a representation of the flow in the stagnation region of the

Space Shuttle, at least at higher altitudes.

2. Modelling of Cross Sections

In order to avoid uncertainties associated with the choice of hard sphere cross-section to best model the "typical" collision, an energy dependent inverse power law cross-section collision code has been incorporated in the program. The power law exponent and reference cross-section is chosen to provide the best fit to the viscosity temperature dependence over the range between the wall and stagnation conditions which represents the energy range of the important collisions. Figure 5 shows that viscosity can be matched to within a few percent over the relevant range with a single choice of exponent and reference conditions

Using the exponent of $N=.552$ and a reference temperature of 1000°K four axisymmetric cases of a hyperbola at 0° have been run to simulate 87KM, 95KM, 105KM and 115KM altitudes. Since in these cases rotational energy was not included these represent a fictitious monatomic gas of $\gamma=5/3$. Figure 6 shows "smoothed" temperature profiles normalized to the free stream temperature. Note that at the two lower altitudes a relatively "flat" Rankine Hugoniot (R.H.) region exists, while at 105KM R.H. conditions are barely reached and at 115KM the temperature peaks at about 70% of R.H. temperature. The shock layer in all cases is, however, very much thicker than the inviscid result based on nose radius R_N . Note, however, that the hyperbola is very blunt and it is not clear that the nose radius is the appropriate dimension or that this models properly the flow about the Space Shuttle Nose.

3. Rotational Energy Exchange

On the basis of previous experience in modelling rotational energy exchange in Reference 4 the model of Larsen and Borgnakke (9) was chosen as appropriate for the blunt geometry under investigation. The External Flow Program was therefore re-coded to include an arbitrary number of internal energy modes but with a single relaxation time related to the parameter ϕ which ranges from $\phi=0$ representing no exchange to $=1$ simulating maximum available exchange at each collision. An indication of the effect at 115KM is shown in Figure 7 and 8 for the most rapid energy exchange ($\phi=1$). Note that, as expected, the peak temperature is lower for the $\gamma=1.4$ case because some of the energy goes into rotation. Also note that significant non-equilibrium exists between translation and rotation even at this maximum energy exchange rate corresponding to $\phi=1$. The effect on surface properties shown in Figure 8 is virtually non-existent for the shear, but noticeable on the heat flux and pressure.

4. Preliminary Attempts at Comparison with Continuum Results

Some preliminary comparisons of both the monatomic ($\gamma = 1.667$) and diatomic ($\gamma = 1.4$) runs for the axisymmetric hyperbola were made with continuum results provided by Prof. C.H. Lewis based on continuum theory (10,11). Among the surface properties only the pressure distribution agreed reasonably well. Since this is the property least sensitive to detailed flow field behavior the agreement is not a very sensitive test. A typical comparison is shown on Figure 8. A comparison of heat flux is not shown on this figure as it is not the same scale at this high altitude. The continuum result of Prof. Lewis gives the stagnation point heat transfer coefficient $C_H = 0.1/2\rho U^3 = 2.8$ which is physically unrealistic. At 105KM the continuum heat transfer coefficient C_H is below one but still appreciably above the Monte Carlo results. At 95KM the Monte Carlo results give heat transfer that is almost twice as large as Prof. Lewis's result for $\gamma = 1.4$.

The most significant discrepancy between the continuum and Monte Carlo results arises in the shock layer thickness. While the definition is somewhat arbitrary, a comparison of the subsonic region in the vicinity of the stagnation point does give a good indication of the extent of influence of the "downstream" portions of the body. Figure 9 shows the Monte Carlo results for the $M=1$ line at 115KM together with the shock line from C.H. Lewis. This discrepancy led to a whole re-examination of the modelling of the nose geometry.

5. Re-examination of Geometric Modelling

Since the subsonic region appears to extend to the "shoulder" wherever the hyperbola is terminated as shown in Figure 9, the applicability of the hyperbola model becomes suspect, at least for the Monte Carlo calculation where upstream influence cannot be eliminated. In order to further assess the effect of the geometric model on the results a comparison of the hyperbola at 0° to the parbola at 32° are presented in Figures 10 and 11. It can be seen from Figure 10 that the subsonic layer in the vicinity of the SUMS orifice is significantly different in the two cases. Figure 11 shows that while the pressure is not dramatically affected by the model the heat flux and shear are significantly altered. On the basis of these results it was concluded that while the runs for the hyperbola may indicate trends they are neither representative of the SUMS region on the Space Shuttle nor good candidates for benchmarking with the

continuum results. All subsequent runs were therefore made using the geometric model of a parabola at 32° angle of attack.

6. Inelastic vibrational Excitation and Dissociations

Because of the high energies of collision between free-stream and reflected molecules, inelastic collisions (certainly vibration and dissociation) are in principle possible. At the highest altitude the number of such collisions is expected to be insignificant because the cross-sections are low enough so that a typical molecule will reach the body with a negligible probability of a previous vibrationally exciting or dissociating collision. At the lower altitudes 87 to 95Km the number of collisions suffered by a typical molecule before reaching the body surface is measured in the tens or hundreds and therefore inelastic cross-sections that are only a few hundredths of the elastic and rotational ones may produce significant effects. The detail needed to properly model these collisions in the Monte Carlo Programs far exceeds the available experimental information, which primarily gives overall rate constants. An investigation of the best combination of analytical and empirical information was initiated early within the grant period. A theoretical attempt to couple low energy vibrational excitation experimental information to the highly non-equilibrium high energy collisions through theoretical work was initiated. Appendix B contains a Master's thesis

presenting the formalism and giving the initial results of a theoretical formalism. Based on those results, coupled to some limited experimental data, a method for determining the probabilities of specific outcomes in individual collisions in the Monte Carlo Programs can be developed. The increase in computing times, however, may make the feasibility of using such a code, for anything but benchmarking, prohibitive.

C. MAJOR RESULTS FOR EXTERNAL FLOW

The primary effort during the grant period up to February 1981 was spent in developing the code for external flow and establishing the appropriate geometry to model the region of the Space Shuttle in the vicinity of the SUMS experiment. The major results to that date are presented in the renewal proposal for the period February 1, 1981 to July 30, 1982 which also served as a progress report on the previous grant period. (12) We will only summarize those results here and update them on the basis of additional external flow computations.

A representation of the shuttle nose geometry as a paraboloid of revolution around an axis 8° from the actual shuttle axis (Figure 3) was found to adequately model the windward side of the shuttle in the vicinity of the SUMS orifice. Computations for both this model at an effective angle of attack of 32° and an alternative axisymmetric flow about hyperboloid model centered on an axis through the nominal stagnation point, demonstrated that the paraboloid model at angle of attack is necessary to adequately model conditions near the SUMS orifice. The typical body and computation cell geometry is shown in Figure 12.

A reasonable indication of the flowfield can be obtained by examining density contours or Mach number contours about the body. Figure 13 shows the sonic lines and the $M = 5$ lines indicating approximately the outer extent of the "shock" layer at 95 Km and 115 Km, within the plane formed by

the velocity vector and the axis of the modelled paraboloid. Note the greater "shock" layer thickness at the higher altitude. Also note the fact that even at 95 Km, the shock layer is a large fraction of the local body dimension such as nose radius. The "shock" thickness is a major portion of the entire "shock layer" casting doubt on any calculation incorporating a thin "shock" assumption. Figure 14 shows some density contours on the "windward" side of the body at 95 Km altitude as well as the sonic line. Note the rather constant density rise towards the body with no discernible separation between "shock" and "boundary layer." Also shown are estimates of the stagnation streamline and another streamline along which the velocity only goes slightly subsonic. Note the rather gradual turning along the latter streamline and the rather diffuse nature of the shock layer even at this altitude where the nominal free stream Knudsen number is 0.04. In order to give a better picture of the three-dimensional aspect of the flowfield, a sketch of one-half of the paraboloid and some contour plots of $M = 1.$ and $M = 5.$ lines are shown on Figure 15.

Information in Figures 13 through 15 gives some indication of the nature of the flowfield in the vicinity of the SUMS orifice. The ultimate objective, however, is to establish properties at the vehicle surface. Figure 16 shows the variation of two surface fluxes (normal pressure and heat flux) at 95 Km along the four cross-planes shown in Figure 15. The normal pressure p_s normalized by $1/2\rho U^2$ ($C_p = p_s / 1/2\rho U^2$) is shown on the right side of the figure while a heat

transfer coefficient ($C_H = Q/1/2\rho U^3$) is shown on the left. Figure 17 shows the variation versus angle in the cross-plane nearest the one containing the SUMS orifice, and compares the results to some theoretical and semi-empirical predictions. The pressure coefficient, as expected, lies approximately between the free molecular (C_p (F.M.)) value and the modified Newtonian (C_p (NEWT) value, up to about 90° . (Note that at 90° around the axis the local angle (β) between the velocity vector and the surface normal is approximately 64° at this cross-plane.) The heat transfer coefficient lies substantially below the free molecular value. It is also compared to a semi-empirical extrapolation of experimental results of stagnation heat transfer on hemispheres presented in reference 13. Direct comparison is clearly dubious due to the substantially different body geometry and the implicit assumption of totally local behavior contained in the $\cos\beta$ variation with local angle of the surface to the velocity vector. Some additional degree of uncertainty is contained in the choice of "body size" that is used to evaluate the correlation parameter K_R^2 .

The ultimate objective of the external flow calculations is to provide information on the properties of the gas entering the SUMS orifice. It is the potential non-equilibrium nature of the entering distribution of molecules that is responsible for the difference between the surface "pressure" and the gas pressure within the internal plumbing around the mass spectrometer. The non-equilibrium aspect is most commonly represented by a temperature jump and a

velocity slip within the continuum formalism. Under the highly rarefied conditions at the upper end of the altitude range of interest, even that description may be inadequate because of the highly non-Boltzmann distribution of the molecules arriving at the surface. Figure 18 shows the distribution function of the flux of molecules arriving at the SUMS orifice at 95 Km and 115 Km obtained from the external program. The flux distribution is plotted both versus molecular velocities normal to the local surface and tangential within the plane formed by the free stream velocity vector and the local normal. Note that at 95 Km one could fit the distribution by a Maxwellian with some temperature different than the body temperature and a slip velocity which is comparable to the free stream mean molecular speed. At 115 Km the distribution is clearly composed of two components, the free stream molecules arriving unperturbed at the surface, and the collided molecules having a broader distribution possibly representable by another Maxwellian. Clearly this potential bimodal character of the incoming molecules must be recognized in the evaluation of the "entrance" problem at the SUMS entrance.

As a companion to the initial version of the external code (EXT) an internal code (INT) was developed (8). The objective of this code was to determine local properties and surface fluxes in a cavity connected by a tube to the exterior surface. The input molecular distribution is obtained from the surface flux information provided by the external code. While the code is in principle general to allow a wide variety of geometric configurations it is optimized for a short tube-large cavity geometry. Early attempts to apply this code directly to the SUMS inlet geometry resulted in very long running times with little assurance that a steady state solution had been achieved. A total recoding of the entrance problem was therefore implemented resulting in a code (INTERNAL) that allows greater flexibility in handling the long tube geometry of SUMS as well as incorporates all of the changes in intermolecular collisions, numbers of species and rotational energy exchange developed for the code EXTERNAL.

A. Entrance Computer Code INTERNAL

The purpose of the program described in this section is to determine the fluid field inside a cavity which consists of a connected sequence of conic sections rotated about an axis of symmetry. The cavity is considered in two parts: a main cavity on whose interior surface the sensor will be

located, and an inlet tube whose orifice is presumed to be at the exterior surface of a spacecraft. A detailed code listing is given in Appendix C.

The input data for this program includes a molecular distribution function which is obtained from an external-flow run. Some of the other data refers to input parameters of the external-flow run. Thus it is seen that the pair of programs can be run as a set, computing the conditions inside a cavity which exist for a given set of conditions in the undisturbed free-stream flow. The programs have not yet been directly coupled, although they are written with this intent.

The program was constructed by turning an external-flow program inside out. In doing this, the basic molecule/body collision mechanism is preserved with only minor changes, while the molecule/molecule collision mechanism are not changed in any way.

A general description of the way in which the internal-flow program conducts the experiment is unnecessary since the description in the preceding section generally applies here. Any important differences will be described as they occur.

The numerical experiment takes place in a cylindrical block of space quite similar to that used in external flow. However, the inlet region must be cylindrical in shape while the cavity region can be defined by conic sections.

There is no subdivision of 1st level cells into smaller sized cells. Also, the number of azimuthal wedges is specified for a full 360° (since there is no plane of symmetry determined by an input flow direction) and is the

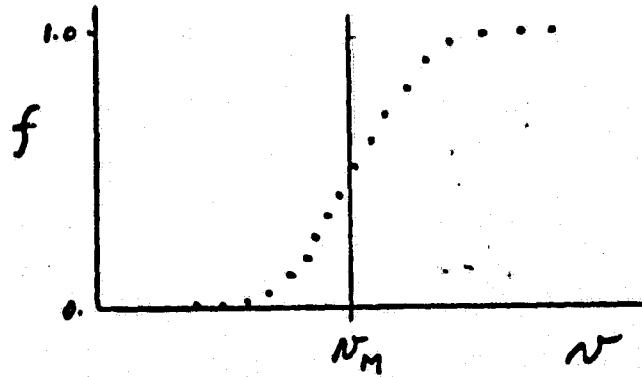
same for both regions. No weighting factors are used as the desired surface fluxes on the walls generally occur near the outer edges of the sample space.

The orientation of the sample space is determined by the geometry of the particular body surface segment in external flow which is used as the orifice area of internal flow. The body surface normal of external flow becomes the x-axis (axis of symmetry) of internal flow; the direction given by the cross product of the x-axis of external flow and the body surface normal becomes the z-axis of internal flow; and the direction given by the cross product of the above cross product and the body surface normal becomes the y-axis.

The zero radial plane is the positive x-y plane, and azimuth angles range from 0° to 360° , with the positive z-axis at 90° .

The sample space is initially populated with gas molecules in a manner similar to that used in external flow. In this program, however, the gas molecules are initially in thermal with the walls locally. The selection of the density profile initially in the tube indeed poses a problem and is very critical in minimizing running time. The code therefore allows an arbitrary initial distribution specified by the user. The choices and procedures for selecting them are described in the next section. The key objective is to provide a distribution that is consistent with the input flux and will not have to be severely altered in magnitude to arrive at the steady state.

The molecular input distributions are generated from data produced by an external-flow run. This data consists of velocity samples of molecules impinging on the external body surface segment which is considered to be the orifice area of the inlet region. The form of the velocity samples is a series of horizontal S-shaped curves, one for each molecule type, in each of three directions, for both UNCOLLIDED (free-stream) and COLLIDED molecules. Each s-curve gives the fraction of molecules (impinging on the body surface segment) with a velocity \leq the given velocity. A typical curve for UNCOLLIDED molecules is given [v_M is the velocity corresponding to the center velocity of the distribution as computed in the external flow program]:



The information from these curves is used directly to generate the molecular distribution of the molecules entering

the inlet.

The population (pressure) inside the cavity region is selected initially as an input. The computer simulation proceeds until a steady profile inside the tube entrance region is generated. Since this does not necessarily require zero net flux, a series of runs with different cavity pressures is generated and a cross plot of flux versus pressure is the output for a single external flow input condition. The equilibrium solution if desired can be obtained from the zero flux point.

B. Scope of "Entrance Problem" for SUMS experiment

Within the free molecular regime the "entrance problem" has been examined by Hughes and deLeeuw⁽²⁾. In order to indicate the potential magnitude of the problem, Figure 20 shows the variation of both the surface pressure (p_s) and the chamber pressure (p_c) versus angle of attack under free molecular conditions at infinite speed ratio and a very long tube (conditions approached at high altitude during Space Shuttle re-entry). For the SUMS location, the local angle of attack β is approximately 28° giving a possible difference of a factor of 10 between the surface pressure (p_s) and the internal chamber (p_c) that directly affects the mass spectrometer reading. While this theoretical result is expected to be accurate at 140 Km and above, at lower altitudes the local flux distribution will have both a directed and a rather diffuse component (see Figure 18), and

also the effect of internal collisions within the tube will begin to play a role (see Figure 1).

Because the tube length to diameter ratio is very large (37) the time constant and therefore the computing time to reach equilibrium is very long. For this reason no attempt is made to simulate the problem all the way to the condition where the chamber pressure behind the tube is at the correct value to nearly balance the net flux. Instead a series of runs with different assumed "cavity" pressures are performed and the zero net flux (or a given small value for the dynamic condition) can be selected to interpolate the correct "cavity" pressure. This procedure if it proves generally successful can of course be automated within the code.

C. Preliminary Results of INTERNAL Code

Preliminary results for the SUMS entrance tube geometry are presented as a couple points on Figure 21 giving to the ratio p_c/p_s versus altitude.

Also shown are some theoretical curves for full continuum (Continuum), for continuum external flow but free molecular flow through the tube (No slip, $Kn_s \rightarrow \infty$), and for fully free molecular flow (Free Molecule). In addition, some points using a best fit Maxwellian to the external flow flux and free molecular internal flow from reference 2, are shown at 95 Km and 115 Km. These results are preliminary, but they do indicate the magnitude of the correction and the fact that no simple use of the currently available results can cover

the entire range of altitudes of interest.

IV. DISCUSSION

If we couple information from INTERNAL (Figure 21) to the results of the external flow computations (summarized as a plot of $p_s / 1/2 \rho U^2$ versus altitude in Figure 22), we can produce a preliminary estimate of the data reduction curve that could be coupled to the calibration of the mass spectrometer to deduce the $q = 1/2 \rho U^2$ during Space Shuttle re-entry. Figure 23 is a cross-plot of the q/p_c versus p_c obtained from Figures 21 and 22 with the 87 and 105 Km results only estimated on the basis of interpolation of Figure 21. The establishment of such a data reduction curve for the nominal re-entry conditions, together with associated error bars as well sensitivities to wall temperature and angle of attack is necessary for the proper interpretation of data to be obtained by the SUMS instrument.

The data measured by the mass spectrometer in the SUMS experiment essentially provides collector currents of charged species of different masses. Calibration of the instrument can relate these to the overall pressure and composition at the entrance to the instrument being calibrated. Since the actual environment of the flux of molecules to be encountered under flight conditions cannot be simulated the calibration is performed with the incident flux essentially in equilibrium with the instrument outer walls (room temperature and no flow). The data reduction procedure must therefore relate the effective environment in the ground test simulation to the desired dynamic pressure q under the

flowing non-equilibrium condition and through the calibration to the instrument measurements.

As discussed above, capability now exists for calculating the surface flux distribution of molecules entering an opening at the SUMS location. The computational procedures for connecting that information to the pressure immediately behind the entrance tube has also been developed, although not fully exercised over the entire range of parameters applicable to the SUMS experiment. That pressure, in turn, can be directly related to the ground test environment used to produce the calibration curves for the mass spectrometer.

The primary objective of future work must be to provide a data reduction procedure that relates the spectrometer reading to the free stream dynamic pressure ($q = 1/2 \rho U^2$). With currently available procedures a relation between the pressure p_c and q such as the preliminary one shown on Figure 23 can be obtained using the best available information on the flight parameters such as velocity, angle of attack, tile temperature, surface conditions, molecular collision parameters, etc. This relation must then be combined with the calibration curve where the instrument readings are related either to p_c directly, or to a calibration pressure which can be related to p_c by conventional means. A single plot of q versus total measured collector current can thus be obtained from the combination of these results.

The relation between the dynamic pressure q and the calibration pressure p_c depends on many parameters of the

problem. Some of these, such as flight velocity, angle of attack, and tile surface temperature are expected to vary only slightly from their nominal values. Since the actual measured values of these quantities will be available on each individual flight, corrections to the data reduction relation should be evaluated in the form of sensitivity coefficients for small changes from the nominal. Studies to determine the effects of these parameters are necessary to establish the significant sensitivity coefficients that must be incorporated into the data reduction scheme.

The parameters such as surface accommodation, surface recombination, and free stream composition can also affect the results. In addition, modeling simplifications of both the geometry of the problem and the molecular collision phenomena can alter the quantitative value of the relation between q and p_c . Because of the unavailability of any in-flight measurements that could lead to an evaluation of these parameters, bounds on the uncertainties they produce should be studied. Sensitivity of the data reduction relation to the most significant of these can then produce bounds on uncertainties on the dynamic pressure q , due to reasonable variations. The final goal of future work is an algorithm for the evaluation of the dynamic pressure q , from the calibration pressure p_c , together with error bounds, due to the uncertainties associated with the external flow and the entrance problems.

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FIGURE 1. RELEVANT KNUDSEN NUMBERS VERSUS ALTITUDE

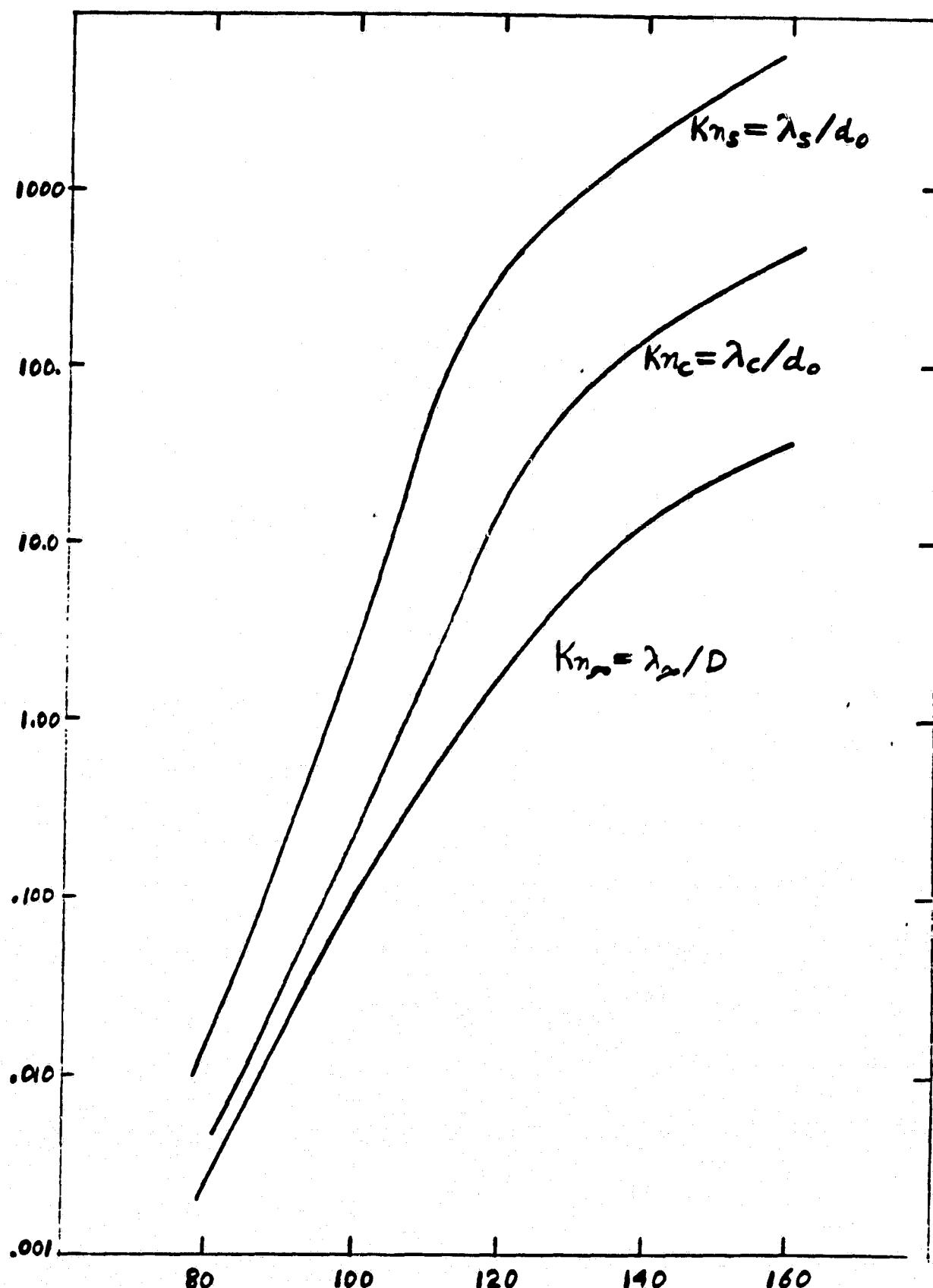


FIGURE 2. SYSTEM FLOW-CHART

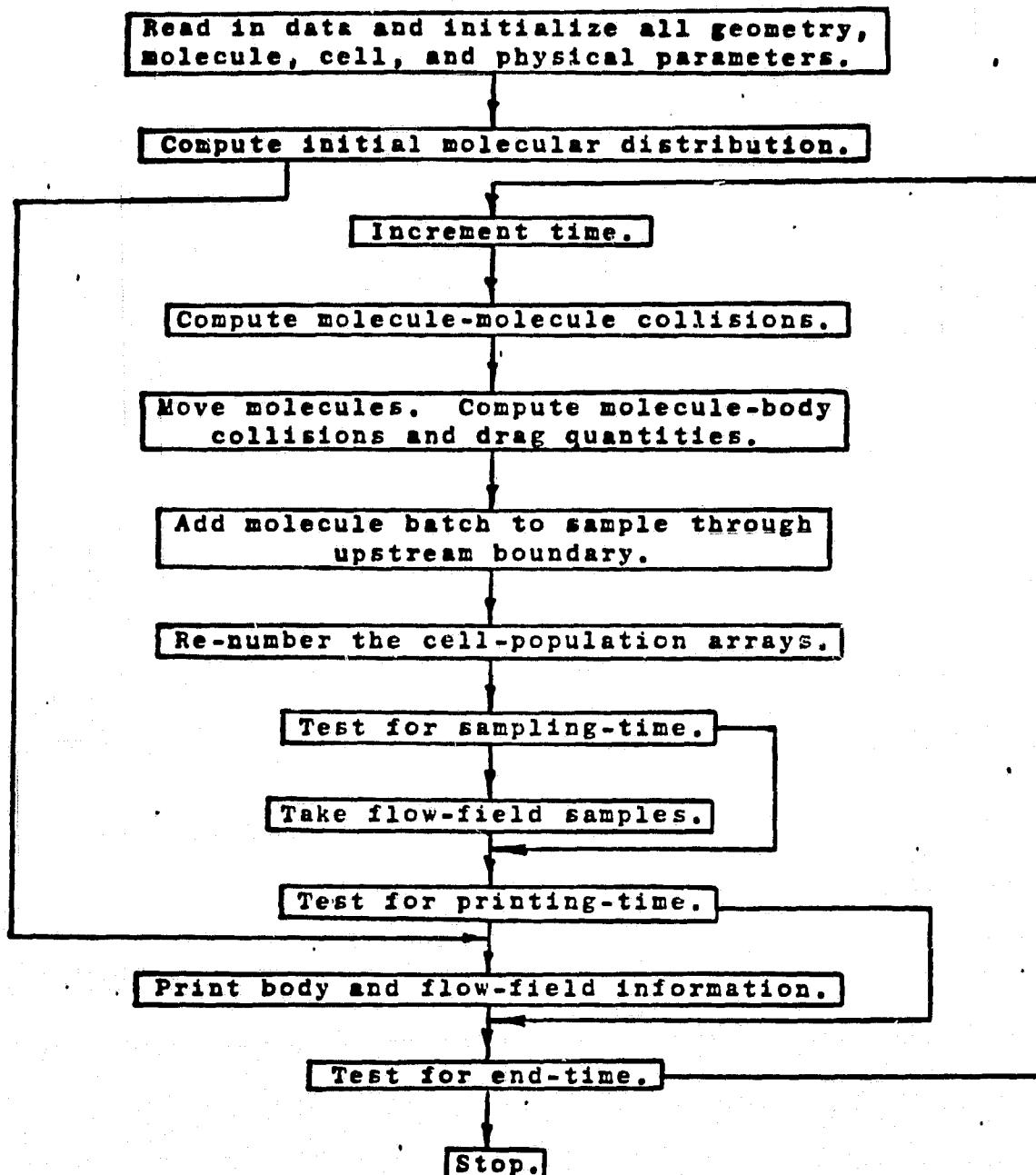
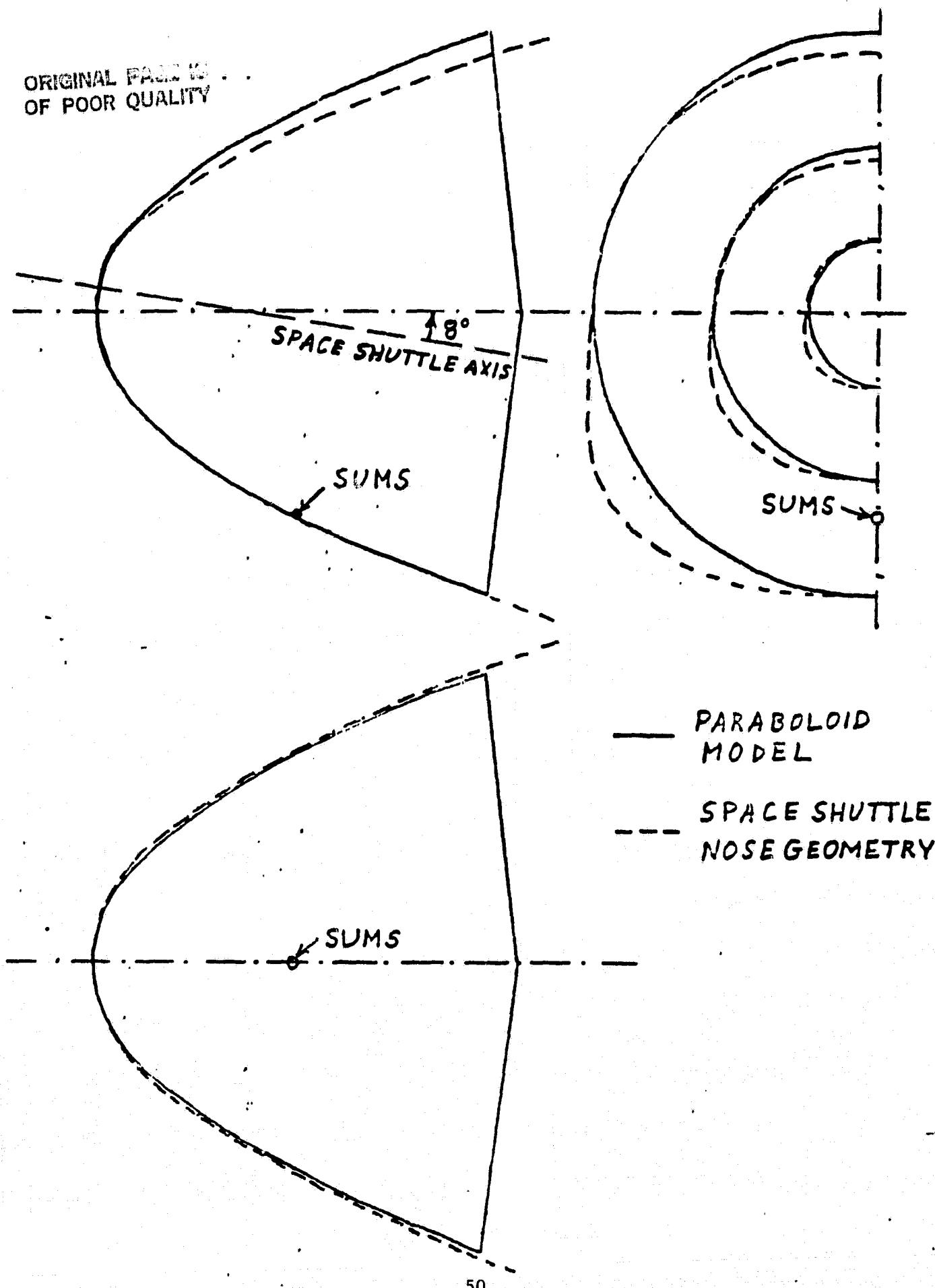


FIGURE 3. MODELING OF SPACE SHUTTLE NOSE (PARABOLA)



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FIGURE 4. TYPICAL CELL CONFIGURATION FOR HYPERBOLA

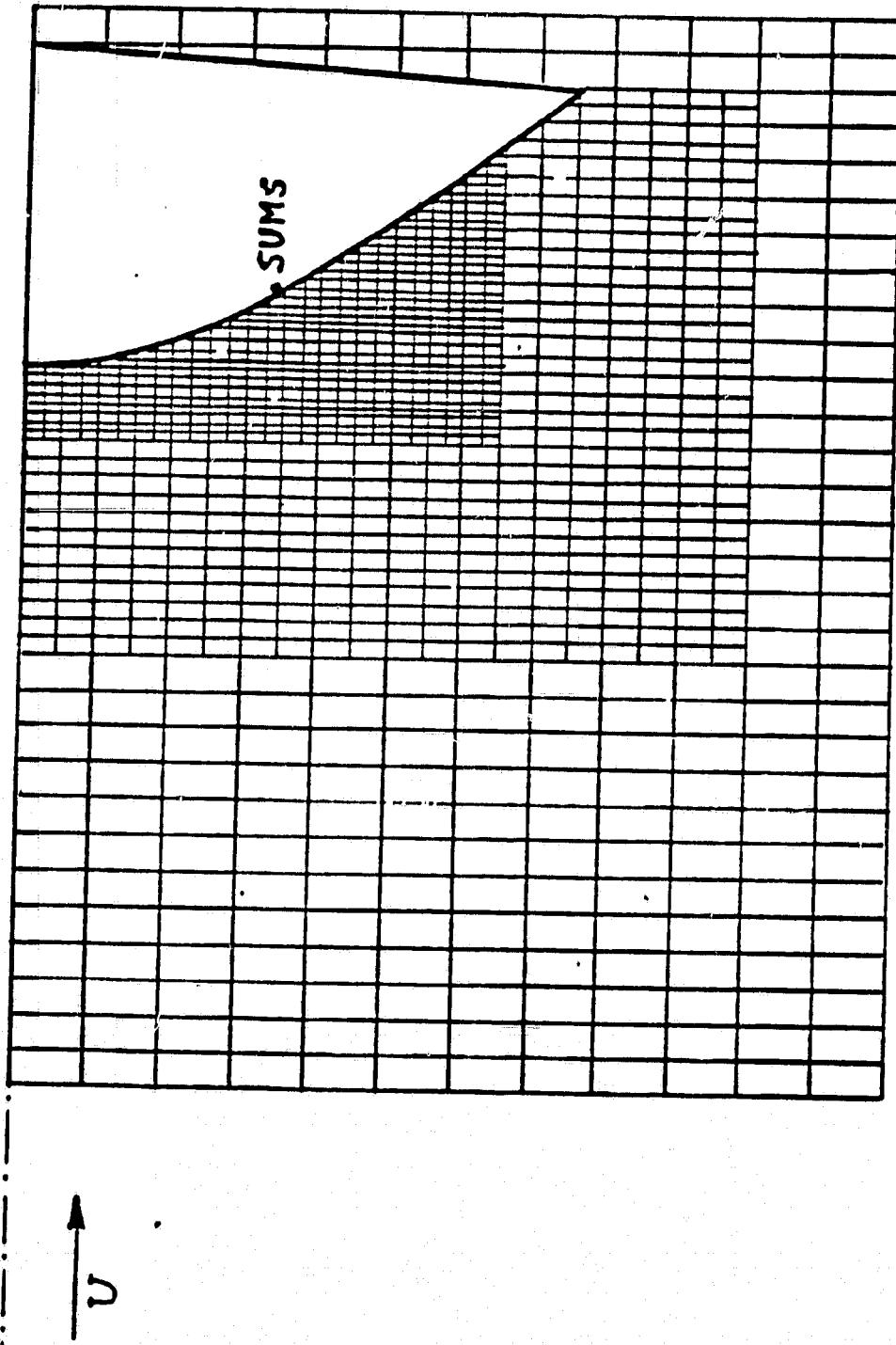


FIGURE 5. MODELLING OF VISCOSITY-TEMPERATURE DEPENDENCE

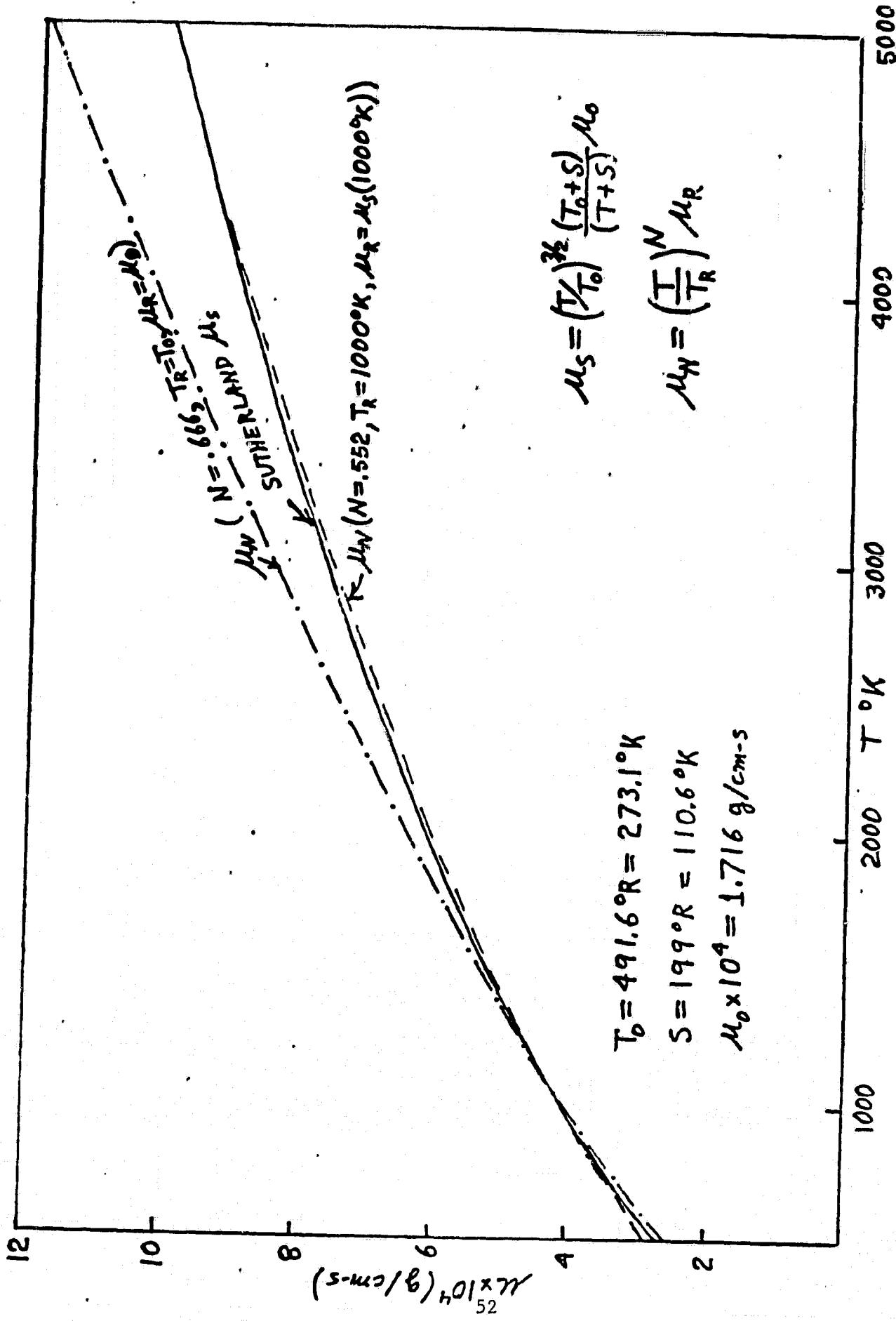
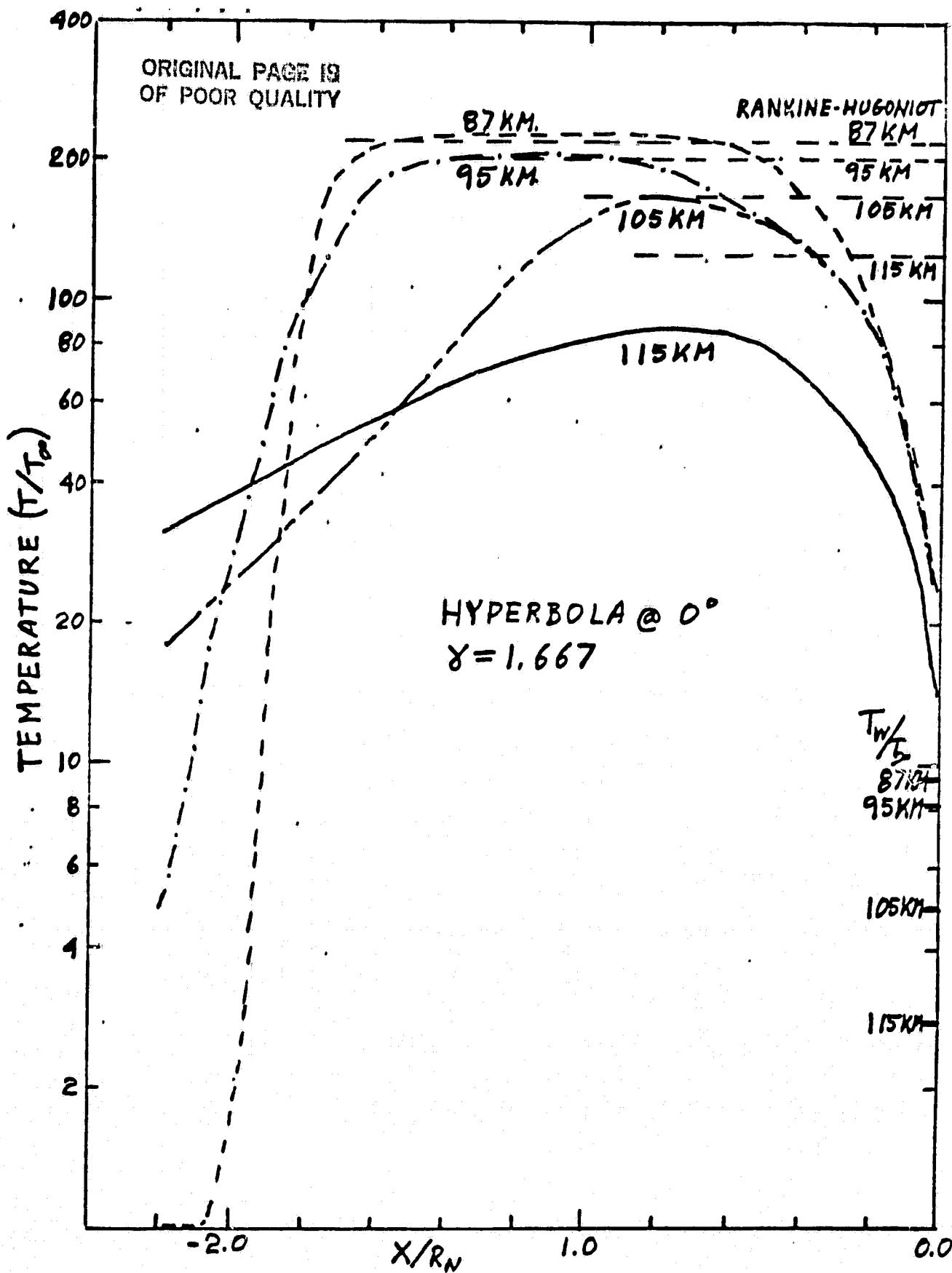


FIGURE 6. TEMPERATURE DISTRIBUTIONS (EFFECT OF ALTITUDE)



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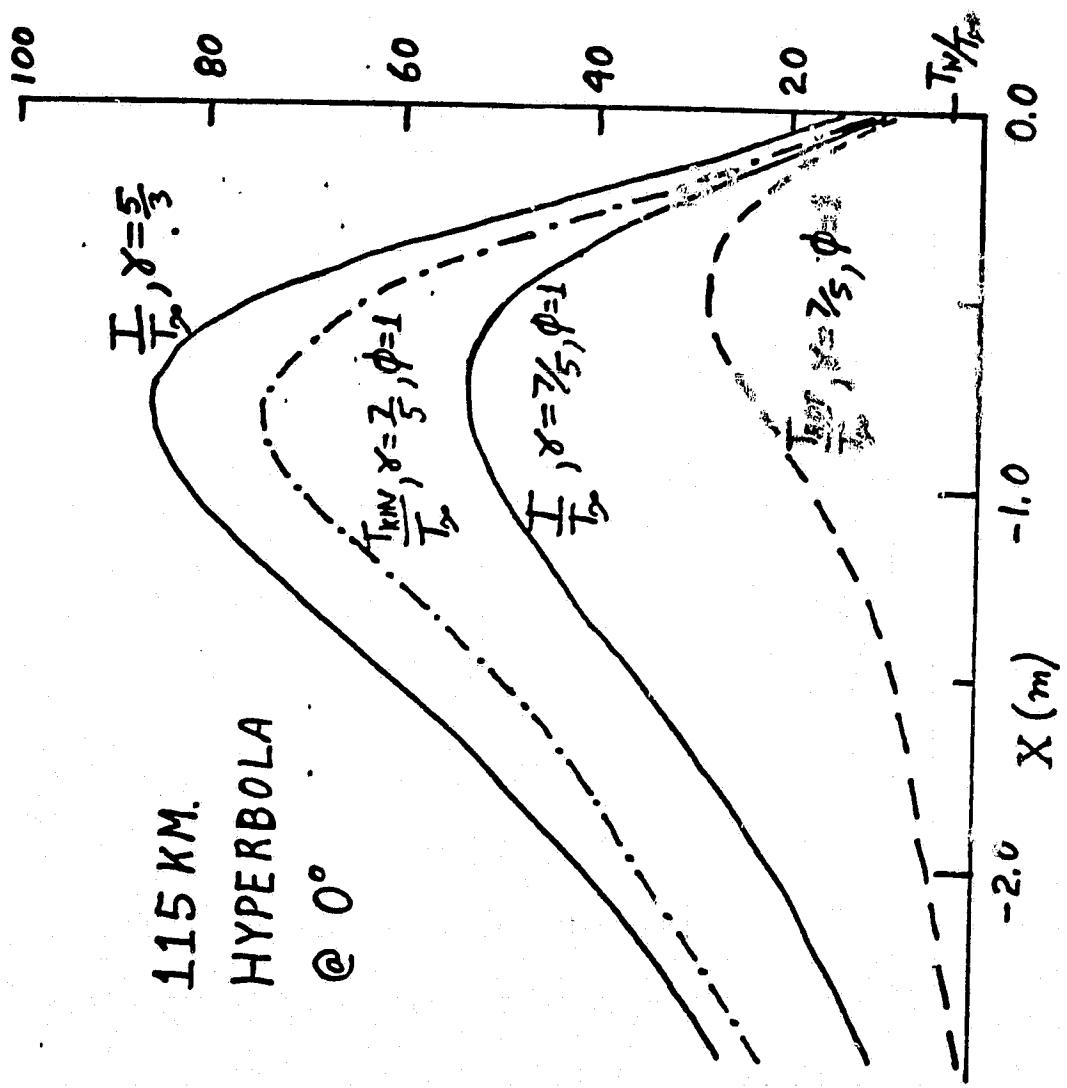


FIGURE 7. TEMPERATURE DISTRIBUTION ALONG STAGNATION STREAMLINE
(EFFECT OF γ)

FIGURE 8. SURFACE FLUXES (EFFECT OF γ)

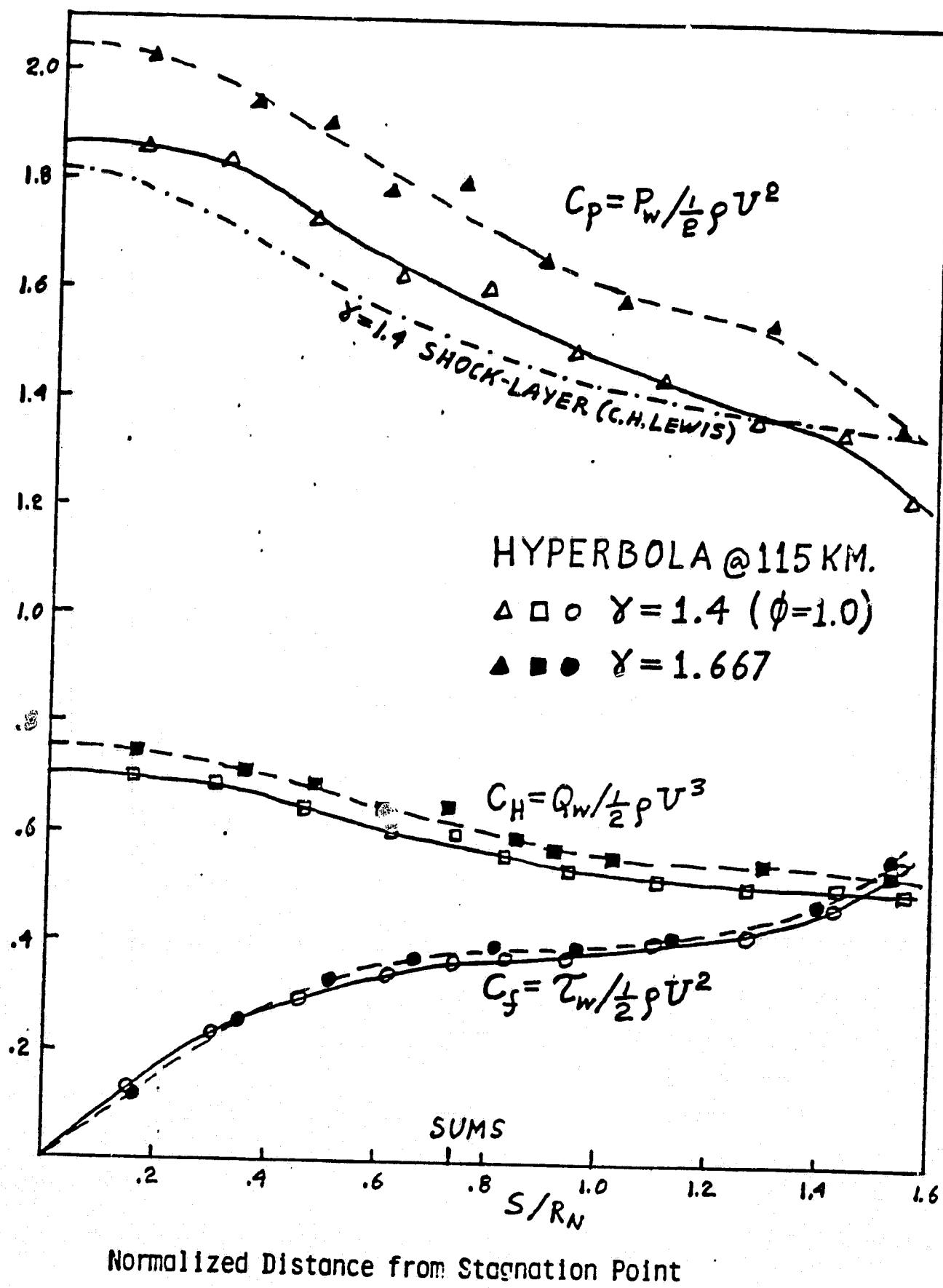
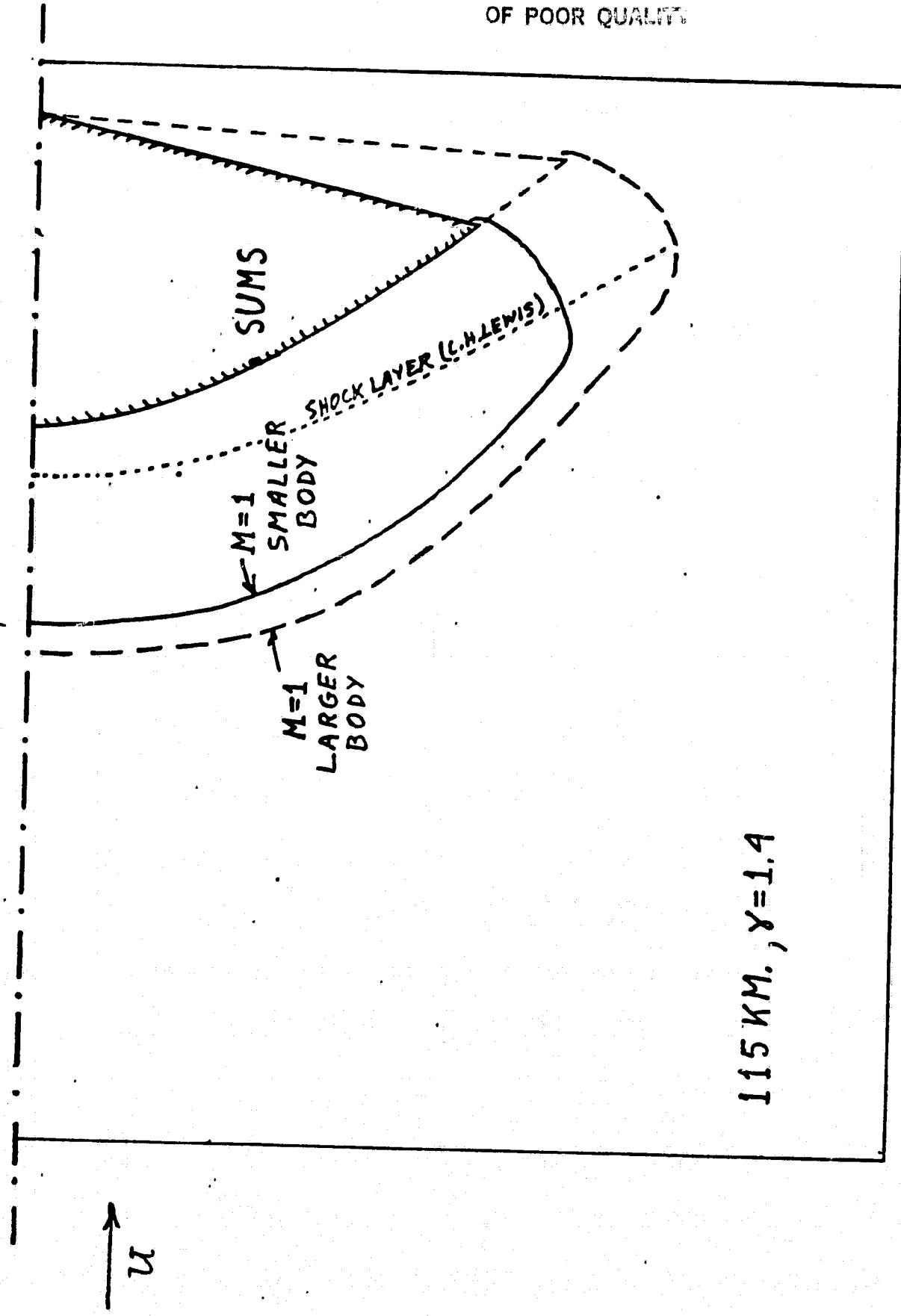


FIGURE 9. SONIC LINES (EFFECT OF BODY SIZE)

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115 KM. , $\gamma = 1.4$

FIGURE 10. SONIC LINES (EFFECT OF GEOMETRIC MODEL)

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32°

SUMS

PARABOLA
@ 32°

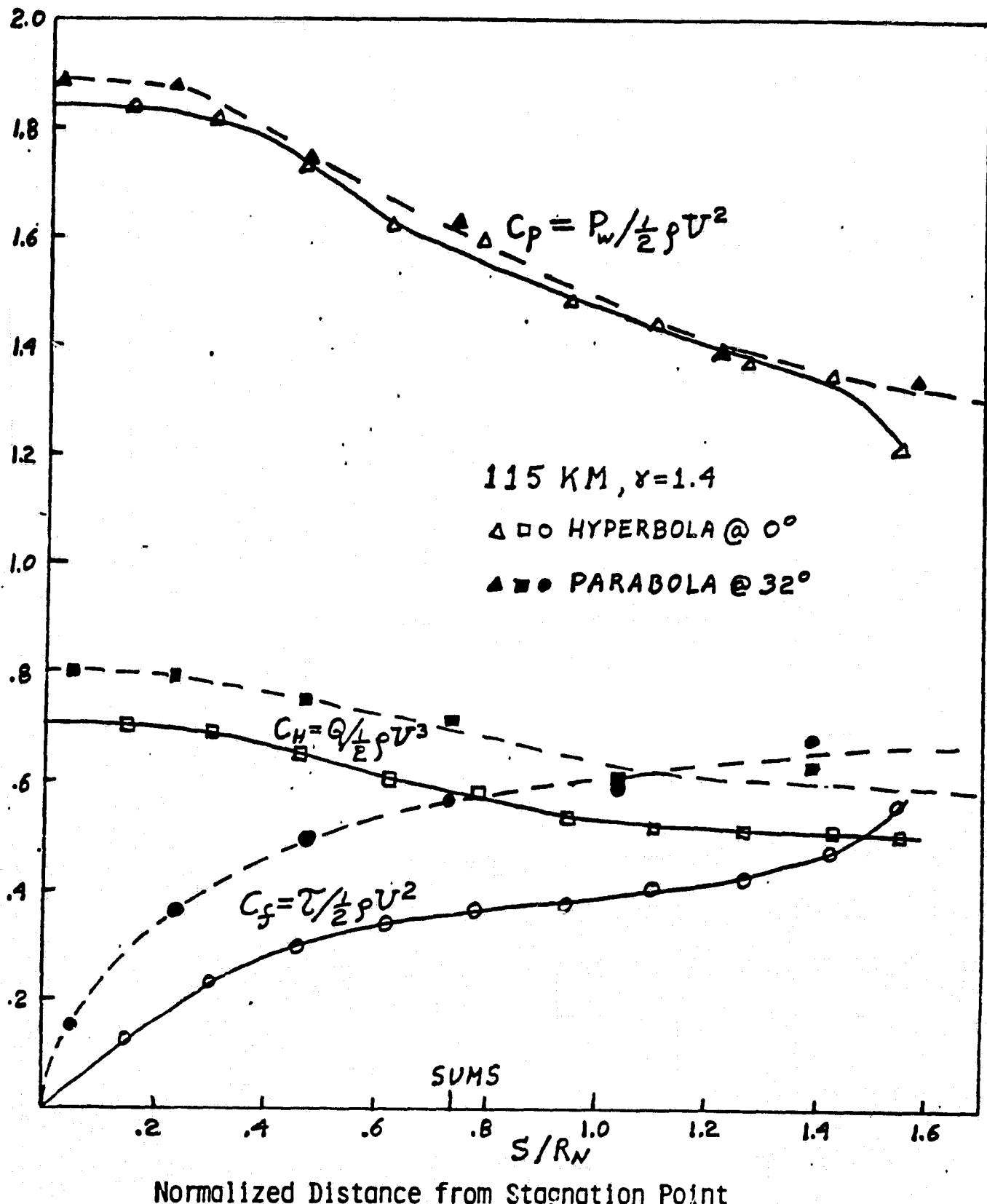
$M=1$

HYPERBOLA @ 0°

115 KM, $\gamma = 1.4$

U

FIGURE 11. SURFACE FLUXES (EFFECT OF GEOMETRIC MODEL)



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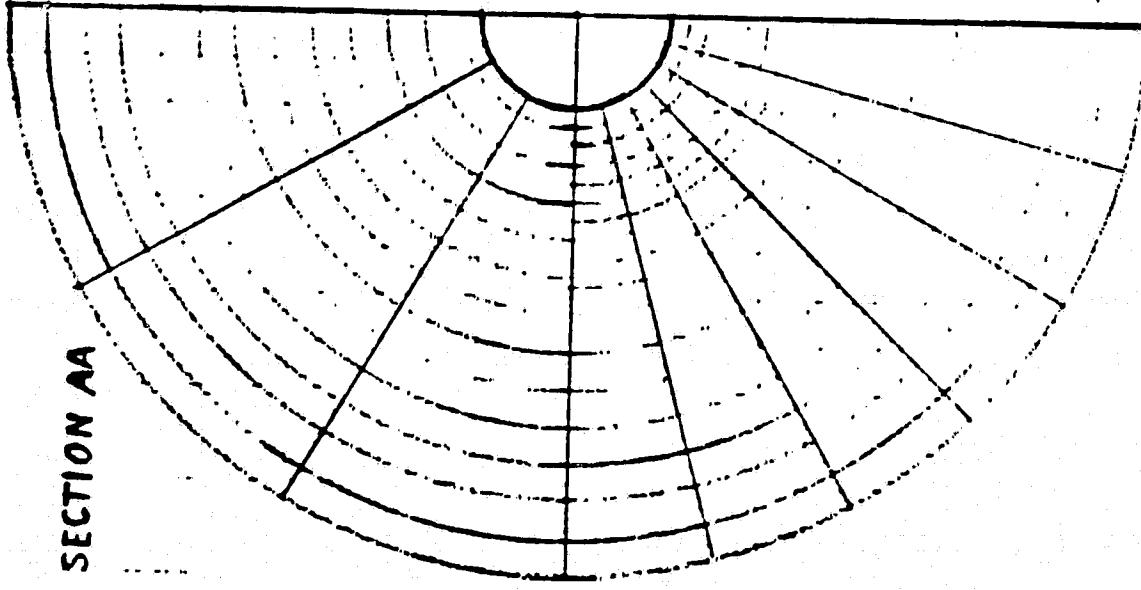
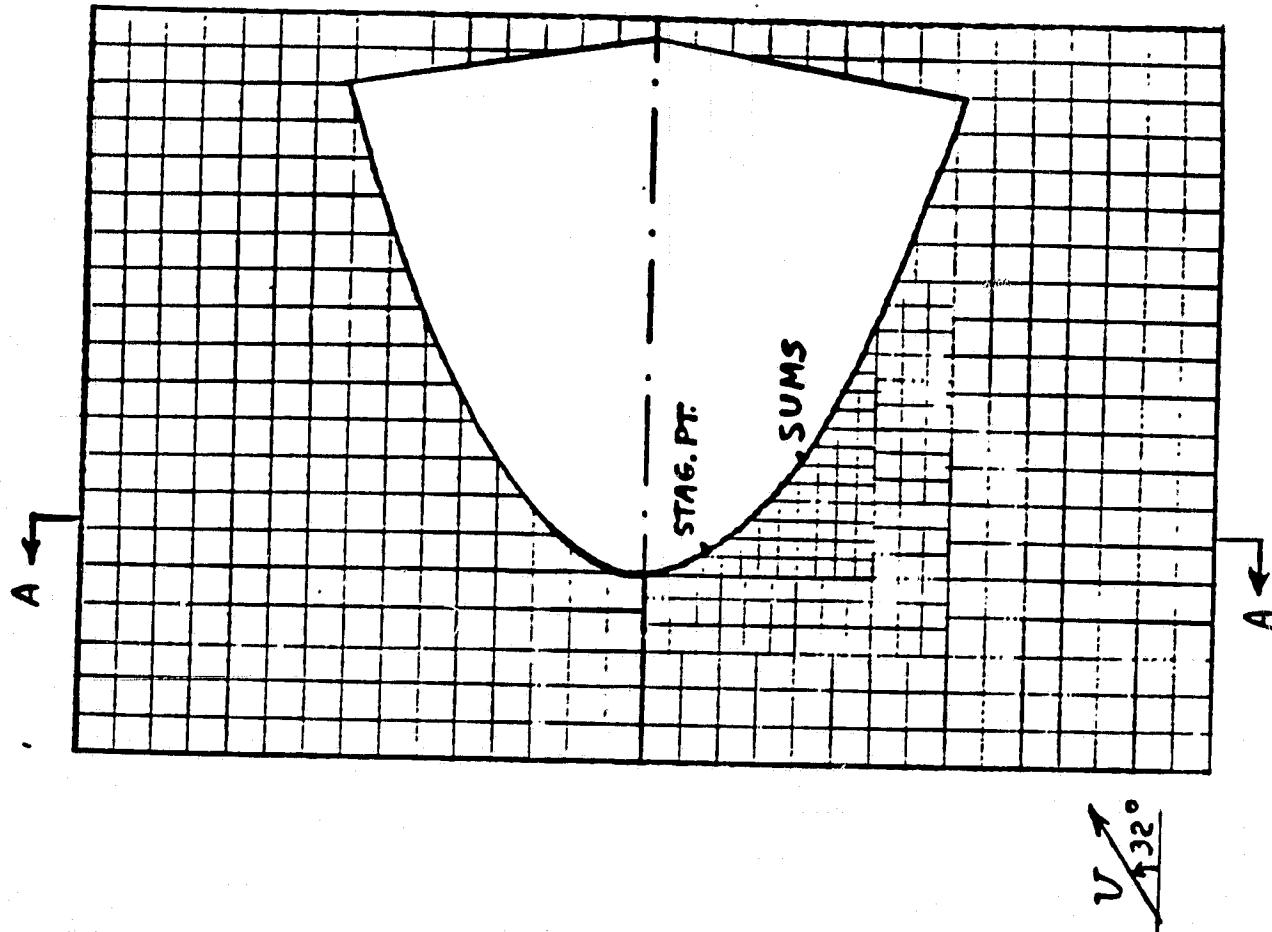


FIGURE 12. TYPICAL CELL CONFIGURATION (95 Km run)

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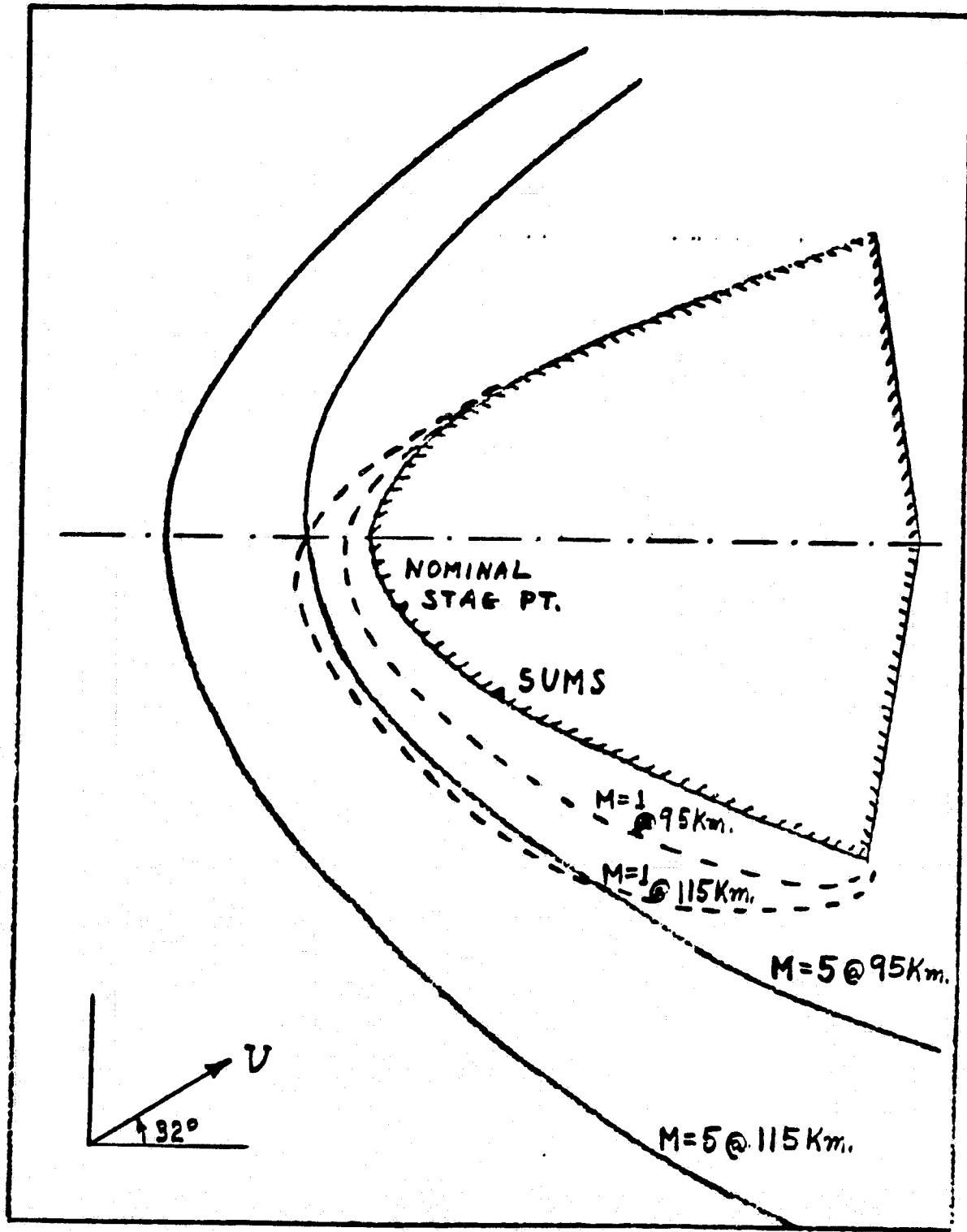


FIGURE 13. MACH NUMBER CONTOURS IN SYMMETRY PLANE AT 95 Km
and 115 Km

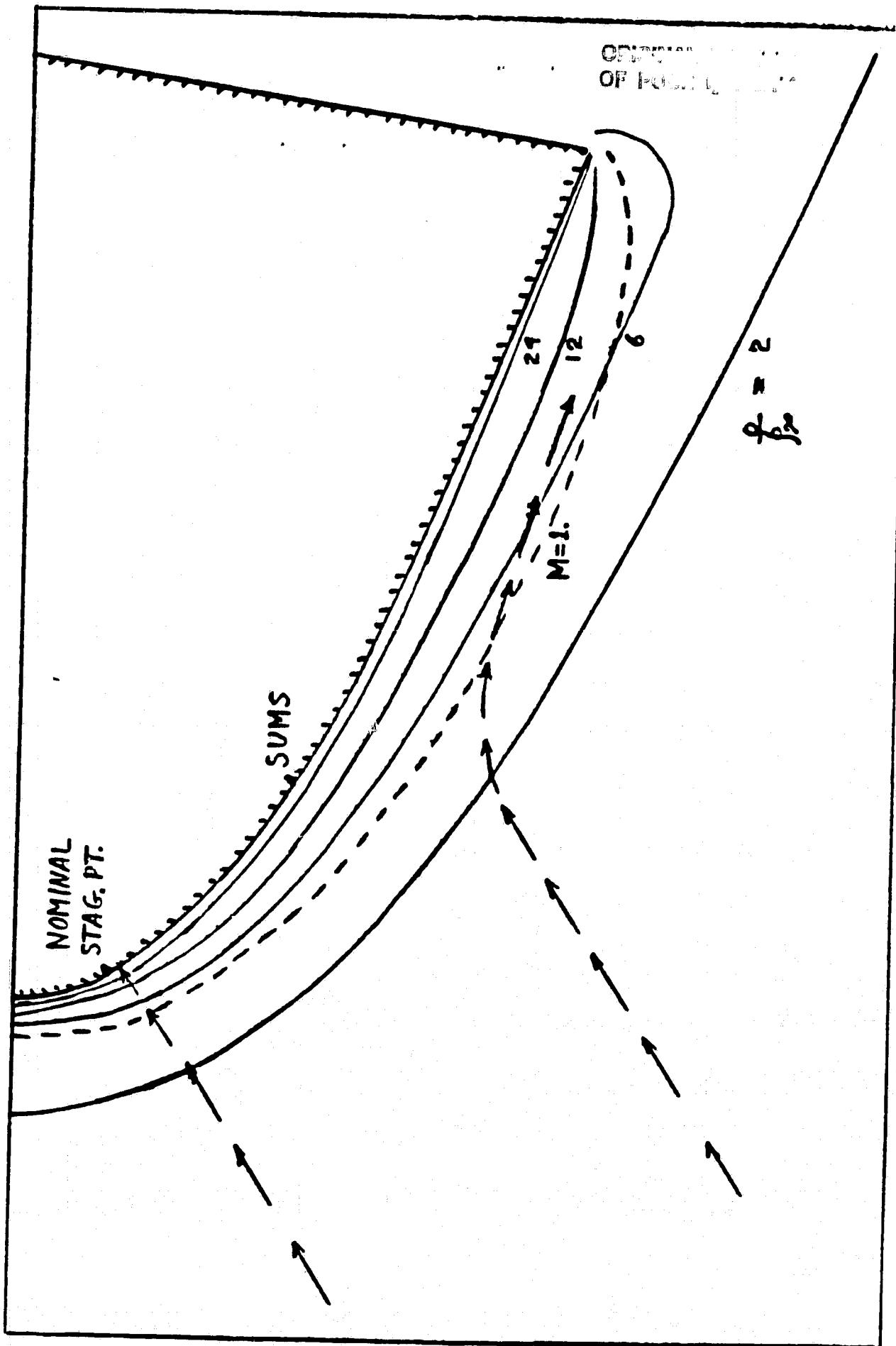
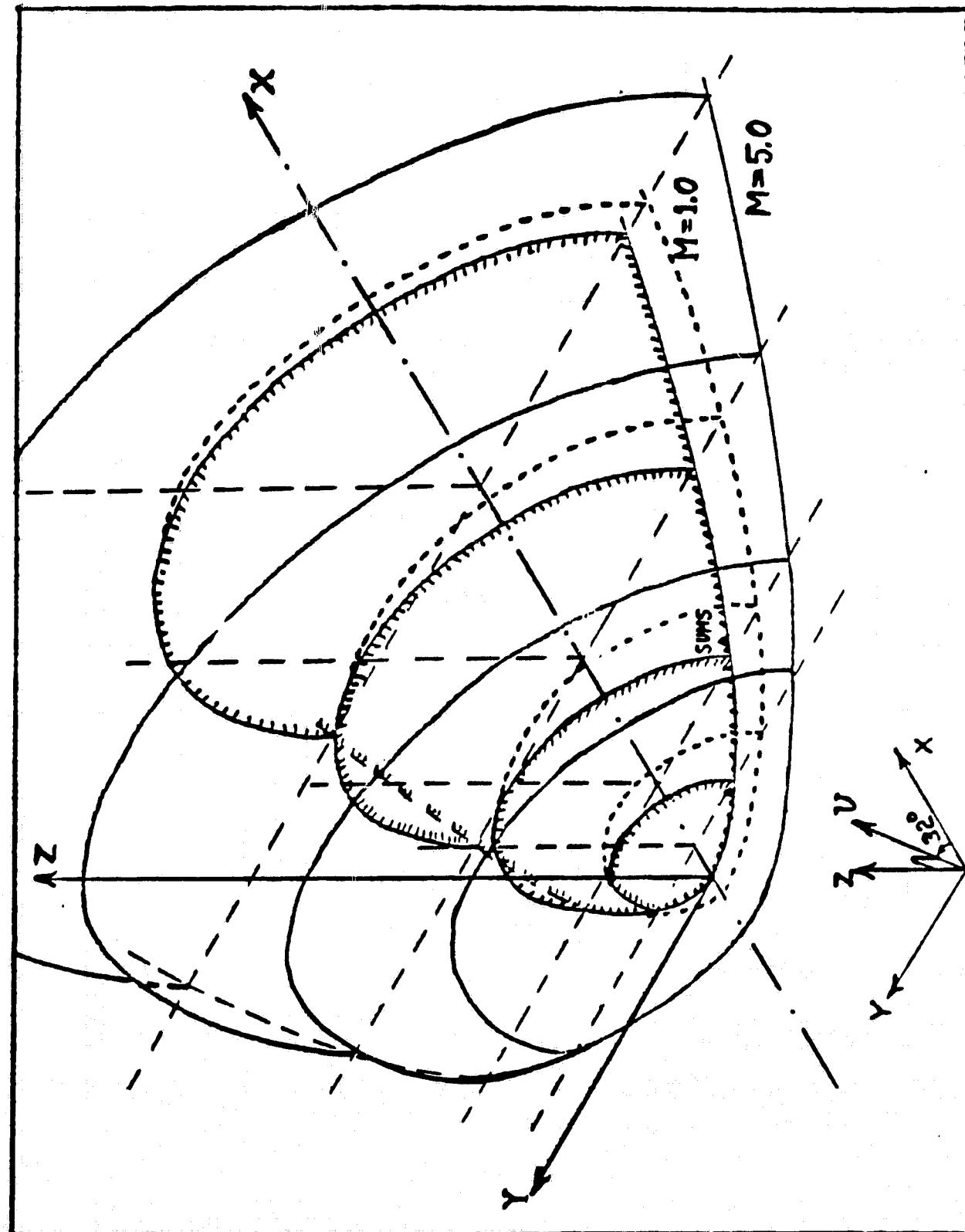


FIGURE 14. DENSITY CONTOURS ON WINDWARD SIDE AT 95 Km ALTITUDE

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MOVING CLOUDS DUE TO MACH NUMBER CINTOURS
AT 45 Km ALTITUDE



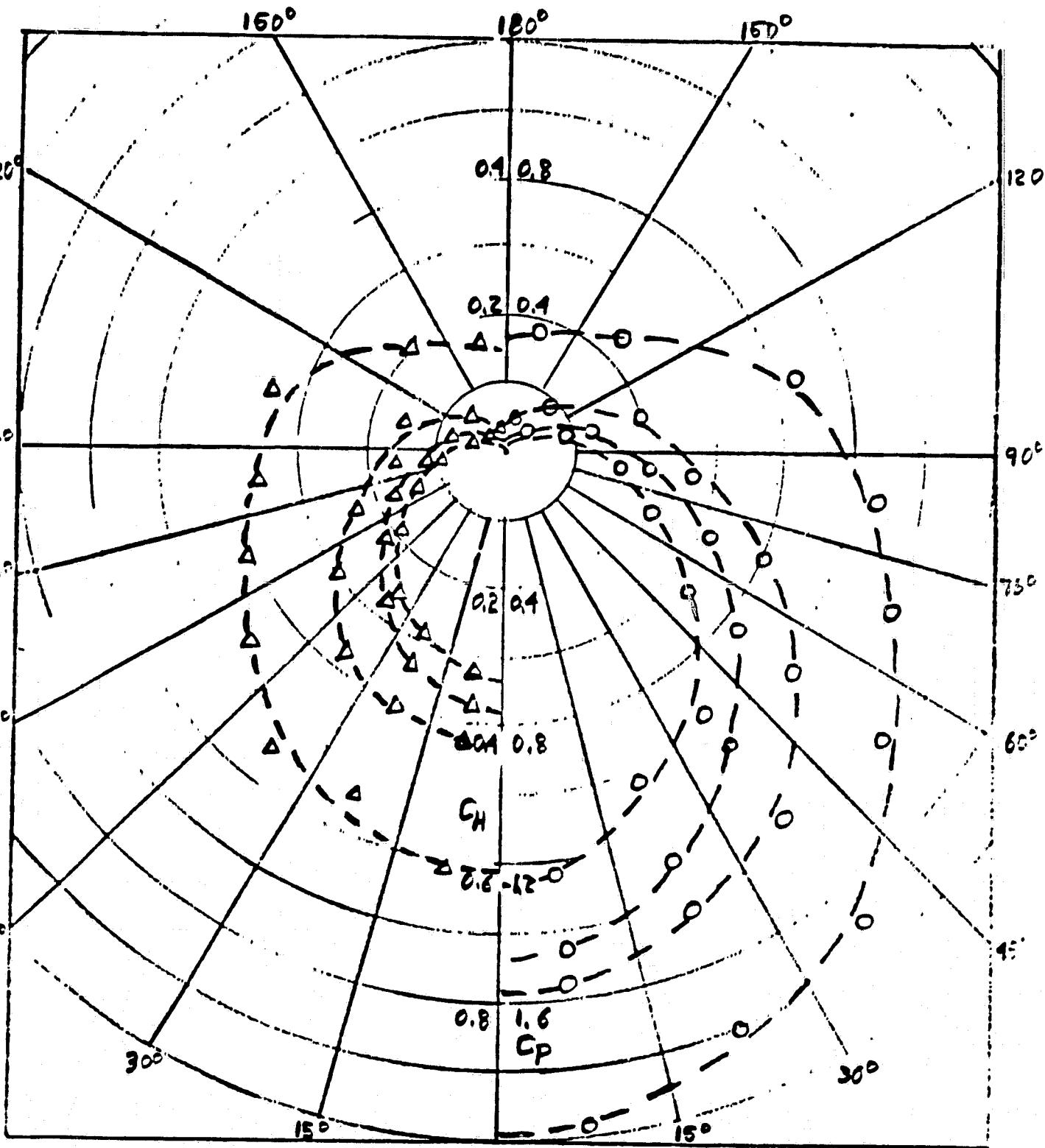


FIGURE 16. SURFACE PROPERTIES; PRESSURE COEFFICIENT $C_p = p_s / 1/2 \rho U^2$ AND HEAT TRANSFER COEFFICIENT $C_H = Q / 1/2 \rho U^3$ IN FOUR CROSS-PLANES AT 95 Km ALTITUDE

FIGURE 17. SURFACE PROPERTIES; PRESSURE C_p AND HEAT TRANSFER COEFFICIENT C_h VERSUS ANGLE IN CROSS-PLANE $x = 75$ cm (NEAR THE SUMS ENTRANCE) AT 95 km ALTITUDE

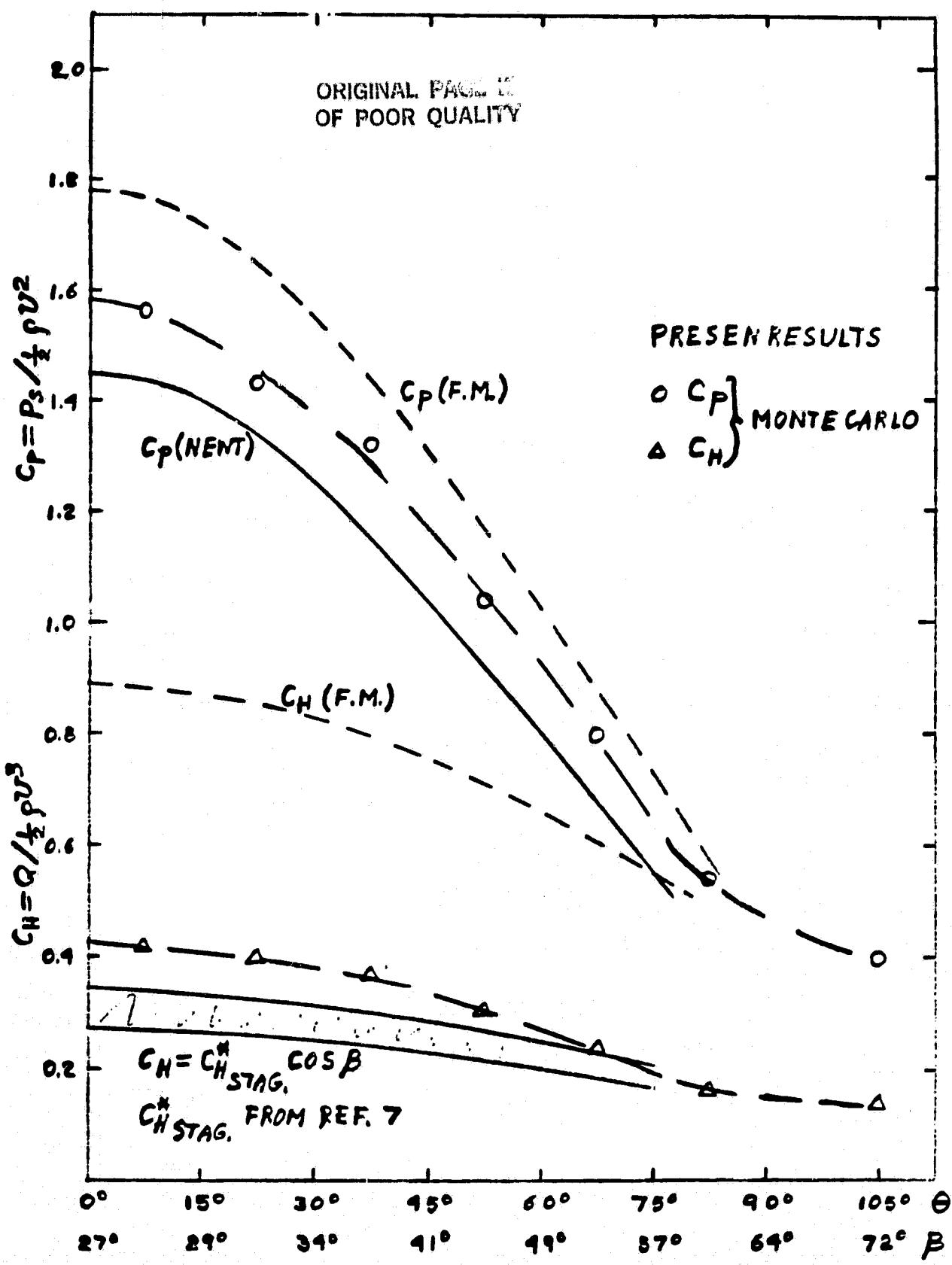


FIGURE 18. FLUX DISTRIBUTION FUNCTION AT SUMS INLET

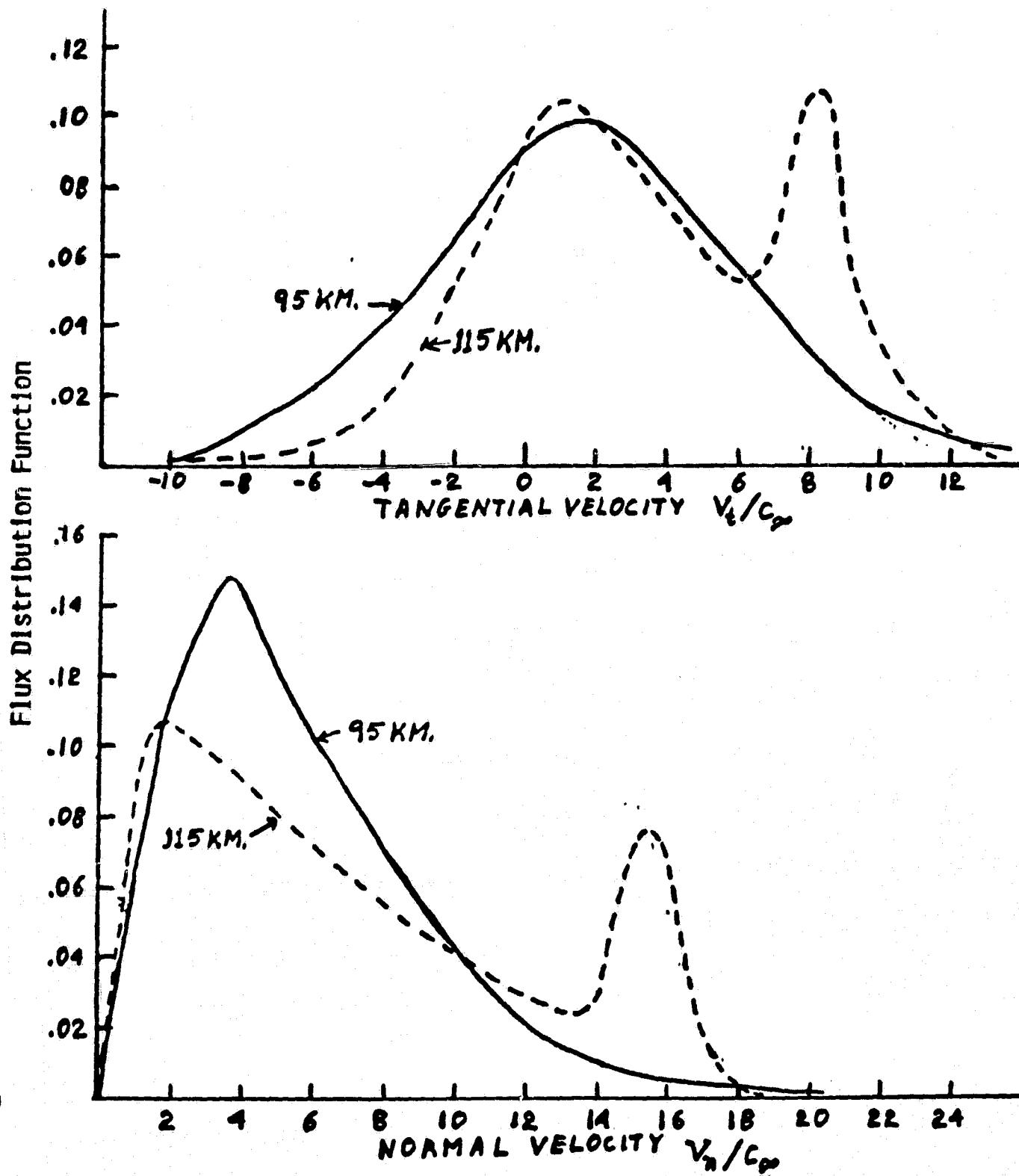
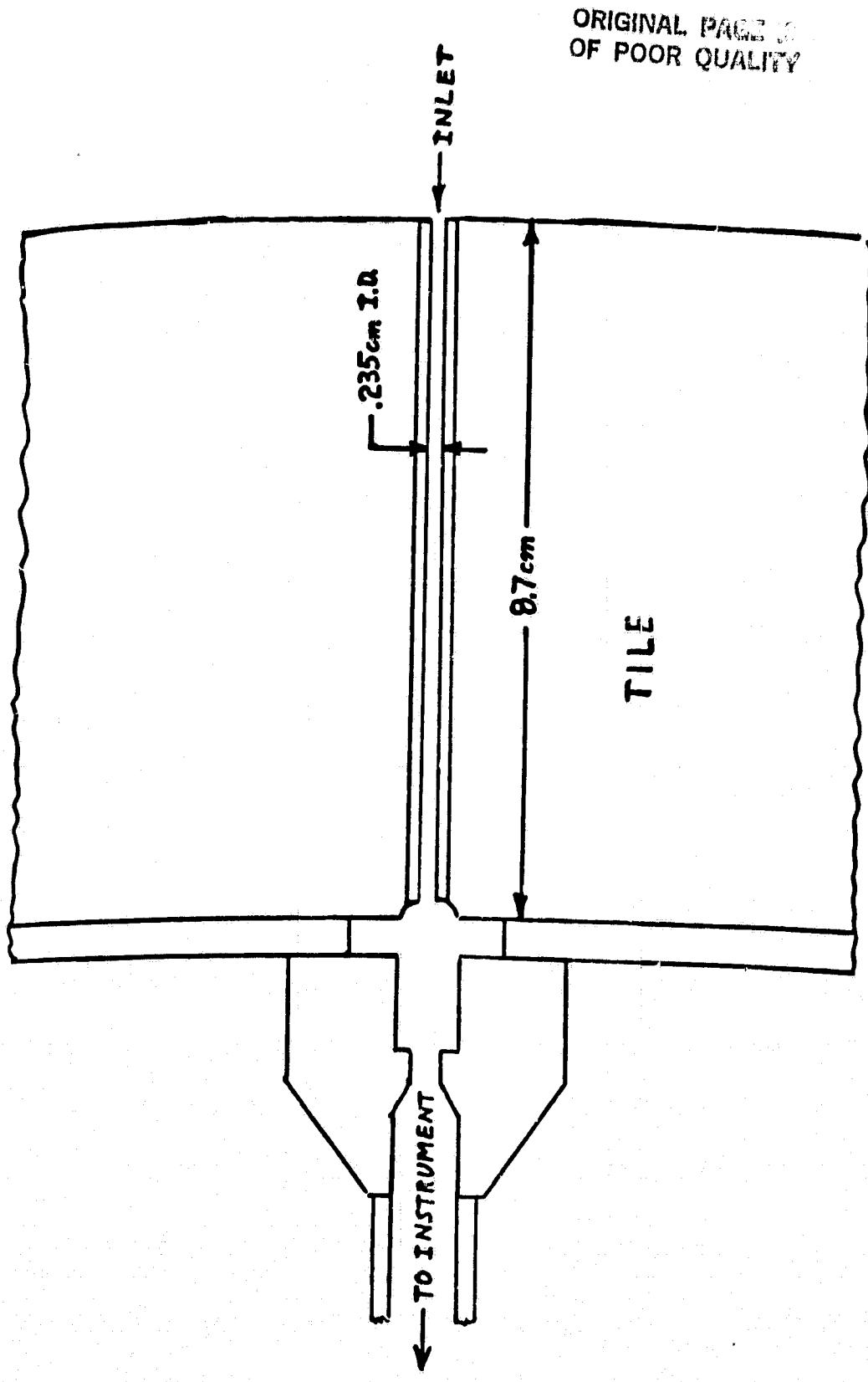
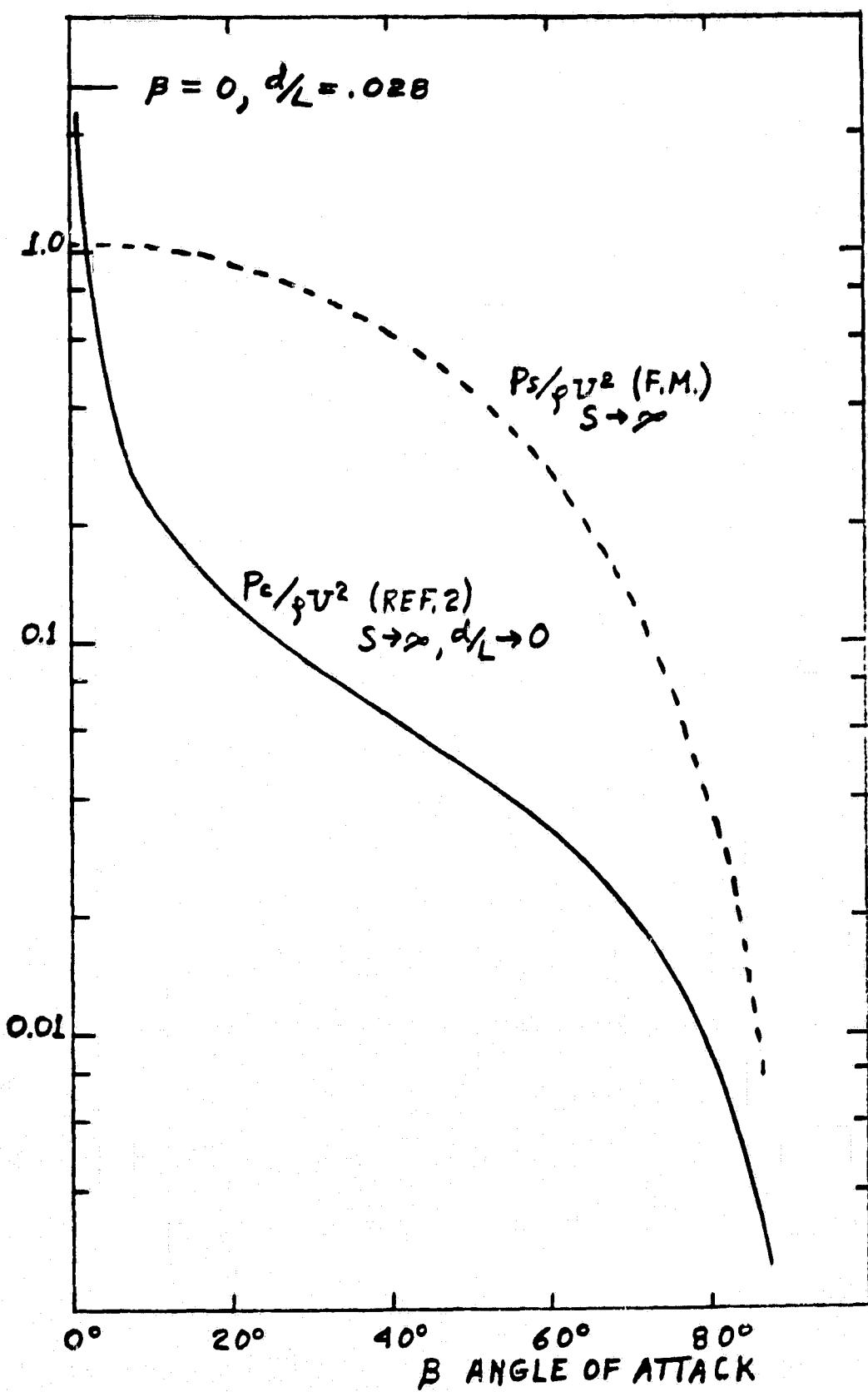


FIGURE 19. SUMS INLET GEOMETRY



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FIGURE 20. FREE MOLECULAR PROBE RESPONSE ($S \rightarrow \infty$, $d/L \rightarrow 0$)



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FIGURE 21. PRESSURE CORRECTION FACTOR (p_c/p_s)

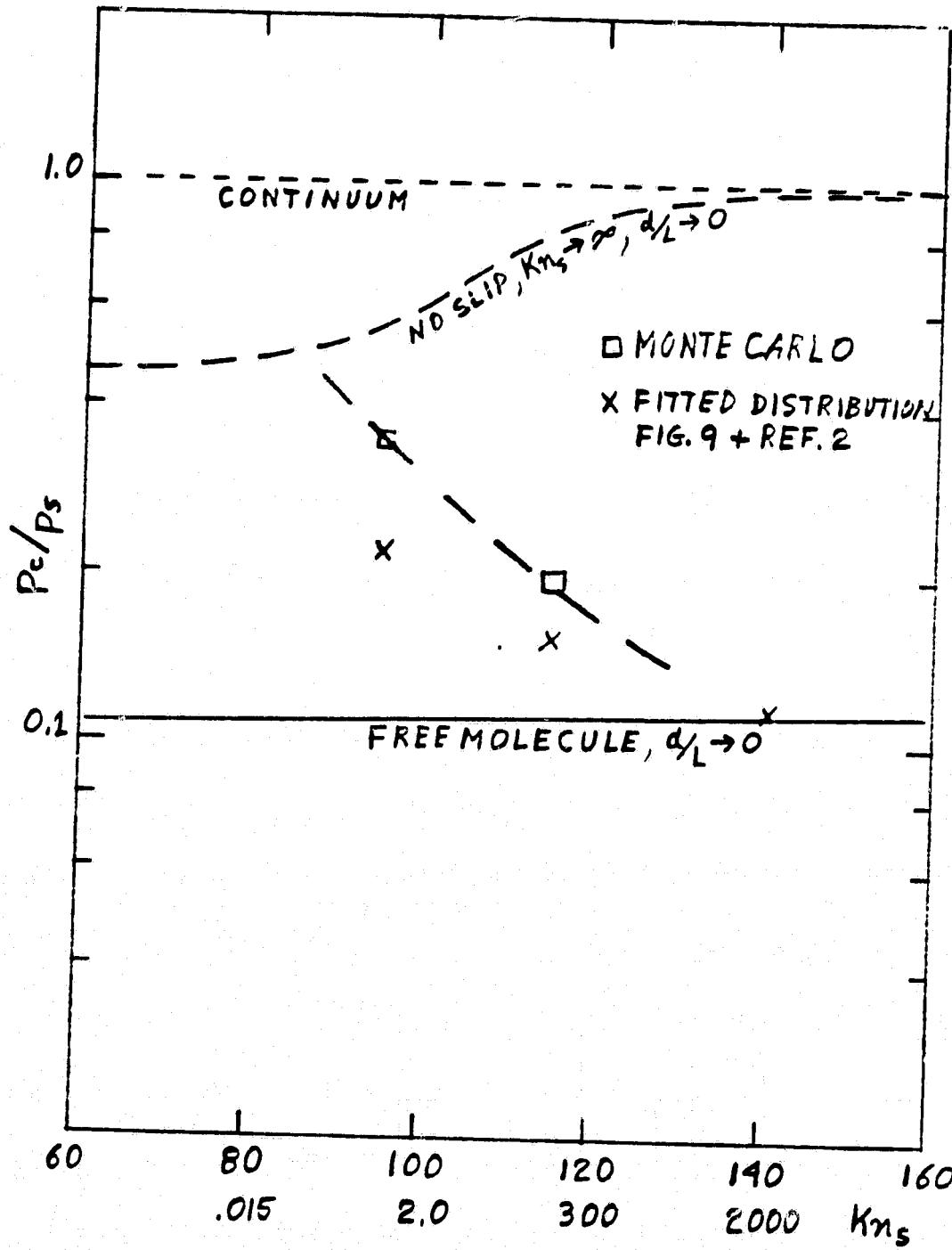


FIGURE 22.

SURFACE PRESSURE AT SUMS INLET
VERSUS ALTITUDE

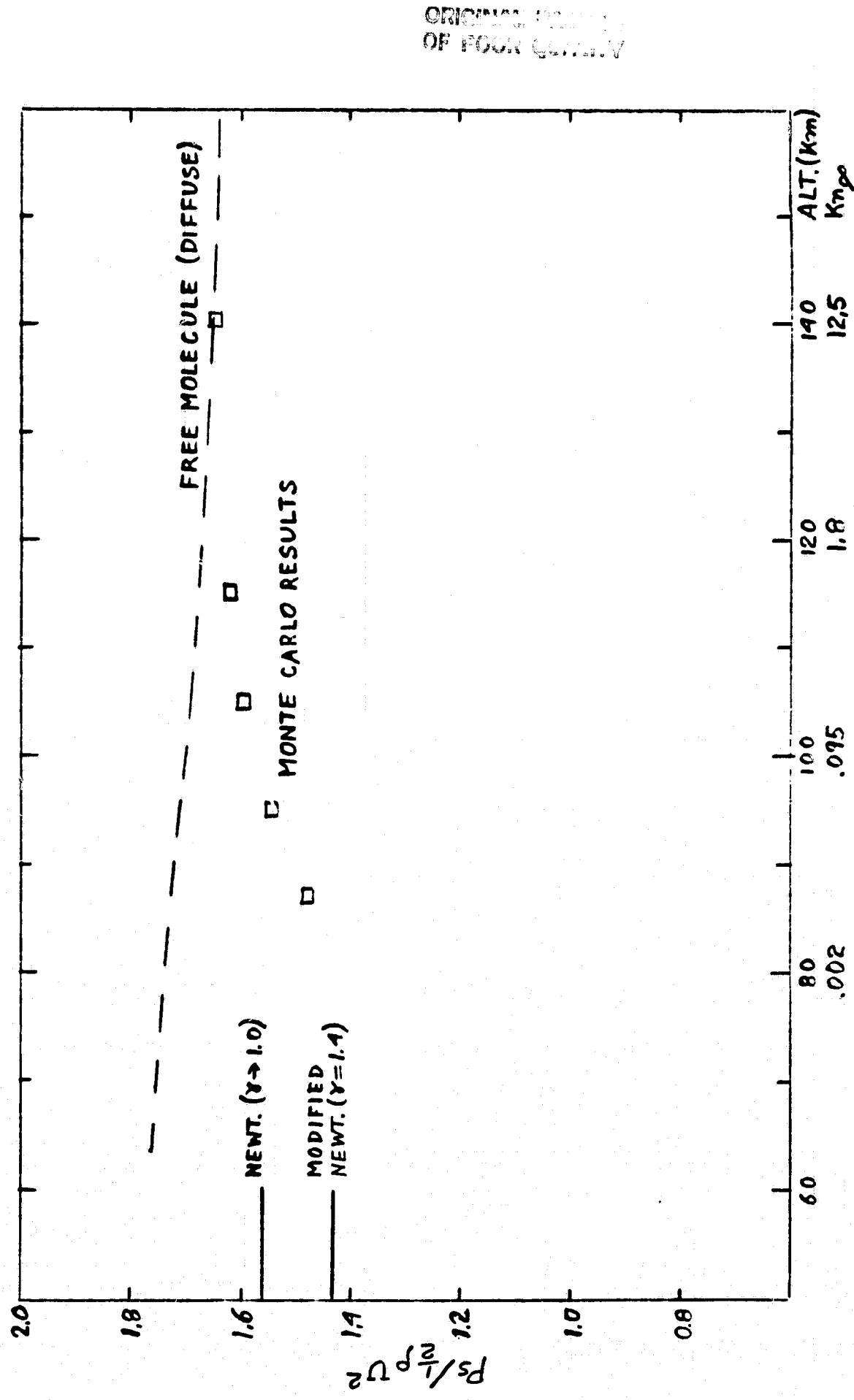
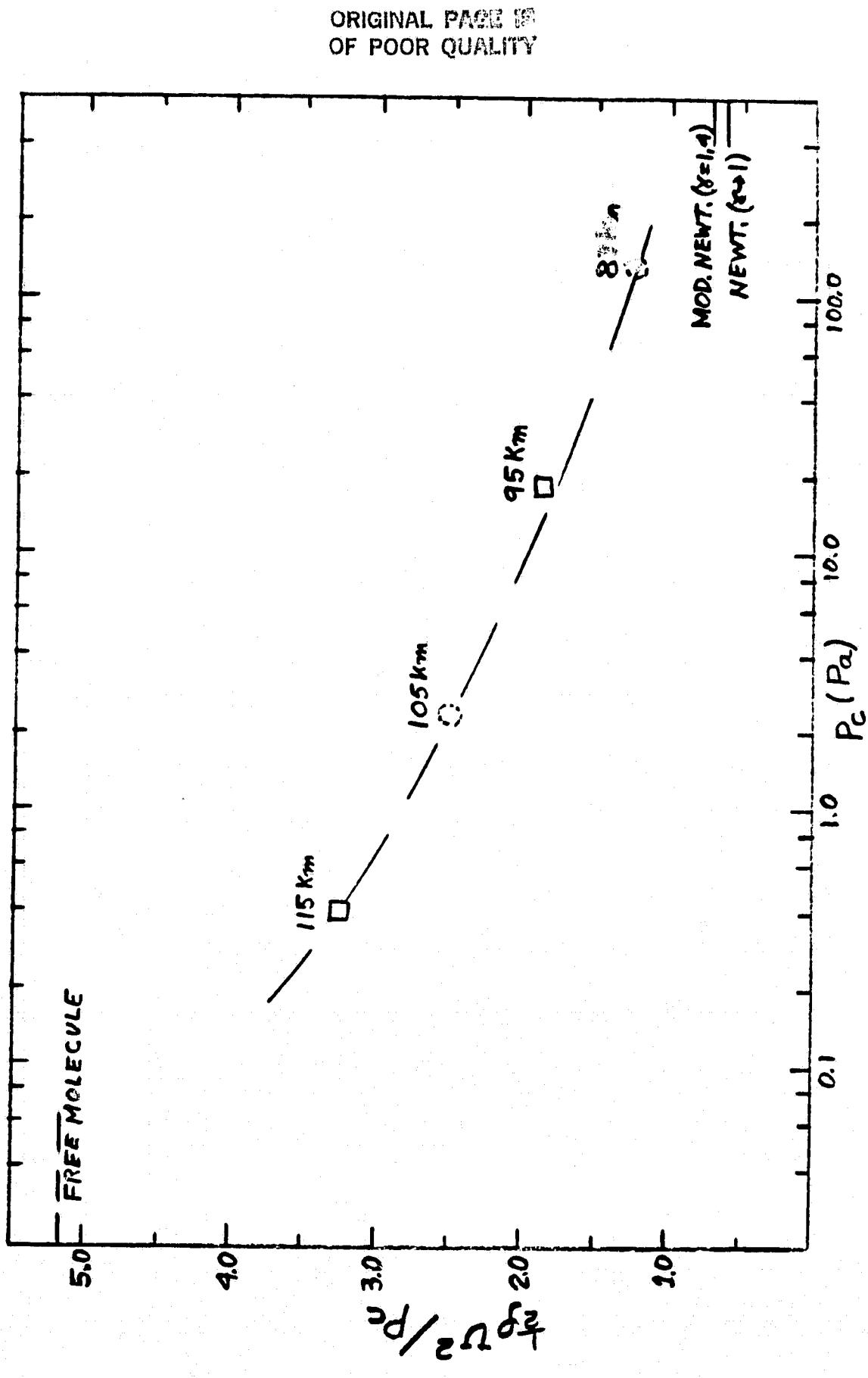


FIGURE 23. PRELIMINARY ESTIMATE OF DATA REDUCTION
RELATION q/P_c VERSUS P_c



APPENDIX A
COMPUTER CODE EXTERNAL

```
/* JOB GKB 0367425.GKBSPACE N=EXTCOMP REG=500 T=.8 P=150 C=0
/*FORMAT PR,DDNAME=SYMSG,DEST=VM370
/* EXEC FORTXCL,PARM.FORT='XREF'
/*FORT.SYSIN DD *
```

MAIN PROGRAM FOR MONTE CARLO 3-D EXTERNAL FLOW CALCULATIONS
 OBJECTIVE OF THIS MAIN PROGRAM IS TO SET THE DIMENSIONS
 MAIN RUNNING PROGRAM IS *** RUN ***

FOLLOWING TWO CARDS HAVE TO BE ELIMINATED FOR NON IBM MACHINES

```
*****  

  INTEGER*2 LB,NBM,NBN,NB,NBF,NBT,NBS,NUMCEL  

  INTEGER*2 LM,LPP,LCOL,LKW
```

THE NEXT CARD IS ASSOCIATED WITH PRINCETON RANDOM NUMBER GENERATOR

```
*****  

  COMMON/RANCOM/NRAN(4)
```

THE FOLLOWING DIMENSION STATEMENTS SET THE MAJOR ARRAY DIMENSIONS
 AND MUST BE CONSISTENT WITH THE FOLLOWING DATA CARD -----

NSP=NUMBER OF SPECIES - EXAMPLE BELOW NSP=1
 NMB=NUMBER OF BOXES WITHOUT COUNTING SUBDIVISIONS OR THOSE
 OCCUPIED BY THE BODY - EXAMPLE BELOW NMB=2600
 NMC=NUMBER OF FINAL CELLS - EXAMPLE BELOW NMC=1500
 NMP=MAX NUMBER OF MOLECULES OF EACH SPECIES ALLOWED IN PROGRAM.
 IF EXCEEDED, PROGRAM EITHER FAILS OR RESTARTS AT BEGINNING
 WITH NUMBER REDUCED BY 10% - EXAMPLE BELOW NMP=20000
 NPB=MAXIMUM NUMBER IN EACH CELL - EXAMPLE NPB=100

```
DIMENSION DBA(1,1500),NB(1,1500),NBF(1,1500),NBT(1,1500),LKW(1500)
DIMENSION TMP(1,1500),TPA(1,1500),XV(1,1500),XVA(1,1500)
DIMENSION XV(1,1500),XVA(1,1500),ZV(1,1500),ZVA(1,1500),DB(1,1500)
DIMENSION TRP(1,1500),TRPA(1,1500),NBS(1,1500),NBM(1,1500)
DIMENSION NBN(1500),T(1,1,1500)
DIMENSION LB(20000),LM(1,20000),ER(1,20000)
DIMENSION LPP(1,20000),PAU(1,20000),PAV(1,20000),PAW(1,20000)
DIMENSION PAX(1,20000),PAY(1,20000),PAZ(1,20000),LCOL(1,20000)
DIMENSION FNB(2600),XC(2600),YC(2600),ZC(2600),NUMCEL(2600)
DATA NSP/1/,NMB/2600/,NMC/1500/,NMP/20000/,NPB/100/
```

2 FORMAT(/17X,'NORMAL TERMINATION OF THE PROGRAM')
 NAMELIST/DIM/NSP,NMB,NMC,NMP,NPB,NRAN

INITIALIZATION OF RANDOM NUMBER GENERATOR - PRINCETON ROUTINE

```
NRAN(1)=0
NRAN(2)=0
NRAN(3)=0
NRAN(4)=0
```

PRINTOUT OF MAJOR ARRAY DIMENSIONS USED ABOVE

ORIGINAL PAGE IS
 OF POOR QUALITY

WRITE(6,DIM)

ORIGINAL PAGE
OF POOR QUALITY

CALL OF MAIN OPERATING PROGRAM WHICH REQUIRES INPUTS:
 &ECONTROL, &TIMES, &FLOREF, &MOLEC, &SHAPES, &GEOM, &COUPLE
 THESE INPUTS ARE ALL CURRENTLY IN THE NAMELIST FORMAT
 AND MAY HAVE TO BE CHANGED IF THAT CONVENTION IS NOT AVAILABLE
 BRIEF DESCRIPTION OF THE PARAMETERS FOLLOWS

ECONTROL - ONE OCCURRENCE (NEW OR RESTART)

PARAMETER	DEFAULT	DEFINITION OR EXPLANATION
NAME	8 BLANKS	ANY ALPHANUMERIC NAME UP TO 8 CHARACTERS
TITLE	24 BLNKS	ANY ALPHANUMERIC TITLE UP TO 24 CHARACTERS
PERCNT	.001	ACCURACY IN INTEGRATION PROCEDURES
ICOPY	1	NUMBER OF ADDITIONAL COPIES OF OUTPUT
DUMP	.TRUE.	IF TRUE WILL CAUSE SYSTEM DUMP FOR ANY OF 12 PROGRAMMER DESIGNED ERROR HALTS.
DEBUG (1)	.FALSE.	IF TRUE WILL PRINT MESSAGE WHEN CELL POP. EXCEEDS MNB
DEBUG (2)	.FALSE.	IF TRUE WILL PRINT CPU TIME AROUND EACH PART OF LOOP
DEBUG (3)	.TRUE.	IF TRUE WILL PRINT CPU TIME REMAINING AT END OF LOOP
NEW	.TRUE.	IF TRUE - NEW RUN - IF FALSE - RESTART OF RUN
SAVE	.FALSE.	IF TRUE - SNAPSHOT SAVED ON TAPE(9) FOR RESTART
REDO	.FALSE.	IF TRUE PROGRAM WILL AUTOMATICALLY RESTART WITH 90% OF TOTAL IF TOTAL CELL POPULATION EXCEEDS MNM

&TIMES - ONE OCCURRENCE (NEW OR RESTART)

PARAMETER	DEFAULT	DEFINITION OR EXPLANATION
DTM	---	REAL NUMBER - FRACTION OF MEAN FREE TIME PER CYCLE
ITS	---	INTEGER - NUMBER OF CYCLES PER SAMPLE
ITP	---	INTEGER - NUMBER OF CYCLES BETWEEN PRINTOUTS
TST	---	INTEGER - ESTIMATE OF NUMBER OF CYCLES TO STEADY STATE
TLIM	---	INTEGER - TOTAL NUMBER OF CYCLES TO END OF RUN - WILL TERMINATE SOONER IF CPU TIME IS TO BE EXCEEDED

&FLOREF - ONE OCCURRENCE (NEW RUN ONLY)

PARAMETER	DEFAULT	DEFINITION
INM	---	INITIAL NUMBER OF MOLECULES IN M< MNM > OR = NMP
MNM	---	MAXIMUM NUMBER OF MOLECULES PER SPECIES
MNB	---	MAXIMUM NUMBER PER CELL - DIAGNOSTIC ONLY
MSP	---	NUMBER OF MOLECULAR SPECIES (MAX. IS 3)
MET	0	IF 0 - DATA IS IN SI (METRIC) UNITS IF >0 - DATA IS IN ENGLISH UNITS
U	---	FLOW VELOCITY (M/SEC) OR (FT/SEC)
ANGLE	---	ANGLE OF ATTACK (DEGREES)
RNU	0.0	ARRAY GIVING MOLE FRACTIONS OF SPECIES IN FREE STREAM
RMA	0.0	ARRAY GIVING MOLECULAR WEIGHTS OF SPECIES ABOVE
TF	---	FREE STREAM TEMPERATURE (K OR R)
DENF	---	FREE STREAM NUMBER DENSITY (NUM/M**3 OR NUM/FT**3)

&MOLEC - ONE OCCURRENCE (NEW RUN ONLY)

PARAMETER	DEFAULT	DEFINITION
TRF	---	REFERENCE TEMPERATURE FOR MOLECULAR DATA
DIR	0.0	CROSS-SECTIONS AT REFERENCE TEMP. (MSPXMSP)
ETA	0.0	PARAMETERS IN DIFFUSION AND VISCOSITY LAW (MSPXMSP)
PHI	0.0	PARAMETERS FOR ROTATIONAL RELAXATION (MSPXMSP)
CHI	0.0	ROTATIONAL DEGREE OF FREEDOM PARAMETER (NROT/2 - 1)
ACR	.001	ACCURACY IN MOLECULAR COLLISION CALCULATIONS

&SHAPES - ND+1 OCCURRENCES WHERE ND=NUMBER OF BODY SEGMENTS (NEW RUN)
 PARAMETER DEFAULT DEFINITION
 FIRST OCCURRENCE
 BODY(1) 0.0 STARTING POINT OF BODY FROM FRONT OF CELLS (M OR FT)
 BODY(I) I>1 - - NEED NOT BE SPECIFIED
 SUBSEQUENT OCCURRENCES (ND)
 BODY(1) - - - X COORDINATE FROM FRONT OF BODY OF THE DOWNSTREAM
 EDGE OF THE CURRENT BODY SEGMENT
 BODY(2) - - - TEMPERATURE OF THIS BODY SEGMENT
 BODY(3) - - - SURFACE AREA/TOTAL CROSS-SECTIONAL AREA FOR SEGMENT.
 IF 0.0 PROGRAM WILL COMPUTE THIS QUANTITY
 BODY(4) - - - SWITCH - IF 0.0 THIS SEGMENT'S EQ. WILL APPEAR LATER
 IF >0.0 THE EQ. OF THIS AND PRECEDING SEGMENTS IS
 GIVEN BY BODY(6+2*MSP) TO BODY(9+2*MSP)
 SWITCH - IF NOT 0.0 THIS IS THE LAST SHAPES CARD
 BODY(I) I EVEN ALPHA - ENERGY ACCOMODATION COEFFICIENT FOR SPECIES
 BODY(J) J ODD SIGMA - TANGENTIAL ACCOMODATION COEFF. FOR SPECIES
 I AND J < (6+2*MSP)
 BODY(6+2*MSP) - ORIGIN OF COORDINATES WITH RESPECT TO BODY START
 FOR THE EQUATION OF THIS BODY SECTION
 BODY(7+2*MSP) - COEFFICIENTS A,B,C IN THE EQUATION
 BODY(9+2*MSP) R**2+A*X**2+B*X+C=0.0 FOR THIS BODY SECTION

&GEOM - ONE OCCURRENCE (NEW RUN ONLY)
 PARAMETER DEFAULT DEFINITION
 NWEDGE - - - TWO INTEGERS GIVING THE NUMBER OF WEDGES BELOW
 AND ABOVE THE ANGLE THETAZ
 THETAZ - - - ANGLE FROM NEG. Y AXIS DIVIDING DIFF. WEDGE SIZES
 RMB - - - MAX. BODY RADIUS - IF 0. WILL BE COMPUTED BY PROGRAM
 BW - - - WIDTH (DEL X) OF FIRST LEVEL CELLS (M OR FT)
 BH - - - HEIGHT (DEL R) OF FIRST LEVEL CELLS (M OR FT)
 NW - - - NUMBER OF FIRST LEVEL CELLS IN X DIRECTION
 NH - - - NUMBER OF FIRST LEVEL CELLS IN RADIAL DIRECTION
 NL 1 NUMBER OF LEVELS OF CELLS
 NPA 0 NUMBER OF FIRST LEVEL CELLS AHEAD OF LEVEL 2
 NCA 0 NUMBER OF FIRST LEVEL CELLS SUBDIVIDED INTO SECOND
 LEVEL CELLS ALONG THE X DIRECTION
 NHA 0 AS ABOVE BUT IN RADIAL DIRECTION
 MW 0 NUMBER OF SECOND LEVEL CELLS IN THE X DIRECTION
 MH 0 NUMBER OF SECOND LEVEL CELLS IN THE RADIAL DIRECTION
 NFB 0 NUMBER OF SECOND LEVEL CELLS AHEAD OF LEVEL 3
 NCB 0 NUMBER OF SECOND LEVEL CELLS SUBDIVIDED INTO THIRD
 LEVEL CELLS ALONG THE X DIRECTION
 NHB 0 AS ABOVE BUT IN RADIAL DIRECTION
 LW 0 NUMBER OF THIRD LEVEL CELLS IN X DIRECTION
 LH 0 NUMBER OF THIRD LEVEL CELLS IN RADIAL DIRECTION
 LD - - - NUMBER OF FIRST LEVEL CELLS FROM AXIS IN R DIRECTION
 FOR WEIGHTING FACTOR BOUNDARIES (5 INTEGERS)
 LF - - - WEIGHTING FACTOR RATIOS AT BOUNDARIES ABOVE (5 INTEGERS)

&COUPLE - ONE OCCURRENCE (NEW RUN ONLY) - DISTRIBUTION FUNCTION
 PARAMETER DEFAULT DEFINITION
 NS - - - THE NUMBER OF BODY SEGMENTS FOR ACCUMULATING VELOCITY
 DISTRIBUTION FUNCTION INFORMATION

ORIGINAL PAGE 1
OF POOR QUALITY PRINCETON UNIVERSITY TIME-SHARING SYSTEM

C	MS	- - -	ARRAY(NS) OF AXIAL SEGMENT NUMBERS
C	IWS	- - -	ARRAY(NS) OF AZIMUTHAL WEDGE NUMBERS
C	VEL	3., 3., 4.	THERMAL VELOCITY SPREAD FOR THE UNCOLLIDED MOLECULES
C	MJ	20	NUMBER OF DEL V REGIONS FOR SAMPLING VELOCITY SPACE
C	SL	- - -	ARRAY GIVING LOWER BOUND ON THE VELOCITY SAMPLE OF COLLIDED MOLECULES (MSP X NS X 3)
C	DELS	- - -	THE RANGE (SL < V < SL+DELS) FOR SAMPLE OF COLLIDED MOLECULES (MSP X NS X 3)

IF DISTRIBUTION FUNCTION INFORMATION IS NOT DESIRED USE:
&COUPLE NS=0 &END

A SAMPLE INPUT DECK IS GIVEN BELOW:

```

&CONTRL NAME='SHUT', 'TLE ', TITLE='HYPE', 'RBOL', 'A AT', ' 95K', 'M  M', 'ON. ',
  DEBUG=.F.,, T.,, T.,, NEW=.T.,, SAVE=.T.,, ICOPY=0, REDO=.T. &END
&TIMES DTM=.025, ITS=5, ITP=1000, TST=400, TLIM=1000 &END
&FLOREF INM=4500, MNM=20000, MNB=100, MSP=1, MET=0, U=7485.9, ANGLE=0.0, RNU=1., 2*0.,
  RMA=28.94, 0., 0., TF=195.51, DENF=2.52E+19 &END
&MOLEC TRF=1000, DIR=3.5E-19, ETA=.104, PHI=0.0, CHI=-1., ACR=.001 &END
&SHAPES BODY=1.00 &END
&SHAPES BODY=.0173, 1590., 3*0.0, 2*1.0 &END
&SHAPES BODY=.0672, 1590., 3*0.0, 2*1.0 &END
&SHAPES BODY=.1444, 1590., 3*0.0, 2*1.0 &END
&SHAPES BODY=.2432, 1590., 3*0.0, 2*1.0 &END
&SHAPES BODY=.3579, 1590., 3*0.0, 2*1.0 &END
&SHAPES BODY=.4842, 1590., 3*0.0, 2*1.0 &END
&SHAPES BODY=.6192, 1590., 3*0.0, 2*1.0 &END
&SHAPES BODY=.7405, 1590., 3*0.0, 2*1.0 &END
&SHAPES BODY=.7500, 1590., 0., 1., 0., 2*1., 0., -1.423278, -2.286, 0. &END
&SHAPES BODY=.9000, 1590., 0.0, 4*1.0, 0.9000, -111.82, 2*0.0 &END
&GEOM NWEDGE=1, 0, THETAZ=180., RMB=0.0, BW=.05, BH=.1, NW=40, NH=24, NL=2,
  NPA=15, NCA=20, NHA=16, MW=40, MH=32, LD=1, 2, 4, 8, 12, LF=3, 2, 2, 2, 1 &END
&COUPLE NS=3, MS=1, 5, 8, IWS=1, 1, 1, VEL=2.5, 2.5, 2.0, MJ=20, SL=9*1., 18*-9.,
  DELS=27*20. &END

```

```

CALL RUN(NSP,NMB,NMC,NMP,NPB,DBA,NB,NBF,NBT,TMP,TMPA,XV,XVA,YV,
1YVA,ZV,ZVA,T,DB,FNB,XC,YC,ZC,NUMCEL,LM,LPP,PAU,PAV,PAW,PAX,PAY,
2PAZ,LCOL,TRP,TRPA,ER,LKW,NBS,LB,NBM,NBN)
WRITE(6,2)
STOP
END

```

```

SUBROUTINE RUN(NSP,NMB,NMC,NMP,NPB,DBA,NB,NBF,NBT,TMP,TMPA,XV,XVA,
1YV,YVA,ZV,ZVA,T,DB,FNB,XC,YC,ZC,NUMCEL,LM,LPP,PAU,PAV,PAW,PAX,
2PAY,PAZ,LCOL,TRP,TRPA,ER,LKW,NBS,LB,NBM,NBN)
MAIN RUNNING PROGRAM ** RUN *** CALLS ALL OTHER SUBROUTINES
*****
```

```

INTEGER*2 LM,LPP,LCOL,LKW
INTEGER*2 LB,NBM,NBN,NB,NBF,NBT,NBS,NUMCEL
INTEGER PRT,SAMP,TST,TLIM,TIME,Q
LOGICAL DUMP,DEBUG(3),SAVE,NEW,REDO

```


FORMATS

ORIGINAL PRINTOUT
OF POOR QUALITY

RUN0690
RUN0700
RUN0710
RUN0720
RUN0730
RUN0740
RUN0750
RUN0760
RUN0770

1 FORMAT(1H1)
2 FORMAT(1H1/17X,'RARIFIED SUPERSONIC FLOW OF BINARY GAS',T74,'I')
3 FORMAT('+',103X,'COPY ',I2)
4 FORMAT(/17X,'FLOW THROUGH ALL THE BOUNDARIES'//)
5 FORMAT(3X,3I4,6E18.6/)
6 FORMAT(7X,2I4, E18.6,72X,E18.6)
30 FORMAT(' TIME = ',F6.3,60X,'RANDOM NUMBER GENERATOR HAS BEEN CALLED
1 ',I10,' TIMES')
31 FORMAT(' CPU TIME LEFT- ',F8.3)
32 FORMAT(7X,'-MOLECULES-/3X,3I6)
33 FORMAT(' TIME = ',F8.3,5X,'COLLISION LOOP=',F8.3,5X,'MOVE LOOP = '
1,F8.3,5X,'TOTAL TIME = ',F8.3/21X,'2ND MOVE LOOP = ',F8.3,5X,
2'CLEANUP LOOP=',F8.3,4X,'PARTICLE NUMBERS = ',4I6)
34 FORMAT(9X,'-MOLECULAR COLLISIONS-/3(3I14/))
35 FORMAT(2X,'-COLLISIONS WITH SURFACE-/3X,3I8)
36 FORMAT(' MAXIMUM NUMBER OF MOLECULES SO FAR- ',I6//) RUN0880
38 FORMAT(' EXCESS MOLECULES OCCURRED IN ',3A4) RUN0890
40 FORMAT(/' SOMETHING IS WRONG WITH BOX NUMBERING IN RUN ',/5I5,2E17, RUN0900
17,3I5,2E17.7,5I5/2E17.7) RUN0910
42 FORMAT(/' SOMETHING WRONG IN COMPUTING XSTART'/1X,8E16.8) RUN0920
44 FORMAT(' NB(',I2,',',I4,') POPULATION EXCEEDED ',I3,' IN MAIN AT TRUN0930
TIME = ',F7.3) RUN0940
50 FORMAT(///' SNAP SAVED ON TAPE') RUN0950
RUN0960
RUN0970
RUN0980

CPA=ELTIME(0)
CALL NOUNDF
LIMIT(4)=NMC
LIMIT(5)=NMP
LIMIT(6)=NPB
LIMIT(7)=NMB
LIMIT(10)=NSP
KAWLS=0
PI=3.141593
PIROOT=SQRT(PI) RUN1030
MET=0
LARGE=0 RUN1060
NL=1 RUN1080
NFA=0 RUN1090
NCA=0 RUN1100
NHA=0 RUN1110
MW=0 RUN1120
MH=0 RUN1130
NFB=0 RUN1140
NCB=0 RUN1150
NHB=0 RUN1160
LW=0 RUN1170
LH=0 RUN1180

MJ=20
 DUMP=.TRUE.
 DEBUG(1)=.FALSE.
 DEBUG(2)=.FALSE.
 DEBUG(3)=.TRUE.
 SAVE=.FALSE.
 NEW=.TRUE.
 REDO=.FALSE.
 PERCNT=.001
 ACR=.001
 DO 58 I=1,15
 58 BODY(I)=0.0
 DO 60 I=1,3
 RNU(I)=0.0
 RMA(I)=0.0
 CHI(I)=0.0
 DO 59 J=1,18
 ALPHA(I,J)=1.0
 59 SIGMA(I,J)=1.0
 DO 60 K=1,3
 ETA(I,K)=0.0
 PHI(I,K)=0.0
 60 DIR(I,K)=0.0
 VEL(1)=3.
 VEL(2)=3.
 VEL(3)=4.
 WRITE(6,1)
 READ(5,CTRL)
 WRITE(6,CTRL)
 IF(NEW) GO TO 103
 REWIND 9
 READ(9) DENP,U,XREF,TRF, KAWLS,NL,NW,NH,MW,MH,LW,LH,NXA,NXB,NCA RUN1380
 1 ,NCB,NFA,NFB,NHA,NHB,BW,BH,BWB,BHB,BWC,BHC,XLB,XLC,PI,NREGRUN1390
 2 ,S,SINANG,COSANG,AKN,NBX,RM,XR,ND,TIME,DTM,TI,ITS,ITP,TST RUN1400
 3 ,TLIM,RMA,RNU,DIR,XSTART,JNM,MNM,MNB,TR,BZC,CN7,DRF,FCF RUN1410
 4 ,FNA,HTF,INM,ITYPE,JTYPE,MJ,NAV,NMAX,NS,NWEDG,PRT,SAMP RUN1420
 5 ,BTA,C1,C2,C3,C7,C8,DAM,DPA,FL,DELANG,FDN,HTI,HTR,JNT,KNM RUN1430
 6 ,NM,WTM,C4,VRM,NCOL,LD,LF,LWF,RLD,CTI,CTR,CNI,CNR,LEV,SN RUN1440
 7 ,ST,D1,D2,D3,D4,SSA,SSB,MS,NSP,NMB,NMC,NMP,NPB,NRAN,VELR
 8 ,IWS,TANGN,XLIM,COEFF,XCB,XS,YCB,TB,ALPHA,SIGMA,NTS,NTSF RUN1460
 9 ,UTL,UTT,VTS,HTS,HTSI,ENT,REM,ENTS,REMS,PTH,THETA,DTH,TMPARUN1470
 A ,DBA,NB,NBF,NBT,TMP,XV,XVA,YV,YVA,ZV,ZVA,T,DE,FNB,XC,YC,ZCRUN1480
 B ,NUMCEL,PAU,PAV,PAW,PAX,PAY,PAZ,FV,NTCV,NTCP,LPF,LCOL,LM RUN1490
 C,ETA,PHI,CHI,CN,CM,CNG,CMG,CN8,TRP,TRPA,THETAZ,NWEDGE,MSP,ANGLE,TF
 D,UTLI,UTTI,VTSI,ER,RMB,LKW,NBS,LB,NBM,NBN
 REWIND 9
 DTMO=DTM
 READ(5,TIMES)
 WRITE(6,TIMES)
 IF(DTM.EQ.DTMO) GO TO 100
 AIME=TIME*DTMO
 TIME=AIME/DTM+0.1
 DO 99 J=1,NSP
 DO 99 L=1,NWEDG
 ENTS(J,L)=ENTS(J,L)*DTM/DTMO

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OF POOR QUALITY

RUN1190
RUN1230
RUN1260
RUN1240
RUN1250

RUN1300
RUN1310
RUN1320
RUN1330
RUN1340

RUN1360
RUN1370

RUN1500

```

DO 98 K=1,6
DO 98 I=1,2
ENT(I,J,K,L)=ENT(I,J,K,L)*DTM/DTMO
98 CONTINUE
99 CONTINUE
100 IF(TI.GT.0.0) TST=TI/DTM
WRITE(6,2)
WRITE(6,4)
WRITE(6,5) (((I,J,L,(ENT(I,J,K,L),K=1,6),L=1,NWEDG),J=1,MSP),I=1,2)
WRITE(6,6) ((J,L,ENTS(J,L),PTH(J,L),L=1,NWEDG),J=1,MSP)
WRITE(6,2)
CALL PRINTA(THETAZ,NWEDGE,TITLE,NAME,XCB,YCB,TB,ALPHA,SIGMA,LD,LF,RUN1540
1XLIM,COEFF,LIMIT,MSP)
CALL PRINTB(PNA,MSP,PNB,LEV,LWF,NM,RLD,XLIM,XC,YC,ZC,NB,NUMCEL,LKW
1,MSP)
GO TO 280
103 READ(5,TIMES)
WRITE(6,TIMES)
READ(5,FLOREF)
WRITE(6,FLOREF)
READ(5,MOLEC)
WRITE(6,MOLEC)
IF(MSP.GT.LIMIT(10)) CALL DIAG(10,LIMIT(10),MSP)
CHIM=0.0
RMR=0.0
DMR=0.0
DO 115 M=1,MSP
RMR=RMR+RMA(M)*RNU(M)
CHIM=CHIM+CHI(M)*RNU(M)
DO 115 K=1,MSP
115 DMR=DMR+RNU(M)*RNU(K)*DIR(M,K)*(TRP/TF)**(ETA(M,K)/2.)
XREF=1.0/(DENF*DMR*1.414214)
VELR=SQRT(16628.64*TF/RMR)
IF(MET.NE.0) VELR=SQRT(99437.92*TF/RMR)
TMR=XREF/VELR
S=U/VELR
XLIM(1)=0.
TR=0.
YR=0.
NREG=0
ND=0
READ(5,SHAPES)
WRITE(6,SHAPES)
XO=BODY(1)/XREF
104 READ(5,SHAPES)
WRITE(6,SHAPES)
ND=ND+1
IF(ND.GT.LIMIT(3)) CALL DIAG(3,LIMIT(3),ND)
XCB(ND)=BODY(1)/XREF+XO
TB(ND)=BODY(2)/TF
YCB(ND)=BODY(3)
DO 1104 M=1,MSP
ALPHA(M,ND)=BODY(4+2*M)
1104 SIGMA(M,ND)=BODY(5+2*M)
IF(TB(ND).GT.TR) TR=TB(ND)

```

ORIGINAL PAGE IS
OF POOR QUALITY

RUN1550

RUN1590

RUN1610

RUN1620

RUN1630

RUN1640

RUN1650

RUN1670

RUN1680

RUN1690

RUN1770

```

IF (YCB(ND) .GT. YR) YR=YCB(ND) RUN1780
IF (BODY(4) .EQ. 0.0) GO TO 104 RUN1800
NREG=NREG+1
XLIM(NREG+2)=XCB(ND)
XT=BODY(6+2*MSP)/XREF+X0
A=BODY(7+2*MSP)
B=BODY(8+2*MSP)/XREF
C=BODY(9+2*MSP)/XREF**2
COEFF(1,NREG)=A
COEFF(2,NREG)=1.0
COEFF(3,NREG)=B-2.*A*XT
COEFF(4,NREG)=A*XT**2-B*XT+C
IF (BODY(5) .EQ. 0.0) GO TO 104
IF (NREG.GT.LIMIT(2)) CALL DIAG(2,LIMIT(2),NREG) RUN1860
NSTEP=NREG+2 RUN1870
A=COEFF(1,1) RUN1880
B=COEFF(3,1) RUN1890
C=COEFF(4,1) RUN1900
DISC=B*B-4.*A*C RUN1910
IF (DISC.LT.0.) DISC=0. RUN1920
DISC=SQRT(DISC) RUN1930
IF (A.NE.0.) GO TO 108 RUN1940
XSTART=-C/B RUN1950
GO TO 109 RUN1960
108 X1=.5/A*(-B+DISC) RUN1970
X2=.5/A*(-B-DISC) RUN1980
XSTART=AMIN1(X1,X2) RUN1990
XMAX=AMAX1(X1,X2)
IF (XMAX.LT.XLIM(3)) XSTART=XMAX
IF (XSTART.GT.0.) GO TO 109
WRITE(6,42) XSTART,A,COEFF(2,1),B,C,DISC,X1,X2
IF (DUMP) CALL ABEND(1)
STOP
109 XLIM(2)=XSTART RUN2000
A=COEFF(1,NREG) RUN2010
B=COEFF(3,NREG) RUN2020
C=COEFF(4,NREG) RUN2030
DISC=B*B-4.*A*C RUN2040
IF (DISC.LT.0.) DISC=0. RUN2050
DISC=SQRT(DISC) RUN2060
IF (A.NE.0.) GO TO 111 RUN2070
XLIM(NSTEP)=-C/B RUN2080
GO TO 112 RUN2090
111 X1=.5/A*(-B+DISC) RUN2100
X2=.5/A*(-B-DISC) RUN2110
XLIM(NSTEP)=AMAX1(X1,X2) RUN2120
XMIN=AMIN1(X1,X2) RUN2130
IF (XMIN.GT.XLIM(NSTEP-1)) XLIM(NSTEP)=XMIN RUN2140
IF (XLIM(NSTEP).GT.XLIM(NSTEP-1)) GO TO 112 RUN2150
WRITE(6,42) XLIM(NSTEP),A,COEFF(2,NREG),B,C,DISC,X1,X2
IF (DUMP) CALL ABEND(1)
STOP
112 AKN=1.0/(XLIM(NSTEP)-XSTART) RUN2170
XCB(ND)=XLIM(NSTEP) RUN2180
DO 260 N=1,3

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MS (N)=0
IWS (N)=0
DO 259 M=1,3
NCOL (N,M)=0
DO 258 K=1,3
SL (N,M,K)=0.0
258 DELS (N,M,K)=0.0
259 CONTINUE
260 CONTINUE
READ (5,GEOM)
READ (5,COUPLE)
WRITE (6,GEOM)
WRITE (6,COUPLE)
BW=BW/XREP
BH=BH/XREP
RMB=RMB/XREP
IF (RMB.GT.0.) GO TO 264
DO 262 K=1,NREG
XBEG=XLIM (K+1)
XEND=XLIM (K+2)
A=COEFF (1,K)
B=COEFF (2,K)
C=COEFF (3,K)
D=COEFF (4,K)
REND=SQRT (ABS ((A*XEND**2+C*XEND+D)/B))
XPEAK=0.
IF (A.NE.0.) XPEAK=-.5*C/A
IF ((XPEAK.LE.XBEG).OR.(XPEAK.GE.XEND)) GO TO 261
RTEMP=SQRT (ABS ((A*XPEAK**2+C*XPEAK+D)/B))
IF (RTEMP.GT.REND) REND=RTEMP
261 IF (REND.GT.RMB) RMB=REND
262 CONTINUE
264 CONTINUE
NWEDG=NWEDGE (1)+NWEDGE (2)
IF (NWEDGE (2).EQ.0) THETAZ=180.0
IF (SAVE) REWIND 9
IF (NWEDG.GT.LIMIT (1)) CALL DIAG (1,LIMIT (1),NWEDG)
IF (NNM.GT.LIMIT (5)) CALL DIAG (5,LIMIT (5),NNM)
IF (MNB.GT.LIMIT (6)) CALL DIAG (6,LIMIT (6),MNB)
IF (NS.GT.LIMIT (8)) CALL DIAG (8,LIMIT (8),NS)
IF ((NS.NE.0).AND.(MJ.GT.LIMIT (9))) CALL DIAG (9,LIMIT (9),MJ)
JNM=INM
DELANG (1)=THETAZ/NWEDGE (1)
DELANG (2)=0.0
IF (NWEDGE (2).NE.0) DELANG (2)=(180.-THETAZ)/NWEDGE (2)
SINANG=SIN (ANGLE/180.*PI)
COSANG=COS (ANGLE/180.*PI)
XR=BW*NW
XLIM (NSTEP+1)=XR
RM=BH*NH
VOL=PI*RM*RM*XR
NXA=NW*NH*NWEDG
NXB=MW*MH*NWEDGE (1)
NXC=LW*LH*NWEDGE (1)
IXA=NW*NH*NWEDGE (1)

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RUN2200

RUN2220

RUN2230

RUN2240

RUN2250

RUN2290

RUN2300

RUN2320

RUN2330

RUN2380

RUN2390

RUN2400

RUN2420

RUN2430

RUN2440

RUN2450

```

NBX=NXA+NXB+NXC
IF(NBX.GT.LIMIT(7)) CALL DIAG(7,LIMIT(7),R047)
BR=SQRT(TR)
DO 113 N=1,5
113 RLD(N)=BH*LD(N)
RLD(6)=RM
LWF(1)=1
RWFM=RLD(1)
B=RWFM*RWFM
C=B
DO 114 N=2,6
A=RLD(N)*RLD(N)
LWF(N)=LWF(N-1)*LF(N-1)
B=B+(A-C)/LWF(N)
D=RLD(N)/LWF(N)
IF(D.GT.RWFM) RWFM=D
114 C=A
INM=INM*RM*RM/B
DDN=INM/VOL
DO 140 MT=1,MSP
WTM(MT)=RMA(MT)/RMR
FDN(MT)=RNU(MT)*DDN
DPA(MT)=RNU(MT)
BTA(MT)=SQRT(WTM(MT))
SR(MT)=S*BTA(MT)
SN(MT)=SR(MT)*COSANG
ST(MT)=SR(MT)*SINANG
DO 117 K=1,MSP
DAM(K,MT)=DIR(K,MT)*(TR/TP)**(ETA(K,MT)/2.)/DMR
CH8(K,MT)=DDN/DAM(K,MT)*1.414214
BT=A MIN1(BTA(K),BTA(MT))
VR1=S+3.* (1.+SQRT(TR))/BT
VR2=3.*SQRT((1.+2.*S**2/(5.+CHIM))* (1./WTM(K)+1./WTM(MT)))
CM(K,MT,1)=AMAX1(VR1,VR2)
CN(K,MT,1)=RAND(0)*CM(K,MT,1)
DF=PHI(K,MT)*(CHI(K)+CHI(MT)+2.)-1
DS=PHI(K,MT)*(2.-.5*ETA(K,MT))-1.0
DO 917 N=2,3
XPM=ACR**AMIN1(DF,DS)
IF((DF.GT.0.).AND.(DS.GT.0.)) XPM=(DF/(DF+DS))**DF*(DS/(DF+DS))**DS
XPN=ACR**AMAX1(DF,DS)
IF((DF.LT.0.).AND.(DS.LT.0.)) XPM=(DF/(DF+DS))**DF*(DS/(DF+DS))**DS
CM(K,MT,N)=XPM-XPN
CN(K,MT,N)=RAND(0)*CM(K,MT,N)
DF=CHI(K)
DS=CHI(MT)
917 CONTINUE
117 CONTINUE
ARG=SN(MT)
DO 119 NT=1,2
D=ERRF(ARG)
TEMPA=EXP(-SN(MT)*SN(MT))/3.544908+0.5*ARG*D
TEMPA=TEMPA/BTA(MT)
TEMPC=0
DO 118 N=1,6

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RUN2460
RUN2470
RUN2480
RUN2490
RUN2500
RUN2510
RUN2520
RUN2530
RUN2540
RUN2550
RUN2560
RUN2570
RUN2580
RUN2590
RUN2600
RUN2610
RUN2620
RUN2630

RUN2730
RUN2740
RUN2750

RUN2760
RUN2770
RUN2780
RUN2790
RUN2800
RUN2810
RUN2820

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TEMPB=RLD (N) *RLD (N) RUN2830
TEMPD=INM*TEMPA*DTM/(XR*RM*RM)*(TEMPB-TEMPC)/LWF (N) RUN2840
DO 116 K=1,NWEDG RUN2850
DELTH=DELANG (1) RUN2860
IF (K.GT.NWEDGE (1)) DELTH=DELANG (2) RUN2870
ENT(NT,MT,N,K)=TEMPD*DELTH/180. RUN2880
116 REM(NT,MT,N,K)=0. RUN2890
TEMPB=TEMPB
118 CONTINUE
SSR=ARG*ARG RUN2900
SSA(NT,MT)=ARG+SQRT(SSR+2.) RUN2910
SSB(NT,MT)=SSA(NT,MT)*( .25*SSA(NT,MT)-ARG) RUN2920
119 ARG=-ARG RUN2930
CHT=CHI (MT) RUN2940
IF (CHT.GT.0.) CMG (MT)=CHT**CHT*EXP (-CHT) RUN2950
IF (CHT.EQ.0.) CMG (MT)=1.0
IF (CHT.LT.0.) CMG (MT)=ACR**CHT*EXP (-ACR)
CNG (MT)=RAND (0)*CMG (MT)
ANG=0.
N=0
ARG=ST (MT)
TEMPD=INM/PI*DTM/RM*2./(BTA (MT)*LWF (6))
DO 135 M=1,2
I=NWEDGE (M)
IF (I.EQ.0) GO TO 135
DO 134 K=1,I
N=N+1
AA=ANG
BB=DELANG (M)
CC=AA+BB
ANG=ANG+BB
THETA (N)=AA
DTH (N)=BB
AA=AA*PI/180.
CC=CC*PI/180.
ARGA=ARG*COS (AA)
ARGC=ARG*COS (CC)
ENTS (MT,N)=0.
REMS (MT,N)=0.
PTH (MT,N)=0.
TEMPA=0.
IF (ARGA.LE.-10.) GO TO 120
TEMPA=FNCTM (ARGA,PIROOT,L,COEFF)
PTH (MT,N)=2.*PIROOT*TEMPA
120 TEMPB=0.
IF (ARGC.LE.-10.) GO TO 125
TEMPB=FNCTM (ARGC,PIROOT,L,COEFF)
125 SUM1=TEMPA+TEMPB
TEMPC=.5*SUM1*(CC-AA)*TEMPD
IF (TEMPC.LT.1.E-06) GO TO 134
CALL SIMPSN (ARGA,ARGC,0,INTGRL,PERCNT,COEFF,PIROOT,SUM1,FNCTM)
ENTS (MT,N)=INTGRL*TEMPD*(CC-AA)
134 CONTINUE
135 CONTINUE
140 CONTINUE
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XS (1) = .5 * (XCB (1) - XSTART) * AKN RUN3370
 DO 155 N=2, ND RUN3380
 155 XS (N) = (.5 * (XCB (N) + XCB (N-1)) - XSTART) * AKN RUN3390
 IP (NS.EQ.0) GO TO 160 RUN3400
 DO 159 I=1, NS RUN3410
 N=MS (I) RUN3420
 X=XS (N) / AKN + XSTART RUN3430
 J=0 RUN3440
 157 J=J+1 RUN3450
 IF (X.GT. XLIM (J+2)) GO TO 157 RUN3460
 CALL HEIGHT (X, Y, J, COEFF, 3) RUN3470
 TANGN (I) = (IWS (I) -.5) * DELANG (1) RUN3480
 IF (IWS (I) * DELANG (1).GT. THETAZ) TANGN (I) = THETAZ + (IWS (I) - NWEDGE (1) -.15) * DELANG (2) RUN3490
 Z=Y*SIN (TANGN (I) * PI/180.) RUN3500
 Y=-Y*COS (TANGN (I) * PI/180.) RUN3510
 CALL NORMAL (EYE, JAY, KAY, ONE, COEFF) RUN3520
 SNN=-S*(COSANG*EYE+SINANG*JAY) RUN3530
 ST1=S*(COSANG*ONE-EYE*SINANG*JAY/ONE) RUN3540
 ST2=-S*SINANG*KAY/ONE RUN3550
 DO 159 MT=1, MSP RUN3560
 VL (MT, 1) = AMAX1 (0., SNN-VEL (1) / BTA (MT)) RUN3580
 VL (MT, 2) = ST1-VEL (2) / BTA (MT) RUN3590
 VL (MT, 3) = ST2-VEL (3) / BTA (MT) RUN3600
 DELV (MT, 1) = SNN+VEL (1) / BTA (MT) - VL (MT, 1) RUN3610
 DELV (MT, 2) = 2.*VEL (2) / BTA (MT) RUN3620
 DELV (MT, 3) = 2.*VEL (3) / BTA (MT) RUN3630
 AMJ=MJ-1 RUN3640
 DO 159 K=1, 3 RUN3650
 DO 159 J=1, MJ RUN3660
 FV (MT, I, 1, J, K) = VL (MT, K) + (J-1) / AMJ * DELV (MT, K) RUN3670
 FV (MT, I, 2, J, K) = SL (MT, I, K) + (J-1) / AMJ * DELS (MT, I, K) RUN3680
 159 CONTINUE RUN3690
 160 CONTINUE RUN3700
 IF (YR.GT.0.) GO TO 169 RUN3710
 A=XSTART RUN3720
 B=XCB (1) RUN3730
 L=1 RUN3740
 SUM1=FNCTN (A, PIROOT, L, COEFF) + FNCTN (B, PIROOT, L, COEFF) RUN3750
 CALL SIMPSN (A, B, L, INTGRL, PERCENT, COEFF, PIROOT, SUM1, FNCTN) RUN3760
 INTGRL=INTGRL/(ABS (COEFF (2, L)) * RMB*RMB) * (B-A) RUN3770
 YCB (1)=INTGRL RUN3780
 DO 168 N=2, ND RUN3790
 A=XCB (N-1) RUN3800
 B=XCB (N) RUN3810
 IF (XLIM (L+2).GE.B) GO TO 167 RUN3820
 L=L+1 RUN3830
 167 SUM1=FNCTN (A, PIROOT, L, COEFF) + FNCTN (B, PIROOT, L, COEFF) RUN3840
 CALL SIMPSN (A, B, L, INTGRL, PERCENT, COEFF, PIROOT, SUM1, FNCTN) RUN3850
 INTGRL=INTGRL/(ABS (COEFF (2, L)) * RMB*RMB) * (B-A) RUN3860
 168 YCB (N)=INTGRL RUN3870
 169 LV=2 RUN3880
 IF (NWEDGE (2).EQ.0) LV=1 RUN3890
 CALL CELL (THETAZ, LV, BW, BH, NW, NH, 0., 0, 0, 0, DELANG, NWEDGE, XC, YC, ZC, FNB)
 IF (NL.LT.2) GO TO 170

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CALL ZERO (NW, NI, NFA, NCA, NHA, 0, NWEDGE (1), FNB)	RUN3900
XLB=BW*NFA	RUN3910
BWB=BW*NCA/MW	RUN3920
BHB=BH*NHA/MH	RUN3930
CALL CELL (THETAZ, 1, BWB, BHB, MW, MH, XLB, NXA, 0, DELANG, NWEDGE, XC, YC, ZC, RUN3940 1 FNB)	RUN3950
IF (NL.LT. 3) GO TO 170	RUN3960
CALL ZERO (MW, MH, NFB, NCB, NHB, NXA, NWEDGE (1), FNB)	RUN3970
XLC=XLB+BWB*NFB	RUN3980
BWC=BWB*NCB/LW	RUN3990
BHC=BHB*NHB/LH	RUN4000
CALL CELL (THETAZ, 1, BWC, BHC, LW, LH, XLC, NXA, NXB, DELANG, NWEDGE, XC, YC, ZRUN4010 1C, FNB)	RUN4020
170 CALL SBTRCT (1, IXA, NXA, BW, BH, DELANG, XC, YC, FNB, XLIM, COEFF)	RUN4030
LEV (1)=NBX+1	RUN4040
LEV (2)=NBX+1	RUN4050
IF (NL.LT. 2) GO TO 190	RUN4060
NI=NXA+1	RUN4070
NF=NXA+NXB	RUN4080
LEV (1)=NI	RUN4090
CALL SBTRCT (NI, NP, NF, BWB, BHB, DELANG, XC, YC, FNB, XLIM, COEFF)	RUN4100
IF (NL.LT. 3) GO TO 190	RUN4110
NI=NP+1	RUN4120
NF=NBX	RUN4130
LEV (2)=NI	RUN4140
CALL SBTRCT (NI, NF, NP, BWC, BHC, DELANG, XC, YC, FNB, XLIM, COEFF)	RUN4150
190 FNA=0.0	RUN4160
M=0	RUN4170
DO 210 N=1, NBX	RUN4180
NUMCEL (N)=0	RUN4190
IF (FNB (N). LE. 0.) GO TO 210	RUN4200
M=M+1	RUN4210
NUMCEL (N)=M	ORIGINAL PAGE IS OF POOR QUALITY
FNA=FNA+FNB (N)	
DYC=BH/2.	
IF (N.GT. NXA) DYC=BHB/2.	
IF (N.GT. NXA+NXB) DYC=BHC/2.	
YTC=YC (N) +DYC	
DO 200 LA=1, 6	
IF (YTC.LE. RLD (LA)) GO TO 201	
200 CONTINUE	
201 LKW (M)=LWF (LA)	RUN4220
210 CONTINUE	RUN4230
NPX=M	RUN4240
IF (NPX.GT. LIMIT (4)) CALL DIAG (4, LIMIT (4), NPX)	RUN4250
220 TIME=0	RUN4260
LARGE=0	RUN4270
SAMP=0	RUN4290
PRT=0	RUN4300
NAV=0	RUN4310
AIME=0.	RUN4380
TI=-1.	RUN4410
NMAX=0	
DO 250 MT=1, 3	
C1 (MT)=RAND (0)	

C2(MT)=RAND(0)	RUN4420
C3(MT)=RAND(0)	RUN4430
C7(MT)=RAND(0)	RUN4440
C8(MT)=RAND(0)	RUN4450
D1(MT)=RAND(0)	RUN4460
D2(MT)=RAND(0)	RUN4470
D3(MT)=RAND(0)	RUN4480
D4(MT)=RAND(0)	RUN4490
PL(MT)=0.	RUN4500
HTI(MT)=0.	RUN4510
HTR(MT)=0.	RUN4520
JNT(MT)=0	RUN4530
NM(MT)=0	RUN4540
DO 230 N=1,3	RUN4550
CTI(MT,N)=0.	RUN4560
CTR(MT,N)=0.	RUN4570
CNI(MT,N)=0.	RUN4580
230 CNR(MT,N)=0.	RUN4590
DO 240 N=1,ND	RUN4600
DO 240 K=1,NWEDG	RUN4610
NTS(MT,N,K)=0	RUN4620
NTSF(MT,N,K)=0	RUN4630
HTSI(MT,N,K)=0.	RUN4640
UTLI(MT,N,K)=0.	RUN4650
UTTI(MT,N,K)=0.	RUN4660
VTSI(MT,N,K)=0.	RUN4670
UTL(MT,N,K)=0.	RUN4680
UTT(MT,N,K)=0.	RUN4820
VTS(MT,N,K)=0.	RUN4830
240 HTS(MT,N,K)=0.	RUN4840
DO 250 I=1,3	RUN4850
NTCP(MT,I)=0	RUN4860
DO 250 L=1,2	RUN4870
DO 250 K=1,3	RUN4880
DO 250 J=1,MJ	RUN4890
NTCV(MT,I,L,J,K)=0	RUN4700
250 CONTINUE	RUN4710
DO 245 N=1,NPX	RUN4720
DO 245 MT=1,NSP	RUN4730
NB(MT,N)=0	RUN4740
NBF(MT,N)=0	RUN4750
NBS(MT,N)=0	RUN4760
NBT(MT,N)=0	RUN4770
DBA(MT,N)=0.	RUN4780
XVA(MT,N)=0.	RUN4790
YVA(MT,N)=0.	RUN4800
ZVA(MT,N)=0.	RUN4930
TMPA(MT,N)=0.	RUN4940
TRPA(MT,N)=0.0	
DO 245 NN=1,NSP	
T(MT,NN,N)=0.0	
FND=DDN	
245 CONTINUE	
DRF=2.0/(FND*S*S*RMB*RMB*PI)	
FCF=1.0/(FND*S*RMB*RMB*PI)	

HTP=.5*DRP/S
 C9=RAND(0)*RWPM
 LL=PNA/VOL*INM
 WRITE(6,2)
 WRITE(6,4)
 WRITE(6,5) (((I,J,L,(ENT(I,J,K,L),K=1,6),L=1,NWEDG),J=1,MSP),I=1,2)
 WRITE(6,6) ((J,L,ENTS(J,L),PTH(J,L),L=1,NWEDG),J=1,MSP)
 WRITE(6,2) RUN5000
 CALL PRINTA(THETAZ,NWEDGE,TITLE,NAME,XCB,YCB,TB,ALPHA,SIGMA,LD,LF,RUN5010
 1XLIM,COEFF,LIMIT,MSP)
 CALL GAS(NWEDG,THETAZ,DELANG,NWEDGE,BTA,C1,DFA,NM,RLD,LWP,FNB,DB,NRUN5030
 1B,NBF,LPF,PAU,PAV,PAW,PAY,PAZ,XLIM,COEFF,LH,LIMIT(4),LIMIT(6),RUN5040
 2LARGE,MNM,MNB,DEBUG(1),LCOL,NUMCEL,MSP,ER,CHI,CNG,CMG,NSP,LB,NBM,
 3NBN)
 CPUTYM=TPIND(0)
 IF(LARGE.NE.0) GO TO 345 RUN5060
 DO 265 I=1,MSP
 265 IF(NM(I).GT.NMAX) NMAX=NM(I)
 CALL PRINTE(FNA,MSP,FNB,LEV,LWP,NM,RLD,XLIM,XC,YC,ZC,NB,NUMCEL,LKW
 1,NSP)
 IF(DEBUG(2)) WRITE(6,1)
 CALL ACCUM(NMC,NPB,FNB,NB,PAU,PAV,PAW,ER,TMP,TRP,XV,YV,ZV,LN,MSP,
 1NSP,LPP,NBP,NBM)
 CPA=ELTIME(0)
 CPI=CPA
 GO TO 340 RUN5130
 280 TIME=TIME+1 RUN5140
 285 LARGE=0
 CPI=ELTIME(0)
 AIME=TIME*DTM
 IF(DEBUG(1)) WRITE(6,33) AIME,CPC,CPM,CPI,CPB,CPA,(NM(I),I=1,3),NMAX RUN5170
 PRT=PRT+1 RUN5180
 SAMP=SAMP+1 RUN5190
 CALL COLIDE(CN,CM,WTM,DB,DBA,NB,NCOL,PAU,PAV,PAW,ER,T,LN,MSP,
 1LIMIT(8),LIMIT(6),NUMCEL,ETA,PHI,CHI,CN8,NSP,LPP,LKW,NBP,NBM)
 KNN(1)=0 RUN5220
 KNN(2)=0 RUN5230
 KNN(3)=0
 CPC=ELTIME(0)
 IF(DEBUG(1)) WRITE(6,33) AIME,CPC,CPM,CPI,CPB,CPA,(NM(I),I=1,3),NMAX
 CALL MOVE(0,AKN,MJ,NS,NWEDG,THETAZ,XSTART,LIMIT(3),LIMIT(1),LIMIT RUN5240
 1(8),LIMIT(9),DELANG,NWEDGE,BTA,C2,C3,DFA,FL,HTI,HTR,JNT,KNM,NM,XCB RUN5250
 2,XLIM,MS,IWS,NTCF,NTCV,FV,CTI,CTR,CNI,CNR,ALPHA,SIGMA,COEFF,HTS,HTR RUN5260
 3SI,NTS,NTSP,UTL,UTT,VTS,PAU,PAV,PAW,PAY,PAZ,LPP,LCOL,TB,MSP,ER
 4,CHI,CNG,CMG,NSP,UTLI,UTTI,VTSI)
 KNM(1)=NM(1)
 KNM(2)=NM(2)
 KNM(3)=NM(3)
 CPM=ELTIME(0)
 IF(DEBUG(1)) WRITE(6,33) AIME,CPC,CPM,CPI,CPB,CPA,(NM(I),I=1,3),NMAX
 CALL FLOW(NWEDG,MNM,LARGE,BTA,C1,C7,C8,D1,D2,D3,D4,DTH,NM,SN,ST,THRUN5300
 1ETA,LWP,RLD,PTH,ENTS,REMS,SSA,SSB,PAU,PAV,PAW,PAY,PAZ,LPP,ENT,RUN5310
 2REM,LCOL,MSP,ER,CHI,CNG,CMG,NSP)
 IF(LARGE.NE.0) GO TO 345 RUN5330
 CPB=ELTIME(0)

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IF (DEBUG (1)) WRITE (6,33) AIME,CPC,CPM,CPI,CPB,CPA, (NM(I), I=1,3), NMAX
CALL MOVE (1,AKN,MJ,NS,NWEDG,THETAZ,XSTART,LIMIT(3),LIMIT(1),LIMIT
1(8),LIMIT(9),DELANG,NWEDGE,BTA,C2,C3,DFA,FL,HTI,HTR,JNT,KNM,NM,XCBRUN5350
2,XLIM,MS,IWS,NTCF,NTCV,FV,CTI,CTR,CNI,CNR,ALPHA,SIGMA,COEFF,HTS,HTRUN5360
3SI,NTS,NTSF,UTL,UTT,VTS,PAU,PAV,PAW,PAX,PAY,PAZ,LPP,LCOL,TB,MSP,ER
4,CHI,CMG,CMG,NSP,UTLI,UTTI,VTSI)
DO 330 MT=1,MSP
CPB=CPB+ELTIME (0)
IF (DEBUG (1)) WRITE (6,33) AIME,CPC,CPM,CPI,CPB,CPA, (NM(I), I=1,3), NMAX
M=0
DO 290 N=1,NBX
IF (NUMCEL(N).EQ.0) GO TO 290
M=M+1
NB(MT,M)=0
NBF(MT,M)=0
290 CONTINUE
NG=NM(MT)
N=0
295 N=N+1
IF (N.GT.NG) GO TO 310
X=PAX(MT,N)
Y=PAY(MT,N)
Z=PAZ(MT,N)
R=SQRT (Y*Y+Z*Z);
ARG=Y/R
TANG=180.* (1.-ARCCOS (ARG)/PI)
IWDGE=TANG/DELANG (1)
IF ((IWDGE.GE.NWEDGE(1)).AND.(DELANG(2).NE.0.)) IWDGE=(TANG-THETAZ)
1/DELANG(2)+NWEDGE(1)
IF (IWDGE.LT.0) IWDGE=0
IF (IWDGE.GE.NWEDG) IWDGE=NWEDG-1
L=X/BW+1.
IF (L.GT.NW) L=NW
M=R/BH
IF (M.GE.NH) M=NH-1
K=(IWDGE*NH+M)*NW+L
IF (K.LE.NXA) GO TO 296
WRITE (6,40)L,M,K,MT,N,X,Y,NWEDGE,NWEDG,TANG,Z,IWDGE,NH,NW
IF (DUMP) CALL ABEND(4)
STOP
296 KW=0
IF (NL.EQ.1) GO TO 300
IF (IWDGE.GE.NWEDGE(1)) GO TO 300
IF (PNB(K).GT.0.) GO TO 300
L=(X-XLB)/BWB+1.
IF (L.GT.MW) L=MW
M=P*BHB
IF (M.GE.MH) M=MH-1
K=(IWDGE*MH+M)*MW+L+NXA
IF (K.LE.NXA+NXB) GO TO 297
WRITE (6,40)L,M,K,MT,N,X,Y,NWEDGE,NWEDG,TANG,Z,IWDGE,MH,MW,NXA
IF (DUMP) CALL ABEND(5)
STOP
297 IF (NL.EQ.2) GO TO 300
IF (PNB(K).GT.0.) GO TO 300

```

ORIGINAL PRINT
OF POOR QUALITY

RUN5410
RUN5420
RUN5440
RUN5450
RUN5460
RUN5470
RUN5480
RUN5490
RUN5500
RUN5510
RUN5520
RUN5530
RUN5540
RUN5550
RUN5560
RUN5570
RUN5580
RUN5600
RUN5610
RUN5620
RUN5630
RUN5640
RUN5650
RUN5660
RUN5670
RUN5680
RUN5690
RUN5700
RUN5710
RUN5730
RUN5740
RUN5750
RUN5760
RUN5770
RUN5780
RUN5790
RUN5800
RUN5810
RUN5820
RUN5830
RUN5840

L=(X-XLC)/BWC+1.
 IF(L.GT.LW) L=LW
 M=R/BHC
 IF(M.GE.LH) M=LH-1
 K=(IWDGE*LH+M)*LW+L+NXA+NXB
 IF(K.LE.NBX) GO TO 300
 WRITE(6,40)L,M,K,MT,N,X,Y,NWEDGE,NWEDG,TANG,Z,IWDGE,LH,LW,NXA,NXB
 IF(DUMP) CALL ABEND(6)
 STOP
 300 LA=0
 IF(NUMCEL(K).EQ.0) GO TO 306
 301 LA=LA+1
 IF(R.GT.RLD(LA)) GO TO 301
 KW=LWP(LA)
 KKW=LPP(MT,N)
 IF(KW.EQ.KKW) GO TO 305
 IF(KW.LT.KKW) GO TO 302
 M=KW/KKW
 A=RAND(0)
 B=M
 B=1./B
 IF(A.LT.B) GO TO 305
 306 PAX(MT,N)=PAX(MT,NG)
 PAY(MT,N)=PAY(MT,NG)
 PAZ(MT,N)=PAZ(MT,NG)
 PAU(MT,N)=PAU(MT,NG)
 PAV(MT,N)=PAV(MT,NG)
 PAW(MT,N)=PAW(MT,NG)
 ER(MT,N)=ER(MT,NG)
 LPP(MT,N)=LPP(MT,NG)
 LCOL(MT,N)=LCOL(MT,NG)
 N=N-1
 NM(MT)=NM(MT)-1
 NG=NM(MT)
 GO TO 295
 302 M=KKW/KW-1
 IF((NM(MT)+M).LE.MNM) GO TO 307
 LARGE=3
 GO TO 345
 307 CONTINUE
 DO 304 L=1,M
 NM(MT)=NM(MT)+1
 NG=NM(MT)
 PAX(MT,NG)=PAX(MT,N)
 PAY(MT,NG)=PAY(MT,N)
 PAZ(MT,NG)=PAZ(MT,N)
 PAU(MT,NG)=PAU(MT,N)
 PAV(MT,NG)=PAV(MT,N)
 PAW(MT,NG)=PAW(MT,N)
 ER(MT,NG)=ER(MT,N)
 LCOL(MT,NG)=LCOL(MT,N)
 304 LPP(MT,NG)=KW
 305 LPP(MT,N)=KW
 Q=NUMCEL(K)
 J=NB(MT,Q)+1

ORIGINAL TAPE IS
OF POOR QUALITY

RUN5850
 RUN5860
 RUN5870
 RUN5880
 RUN5890
 RUN5900
 RUN5910
 RUN5920
 RUN5930
 RUN5940
 RUN5960
 RUN5970
 RUN5980
 RUN5990
 RUN6000
 RUN6010
 RUN6020
 RUN6030
 RUN6040
 RUN6050
 RUN6060
 RUN6070
 RUN6080
 RUN6090
 RUN6100
 RUN6110
 RUN6120
 RUN6130
 RUN6140
 RUN6150
 RUN6160
 RUN6170
 RUN6180
 RUN6190
 RUN6230
 RUN6240
 RUN6200
 RUN6210
 RUN6260
 RUN6270
 RUN6280
 RUN6290
 RUN6300
 RUN6310
 RUN6320
 RUN6330
 RUN6340
 RUN6350
 RUN6360

IF (J.LE.MNB) GO TO 308	RUN6370
IF (DEBUG (1)) WRITE (6,44) MT,K,MNB,AIME	
308 NB (MT,Q)=J	RUN6400
NBF (MT,Q)=NBF (MT,Q)+KW	RUN6410
LB (N)=Q	
GO TO 295	RUN6430
310 CONTINUE	
NBM (MT,1)=0	
DO 320 N=1,NBX	
M=NUMCEL (N)	
IF (M.EQ.0) GO TO 320	
A=NBF (MT,M)	RUN6480
DB (MT,M)=A*DFA (MT)/FNB (N)	RUN6490
NBM (MT,M+1)=NBM (MT,M)+NB (MT,M)	
NBN (M)=NBM (MT,M)	
320 CONTINUE	RUN6500
IF (NM (MT).GT.NMAX) NMAX=NM (MT)	RUN6510
DO 325 N=1,NG	
Q=LB (N)	
NBN (Q)=NBN (Q)+1	
NA=NBN (Q)	
325 LM (MT,NA)=N	
330 CONTINUE	RUN6520
IF (SAMP.LT.ITS) GO TO 335	
CALL ACCUM (NMC,NPB,FNB,NB,PAU,PAV,PAW,ER,TMP,TRP,XV,IV,ZV,LM,MSP,	
1NSP,LPF,NBF,NBM)	
SAMP=0	RUN6570
IF (TIME.LE.TST) GO TO 335	
CALL AVRGE (FNB,DB,DBA,NB,NBT,XV,IV,ZV,XVA,YVA,ZVA,TMP,TMPA,TRP,TRP	
1A,MSP,NSP,NBF,NBS)	
NAV=NAV+1	RUN6600
335 CPA=ELTIME (0)	
CPI=CPC+CPM+CPB+CPA	
CPJ=2.*CPI+5.	
340 CPUTYM=TFIND (0)	
IF (DEBUG (2)) WRITE (6,33) AIME,CPC,CPM,CPI,CPB,CPA, (NM (I),I=1,3),NMAX	
IF ((TIME.GE.TLIM).OR.(CPUTYM.LE.CPJ)) GO TO 345	
IF (PRT.LT.ITP) GO TO 280	
PRT=0	RUN6650
345 WRITE (6,30) AIME,KAWLS	RUN6660
IF (DEBUG (3)) WRITE (6,31) CPUTYM	RUN6670
WRITE (6,32) (NM (I),I=1,3)	
WRITE (6,34) ((NCOL (I,J),J=1,3),I=1,3)	
WRITE (6,35) (JNT (I),I=1,3)	
IF (LARGE.NE.0) GO TO 360	RUN6740
WRITE (6,36) NMAX	RUN6750
IF (.NOT.SAVE) GO TO 355	RUN6900
IF (PRT.NE.0.AND.CPUTYM.GT.CPJ.AND.TIME.LT.TLIM) GO TO 355	
WRITE (9) DENP,U,XREP,TRF,KAWLS,NL,NW,NH,MW,MH,LW,LH,NXA,NXB,NCA	
1 ,NCB,NFA,NPB,NBA,NHB,BW,BH,BWB,BHB,BWC,BHC,XLB,XLC,PI,NREG	RUN6930
2 ,S,SINANG,COSANG,AKN,NBX,RM,XR,ND,TIME,DTM,TL,ITS,ITP,TST	RUN6940
3 ,TLIM,RMA,RNU,DIR,XSTART,JNM,MNM,MNB,TR,BZC,CN7,DRF,PCF	RUN6950
4 ,FNA,HTF,INM,ITYPE,JTYPE,MJ,NAV,NMAX,NS,NWEDG,PRT,SAMP	RUN6960
5 ,BTA,C1,C2,C3,C7,C8,DAM,DPA,FL,DELANG,FDN,HTI,HTB,JNT,KNM	RUN6970
6 ,NM,WTM,C4,VRM,NCOL,LD,LF,LWF,RLD,CTI,CTR,CNI,CNR,LEV,SN	RUN6980

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7 ,ST,D1,D2,D3,D4,SSA,SSB,MS,NSP,NMB,NMC,NMP,NPB,NRAN,VELR
 8 ,IWS,TANGN,XLIM,COEFF,XCB,XS,YCB,TB,ALPHA,SIGMA,NTS,NTSF RUN7000
 9 ,UTL,UTT,VTS,HTS,HTSI,ENT,REM,ENTS,REMS,FTH,THETA,DTH,TMPARUN7010
 A ,DBA,NB,NBF,NBT,TMP,XV,XVA,YV,YVA,ZV,ZVA,T,DB,PNB,XC,YC,ZCRUN7020
 B ,NUMCEL,PAU,PAV,PAW,PAX,PAY,PAZ,PV,NTCV,NTCP,LPP,LCOL,LM RUN7030
 C,ETA,PHI,CHI,CN,CM,CNG,CMG,CN8,TRP,TRPA,THETAZ,NWEDGE,NSP,ANGLE,TF
 D,UTLI,UTTI,VTSI,ER,RMB,LKW,NBS,LB,NBM,NBN
 REWIND 9 RUN7040
 WRITE(6,50) RUN7050
 355 CONTINUE
 IF (TIME.LE.TST) GO TO 350
 DT=AIME-TI RUN6770
 CALL PRINT1(DT,COSANG,SINANG,RMA,RNU,DRF,FCP,HTP,FL,HTI,HTR,CTI,
 1CTR,CNI,CNR)
 CALL PRINT2(AKN,XSTART,DT,RNU,RMA,DRF,FCP,HTP,UTLI,UTTI,VTSI,HTSI,
 1DELANG,NWEDGE,XS,XCB,YCB,HTS,NTS,NTSF,UTL,UTT,VTS,LIMIT(3),
 2LIMIT(1),NSP)
 IF (NS.NE.0) CALL PRINT3(MSP,MJ,NS,NWEDG,LIMIT(3),LIMIT(1),
 1LIMIT(8),LIMIT(9), RMA,XS,IWS,MS,TANGN,NTSF,NTCP,NTCV,PV)
 CALL PRINT4(MSP,CHI,RNU,NSP,TRPA,NUMCEL,FDN,WTM,DB,NCB,NB,TMPA,XVA,
 1YVA,ZVA,1,NBT,XC,YC,ZC,LEV,LKW)
 GO TO 353 RUN6860
 350 CONTINUE
 CALL PRINT4(MSP,CHI,RNU,NSP,TRP,NUMCEL,FDN,WTM,DB,NB,TMP,XV,YV,ZV,
 10,NBF,XC,YC,ZC,LEV,LKW)
 353 IF (DEBUG(2)) WRITE(6,1)
 IF ((TIME.LT.TLIM).AND.(CPUTYM.GT.CPJ)) GO TO 280
 IF (IC.EQ.1COPY) RETURN
 IC=IC+1 RUN7080
 WRITE(6,2)
 WRITE(6,4)
 WRITE(6,5) (((I,J,L,(ENT(I,J,K,L),K=1,6),L=1,NWEDG),J=1,NSP),I=1,2)
 WRITE(6,6) ((J,L,ENTS(J,L),FTH(J,L),L=1,NWEDG),J=1,NSP)
 WRITE(6,2)
 WRITE(6,3) IC RUN7130
 CALL PRINTA(THETAZ,NWEDGE,TITLE,NAME,XCB,YCB,TB,ALPHA,SIGMA,LD,LF,RUN7140
 1XLIM,COEFF,LIMIT,NSP)
 CALL PRINTB(FNA,NSP,PNB,LEV,LWP,NB,RLD,XLIM,XC,YC,ZC,NB,NUMCEL,LKW
 1,NSP)
 SAVE=.FALSE. RUN7180
 GO TO 345 RUN7190
 360 WRITE(6,38) (DBG1(I,LARGE),I=1,3) RUN7200
 IF (REDO) GO TO 364
 IF (DUMP) CALL AREND(9) RUN7220
 STOP RUN7230
 364 CONTINUE
 IF (NEW) GO TO 365
 READ(9) DENP,U,XREF,TRF, KAWS,NL,NW,NH,MW,MH,LW,LH,NXA,NXB,NCA
 1 ,NCB,NFA,NFB,NHA,NHB,BW,BH,BWB,BHB,BWC,BHC,XLB,XLC,PI,NREG
 2 ,S,SINANG,COSANG,AKN,NBX,RM,XR,ND,TIME,DTM,TI,ITS,ITP,TST
 3 ,TLIM,RMA,RNU,DIR,XSTART,JNM,MNM,MNB,TR,BZC,CN7,DRF,FCF
 4 ,PNA,HTF,INM,ITYPE,JTYPE,MJ,NAV,NMAX,NS,NWEDG,PRT,SAMP
 5 ,BTA,C1,C2,C3,C7,C8,DAM,DPA,FL,DELANG,FDN,HTI,HTR,JNT,KNM
 6 ,NH,WTM,C4,VRM,NCOL,LD,LF,LWP,RLD,CTI,CTR,CNI,CNR,LEV,SN
 7 ,ST,D1,D2,D3,D4,SSA,SSB,MS,NSP,NMB,NMC,NMP,NPB,NRAN,VELR

8 ,IWS,TANGN,XLIM,COEFF,XCB,XS,YCB,TB,ALPHA,SIGMA,NTS,NTSF
 9 ,UTL,UTT,VTS,HTS,HTSI,ENT,REM,ENTS,REMS,PTH,THETA,DTH,TMPA
 A ,DBA,NB,NBF,NBT,TMP,XV,XVA,YV,YVA,ZV,ZVA,T,DB,FNB,XC,YC,ZC
 B ,NUMCEL,PAU,PAV,PAW,PAX,PAY,PAZ,FV,NTCV,NTCF,LPP,LCOL,L
 C,ETA,PHI,CHI,CN,CN,CNG,CMG,CN8,TRP,TRPA,THETAZ,NWEDGE,MSP,ANGLE,TP
 D,UTLI,UTTI,VTSI,ER,RMB,LKW,NBS,LB,NBM,NBN
 REWIND 9

365 JNM=9*JNM/10

RUN7260

ANM=INM

INM=9*ANM/10

DDN=.9*DDN

DRF=DRF/.9

FCP=FCP/.9

HTP=HTP/.9

DO 370 MM=1,MSP

FDN(MM)=FDN(MM)*INM/ANM

DO 366 KK=1,MSP

366 CN8(KK,MM)=CN8(KK,MM)*.9

DO 370 LT=1,NWEDG

ENTS(MM,LT)=ENTS(MM,LT)*INM/ANM

REMS(MM,LT)=0.0

DO 370 NK=1,2

DO 370 NJ=1,6

ENT(NK,MM,NJ,LT)=ENT(NK,MM,NJ,LT)*INM/ANM

370 REM(NK,MM,NJ,LT)=0.0

IF(NEW) GO TO 220

TST=TIME+TST

TI=-1.

PRT=ITP

WRITE(6,2)

WRITE(6,4)

WRITE(6,5) ((I,J,L,(ENT(I,J,K,L),K=1,6),L=1,NWEDG),J=1,MSP),I=1,2)

WRITE(6,6) ((J,L,ENTS(J,L),PTH(J,L),L=1,NWEDG),J=1,MSP)

WRITE(6,2)

IF((LARGE.EQ.2).OR.(LARGE.EQ.3)) GO TO 280

REDO=.FALSE.

GO TO 360

END

RUN7410

SUBROUTINE DIAG(N,ITEST,NUM)

DIAG010

REAL*8 PAB:*(10) // ' NWEDGE', ' NREG', ' ND', ' NPX',
 1 MNM', ' MNB', ' NBX', ' NS', ' MJ', ' MSP' /

DIAG040

DIAG050

DIAG060

DIAG070

DIAG080

FORMATS

32 FORMAT(9X,'ENT,REM,ENTS,REMS,PTH,THETA,DTH')
 42 FORMAT(//5X,43H ARRAY DIMENSIONS ARE ABOUT TO BE VIOLATED./) DIAG100
 44 FORMAT(5X,18H MAXIMUM VALUE IS 15,20H, WHEREAS YOU INPUT 15,3H (,DIAG110
 1A8,1H)) DIAG120
 56 FORMAT(/5X,78H IF YOU DESIRE TO USE THIS VALUE, THE FOLLOWING ARRADIAG130
 1YS MUST BE RE-DIMENSIONED./) DIAG140
 62 FORMAT(9X,'HTS,HTSI,NTS,NTSF,UTL,UTT,VTS') DIAG150
 64 FORMAT(9X,'XLIM,COEFF'//11X,'NOTE THAT THE XLIM ARRAY MUST BE DIMEDIAG160

1NSIONED TO 3 MORE THAN THE COEFF ARRAY.")	DIAG170
66 FORMAT(9X,'XCB,XS,YCB,TB,ALPHA,SIGMA')	DIAG180
68 FORMAT(9X,'DBA,NB,NBF,NBT,TMP,TMPC,XV,XVA,YV,YVA,ZV,ZVA,T,DB')	DIAG190
70 FORMAT(9X,'LPP,PAU,PAV,PAW,PAX,PAY,PAZ,LCOL')	DIAG200
72 FORMAT(8X,3H LM)	DIAG210
74 FORMAT(//5X,76H IF YOU CHANGE THE ARRAY DIMENSIONS, ALSO CHANGE THIS 1E 'LIMIT' DATA STATEMENT.)	DIAG220
75 FORMAT(9X,'ALL ARRAYS ASSOCIATED WITH SPECIES')	DIAG230
76 FORMAT(9X,'FNB,XC,YC,ZC,NUMCEL')	DIAG240
78 FORMAT(9X,'FV,NTCV,NTCF,MS,IWS,SL,DELS,TANGN')	DIAG250
80 FORMAT(9X,'PV,NTCV')	DIAG260

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WRITE(6,42)		DIAG270
WRITE(6,44) ITEST,NUM,PARAM(N)		DIAG280
WRITE(6,56)		DIAG290
GO TO (1,2,3,4,5,6,7,8,9,10),N		DIAG300
1 WRITE(6,62)		DIAG310
WRITE(6,32)		DIAG320
GO TO 11		
2 WRITE(6,64)	ORIGINAL EXPRESSION OF POOR QUALITY	DIAG340
GO TO 11		DIAG350
3 WRITE(6,66)		DIAG370
WRITE(6,62)		DIAG390
GO TO 11		DIAG400
4 WRITE(6,68)		DIAG420
GO TO 11		
5 WRITE(6,70)		DIAG440
GO TO 11		
6 WRITE(6,72)		DIAG460
GO TO 11		
7 WRITE(6,76)		DIAG480
GO TO 11		
8 WRITE(6,78)		DIAG500
GO TO 11		
9 WRITE(6,80)		DIAG520
GO TO 11		
10 WRITE(6,75)		
11 WRITE(6,74)		
STOP		DIAG540
END		DIAG550

SUBROUTINE PRINTA(THETAZ,NWEDGE,TITLE,NAME,XCB,YCB,TB,ALPHA,SIGMA,PRA0010	
1LD,LF,XLIM,COEFF,LIMIT,MSP)	
INTEGER TST,TLIM,TIME	PRA0030
LOGICAL SAVE,NEW	PRA0040
DIMENSION NWEDGE(2),LIMIT(1),TITLE(6),NAME(2),XCB(1),YCB(1),TB(1)	
DIMENSION ALPHA(3,1),SIGMA(3,1),LD(1),LF(1),XLIM(1),COEFF(4,1)	PRA0060
DIMENSION RNU(3),RMA(3),CHI(3),DIR(3,3),PHI(3,3),ETA(3,3)	
DIMENSION WTM(3),DAM(3,3),VELS(3),XSP(3)	
COMMON /FIRST/NL,NW,NH,MW,MH,LW,LH,NXA,NXB,NCA,NCB,NPA,NFB,NHA,NHB	PRA0070
COMMON /SECND/BW,BH	PRA0080
COMMON /THIRD/PI,NREG,S,SINANG,COSANG,AKN	PRA0090
COMMON /FIFTH/ND,TIME,DTM,TI,ITS,ITP,TST,TLIM,RMA,DIR	PRA100

COMMON /SIXTH/RMB,XSTART,JNN,MNN,MNB,NEW,SAVE,PERCNT,NSR,TR PRA0110
 COMMON/EIGHT/DENP,U,TP,ANGLE,TRF,CHI,PHI,ETA,WTM,DAM,VELR,XREF PRA0120
 DATA NOT/'NOT'/' PRA0130

----- PRA0140

----- PRA0150

----- PRA0160

----- PRA0170

FORMATS

ORIGINAL PAGE IS
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1 FORMAT(16X,40('''),T74,'I'//9X,'3-D',I2,'-FLUID PROGRAM - ') PRA0190
 2 FORMAT('+'',31X,A4) PRA0200
 3 FORMAT('+'',35X, 'A4.' - ' ,I2,' REGIONS',T74,'I',16(/T74,'I')) PRA0210
 4 FORMAT(7X,'FRONT OF BODY =',E12.4,' XSTART MAX HEIGHT =',E12.4,' 1RMB',T74,'I'/7X,'X-LIMIT',T37,'BODY COEFFICIENTS',T74,'I')
 6 FORMAT(5F14.6,3X,'I') PRA0240
 10 FORMAT(1X,72(''')) PRA0250
 12 FORMAT(//14X,'PARAMETERS OF SEGMENTS FOR BODY COLLISIONS',T96,'I'/18X,'X-COORD. TEMP. ALPHA1 ALPHA2 ALPHA3 SIGMA1 SIGMA2
 2 SIGMA3 AREAS',T96,'I')
 14 FORMAT(4X,E12.4,7F9.4,E12.4,T96,'I')
 16 FORMAT(23X,'WEIGHTING FACTORS'/1X,10I6,T96,'I')
 17 FORMAT(///25X,'ARRAY STORAGE USFD'/5X,I6,'*'*,10I6,T96,'I')
 18 FORMAT(1H1/17X,'LENGTH OF CELL IN MEAN-FREE-PATHS = ',F12.4,' BW' A,T76,'I'
 1/17X,'HEIGHT OF CELL IN MEAN-FREE-PATHS = ',F12.4,' BH',T76,'I'
 2/16X,'NUMBER OF L1 CELLS ALONG FLOW AXIS = ',I13,' NW',T76,'I'
 3/17X,'NUMBER OF L1 CELLS IN RADIAL DIR. = ',I13,' NH',T76,'I'
 4/21X,'NUMBER OF LEVELS OF CELL SIZE = ',I13,' NL',T76,'I')
 20 FORMAT(11X,'NUMBER OF L1 CELLS IN FRONT OF L2 CELLS = ',I13,' NFA'
 1,T76,'I'/15X,'NUMBER OF AXIAL SUBDIVIDED L1 CELLS = ',I13,' NCA',T
 276,'I'/14X,'NUMBER OF RADIAL SUBDIVIDED L1 CELLS = ',I13,' NHA',T
 36,'I'/16X,'NUMBER OF L2 CELLS ALONG FLOW AXIS = ',I13,' MW',T76,'I'
 4'/17X,'NUMBER OF L2 CELLS IN RADIAL DIR. = ',I13,' MH',T76,'I')
 22 FORMAT(11X,'NUMBER OF L2 CELLS IN FRONT OF L3 CELLS = ',I13,' NFB'
 1,T76,'I'/15X,'NUMBER OF AXIAL SUBDIVIDED L2 CELLS = ',I13,' NCB',T
 276,'I'/14X,'NUMBER OF RADIAL SUBDIVIDED L2 CELLS = ',I13,' NHB',T
 36,'I'/16X,'NUMBER OF L3 CELLS ALONG FLOW AXIS = ',I13,' LW',T76,'I'
 4'/17X,'NUMBER OF L3 CELLS IN RADIAL DIR. = ',I13,' LH',T76,'I')
 23 FORMAT(3X,'NUMBER OF AZIMUTHAL WEDGES IN LOWER ',I3,' DEGREES = ',I13,' NWEDGE1 I'/3X,'NUMBER OF AZIMUTHAL WEDGES IN UPPER ',I3,' DE
 2GREES = ',I13,' NWEDGE2 I') PRA0470
 24 FORMAT(16X,'BASIC TIME INTERVAL FOR COLLISIONS = ',E13.4,' DTM
 1 I'/8X,'TIME INTERVAL FOR SAMPLING FLOW FIELD INFO = ',E13.4,' DTS
 2 I'/24X,'TIME INTERVAL FOR PRINTING = ',E13.4,' DTP I'/9X,
 3' TIME TO STEADY-STATE CONDITIONS (ASSUMED) = ',E13.4,' TST I'/
 419X,'TIME AT WHICH RUN IS TERMINATED = ',E13.4,' TLIM I')
 26 FORMAT(9X,'INITIAL NUMBER OF MOLECULES - EITHER TYPE = ',I13,' INM
 1 I'/9X,'MAXIMUM NUMBER OF MOLECULES - EITHER TYPE = ',I13,' MN
 2M I'/1X,'MAX NUMBER OF MOLECULES IN ANY CELL - EITHER TYPE = ',PRA0570
 3I13,' MNB I')
 27 FORMAT(//22X,'VELOCITY OF FREE STREAM FLOW = ',E13.4,' U',T76,'I'/19X,'SPEED RATIO OF FREE STREAM FLOW = ',E13.4,' S',T76,'I'/19X,'MAC
 AH NUMBER OF FREE STREAM FLOW = ',E13.4,' N',T76,'I'/19X,'SPECIFIC H
 BEAT RATIO (CALCULATED) = ',E13.4,' GAMMA',T76,'I'/35X,'ANG
 2LE OF ATTACK = ',F13.4,' ANGLE I'/16X,'NUMBER DENSITY OF FREE ST

3 STREAM FLOW =', E13.4, ' N', T76, ' I' / 19X, ' TEMPERATURE OF FREE STREAM FL
 40W =', F13.4, ' TF', T76, ' I' / 16X, ' MOLE FRACTIONS OF FREE STREAM FLOW
 5 =', 3E13.4, ' RNU I' / 16X, ' MOLECULAR WEIGHTS OF SPECIES ABOVE =', 3F13
 6.4, ' RMA I')

28 FORMAT (//10X, 'REFERENCE TEMPERATURE FOR MOLECULAR DATA =', F12.4, '
 1TRF', T90, ' I' / 14X, ' CROSS-SECTION', 26X, ' TEMP EXPONENT', T90, ' I' / 3 (3X,
 23E12.4, 3X, 3F12.6, T90, ' I' /) / 5X, ' CHI/2-1', 11X, ' ROTATIONAL PARAMETER
 3 PHI', T90, ' I' / 3 (F12.4, 5X, 3F12.6, T90, ' I' /))

29 FORMAT (9X, 'DATA SAVED ON TAPE 9')

PRA0670

30 FORMAT (//31X, 'REF MOLECULAR SPEED =', E13.4, ' VELR', T76, ' I' / 20X, ' SP
 ECIES FREE STREAM MOLECULAR SPEEDS', T76, ' I' / 14X, 3E16.6, T76, ' I' / 26X
 2, 'REFERENCE MEAN FREE PATH =', E13.4, ' XREF', T76, ' I' / 26X, 'SPECIES M
 3EAN FREE PATHS', T76, ' I' / 14X, 3E16.6, T76, ' I' / 11X, 'LONGITUDINAL KNUDS
 4EN NUMBER (CALCULATED) =', E13.4, ' AKN', T76, ' I' / 13X, 'TRANSVERSE KNUD
 5SEN NUMBER (CALCULATED) =', E13.4, ' AKT', T76, ' I')

PRA0680

----- PRA0690

PRA0700

IARRAY=708+LIMIT(3)*(32+56*LIMIT(1))+20*LIMIT(2)+LIMIT(4)*(120+4*LIM PRA0710
 1LIMIT(6))+56*LIMIT(5)+LIMIT(8)*(68+96*LIMIT(9))+20*LIMIT(7)+224*LIMIT PRA0720
 2IT(1) PRA0730

WRITE(6, 1) MSP

PRA0750

IF (NEW) WRITE(6, 2) NOT

PRA0760

WRITE(6, 3) NAME, TITLE, NREG

PRA0770

WRITE(6, 4) XSTART, RMB

PRA0780

DO 100 I=1, NREG

PRA0790

100 WRITE(6, 6) XLIM(I+2), (COEFF(J, I), J=1, 4)

PRA0800

WRITE(6, 10)

PRA0810

WRITE(6, 12)

PRA0820

DO 110 I=1, ND

110 WRITE(6, 14) XCB(I), TB(I), (ALPHA(J, I), J=1, 3), (SIGMA(J, I), J=1, 3), YCB
 1(I)

PRA0850

WRITE(6, 10)

PRA0860

WRITE(6, 16) (LD(N), LP(N), N=1, 5)

WRITE(6, 17) IARRAY, (LIMIT(I), I=1, 10)

PRA0880

WRITE(6, 18) BW, BH, NW, NH, NL

PRA0890

IF (NL.GT.1) WRITE(6, 20) NFA, NCA, NHA, MW, MH

PRA0900

IF (NL.GT.2) WRITE(6, 22) NFB, NCB, NHB, LW, LH

PRA0910

IETAZ=THETAZ

PRA0920

JETAZ=180-IETAZ

PRA0930

WRITE(6, 23) IETAZ, NWEDGE(1), JETAZ, NWEDGE(2)

PRA0940

DT5=DTM*ITS

PRA0950

DTP=DTM*ITP

PRA0960

AST=DTM*TST

PRA0970

ALIM=DTM*TLM

PRA0980

CHT=0.0

DO 120 J=1, MSP

120 CHT=CHT+CHI(J)*RNU(J)

PRA0980

GAMMA=(7.+2.*CHT)/(5.+2.*CHT)

AM=S*SQRT(2./GAMMA)

WRITE(6, 24) DTM, DT5, DTP, AST, ALIM

WRITE(6, 26) JNM, MNM, MNB

WRITE(6, 27) U, S, AM, GAMMA, ANGLE, DENP, TF, (RNU(I), I=1, 3), (RMA(I), I=1, 3)

1)

WRITE(6, 28) TRP, ((DIR(I, K), K=1, 3), (ETA(I, K), K=1, 3), I=1, 3), (CHI(I),

```

1 (PHI (I,K) , K=1,3) , I=1,3)
DO 210 I=1,3
VELS(I)=0.0
210 XSP(I)=0.0
DO 220 J=1,MSP
VELS(J)=VELR/SQRT(WTM(J))
XT=0.0
DO 215 M=1,MSP
215 XT=XT+RNU(M)*DAM(J,M)*SQRT(1.+WTM(J)/WTM(M))
220 XSP(J)=1.414214*XREF/XT
AKT=1./RMB
WRITE(6,30) VELR, (VELS(I), I=1,3), XREF, (XSP(I), I=1,3), AKN, AKT
IF(SAVE) WRITE(6,29)
RETURN
END

;-----  

SUBROUTINE PRINTB(FNA,MSP,FNB,LEV,LWF,NM,RLD,XLIM,XC,YC,ZC,NB,
1NUMCEL,LKW,N)
INTEGER*2 LKW,NB,NUMCEL
DIMENSION FNB(1),LEV(1),LWF(1),NM(1),RLD(1),XLIM(1),XC(1)
DIMENSION YC(1),ZC(1),NB(N,1),NUMCEL(1),LKW(1)
COMMON /FIRST/NL
COMMON /THIRD/PI,NREG,S,SINANG,COSANG,AKN
COMMON /FORTH/NBX
1 FORMAT(1H1) PRB0050
2 FORMAT(2X,-----CELL GEOMETRY-----PRB0090
1-----'2X, 'BOX LEVEL POSITION OF CENTER VOLUME WPRB0100
2EIGHTING POPULATION'/2X,'NUM.',12X,'X',7X,'Y THETA',12X,'FACTOPRB0110
3R ',15X,' CELL#')
3 FORMAT(1X,I4,I5,3X,2F8.3,F7.1,E12.3,2X,I2,4X,3I5,3X,I4)
4 FORMAT(2X,-----TOTALS-----',E12.4,8X,3I5) PRB0150
WRITE(6,1)
WRITE(6,2)
DO 200 I=1,NBX PRB0170
IF(NUMCEL(I).EQ.0) GO TO 200
X=(XC(I)-XLIM(2))*AKN
Y=YC(I)*AKN
LEVEL=1 PRB0180
IF(NL.LT.2) GO TO 120 PRB0190
IF(I.LT.LEV(1)) GO TO 120 PRB0200
LEVEL=2 PRB0210
IF(NL.LT.3) GO TO 120 PRB0220
IF(I.LT.LEV(2)) GO TO 120 PRB0230
LEVEL=3 PRB0240
120 CONTINUE PRB0250
J=NUMCEL(I)
M1=NB(1,J)
M2=0 PRB0260
IF(MSP.GE.2) M2=NB(2,J)
M3=0 PRB0320
IF(MSP.GE.3) M3=NB(3,J)
140 WRITE(6,3) I,LEVEL,X,Y,ZC(I),FNB(I),LKW(J),M1,M2,M3,J PRB0370
200 CONTINUE
NM2=0
IF(MSP.GE.2) NM2=NM(2)

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```

NM3=0
IF(MSP.GE.3) NM3=NM(3)
WRITE(6,4) PNA,NM(1),NM2,NM3
RETURN
END

SUBROUTINE SIMPSN(A,B,L,INTGRL,PERCNT,COEFF,PIROOT,SUM1,FUN)      SIMP010
REAL INTGRL
DIMENSION COEFF(4,1)                                                 SIMP020
                                                               SIMP030
                                                               SIMP040
                                                               SIMP050
                                                               SIMP060
                                                               SIMP070
                                                               SIMP080
                                                               SIMP090
                                                               SIMP100
THE PURPOSE OF THIS SUBROUTINE IS TO PERFORM A SIMPSON'S RULE      SIMP110
INTEGRATION.                                                       SIMP120
FORMAT                                                       SIMP130
4 FORMAT(/30H TOO MANY ITERATIONS. TEST IS E15.7,14H, INTEGRAL IS E1SIMP110
15.7,14H IN THE RANGE E15.7,4H TO E15.7)                           SIMP120
                                                               SIMP130
                                                               SIMP140
PREV=0.0
N=4
N1=3
G=.5*(B-A)+A
SUME=FUN(G,PIROOT,L,COEFF)
SUMO=0.0
K=0
235 SUME=SUME+SUMO
N1=N1+K
K=N1+1
IF(K.LT.5000) GO TO 237
WRITE(6,4) TEST,PREV,A,B
STOP
237 SUMO=0.0
DO 240 I=1,N1,2
C=(I*(B-A))/K+A
240 SUMO=SUMO+FUN(G,PIROOT,L,COEFF)
INTGRL=(SUM1+4.*SUMO+2.*SUME)/(3.*K)
TEST=ABS(2.-4.*PREV/(INTGRL+PREV))
PREV=INTGRL
IF(TEST.GT.PERCNT) GO TO 235
RETURN
END

FUNCTION FNCTN(X,PIROOT,L,COEFF)                                     FNCN010
DIMENSION COEFF(4,1)                                                 FNCN020
                                                               FNCN030
                                                               FNCN040
                                                               FNCN050
                                                               FNCN060
THE PURPOSE OF THIS FUNCTION IS TO EVALUATE THE INTEGRAND USED IN FNCN070
THE SIMPSON-S RULE INTEGRATION ROUTINE.                               FNCN080
                                                               FNCN090
                                                               FNCN100
                                                               FNCN110
A=COEFF(1,L)
B=COEFF(2,L)
C=COEFF(3,L)
D=COEFF(4,L)
AA=4.* (A-B) * (A*X+C) *X+C*C-4.*B*D

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IP(AA.LT.0.) AA=0.                                              FNCN120
FNCTN=SQRT(AA)                                                 FNCN130
RETURN                                                       ORIGINAL PAGE IS
END                                                       OF POOR QUALITY
FNCN140
FNCN150

SUBROUTINE HEIGHT(X,R,L,COEFF,I)                                 HGHT010
DIMENSION COEFF(4,1),DBG2(2,3)                                 HGRT020
DATA DBG2/'SBTR','CT ','GAS ',' ','RUN',' '/
HGHT030
HGHT040

THE PURPOSE OF THIS SUBROUTINE IS TO COMPUTE THE R-COORDINATE OF HGHT050
THE BODY CURVE AT ANY GIVEN X-COORDINATE, IN THE X-R PLANE. HGHT060
HGHT070

2 FORMAT(//'* MESSAGE ',I2,E17.8,3X,2A4)                      HGHT080
ARG=-(COEFF(1,L)*X+COEFF(3,L))*X+COEFF(4,L))/COEFF(2,L)
HGHT090
IF(ARG.GE.0.) GO TO 100
HGHT100
WRITE(6,2)L,X,(DBG2(M,I),M=1,2)
HGHT110
ARG=0.
HGHT120
100 R=SQRT(ARG)
HGHT130
RETURN
HGHT140
END
HGHT150

FUNCTION FNCTM(ARG,PIROOT,L,COEFF)                            FNCM010
DIMENSION COEFF(4,1)                                         FNCM020
D=0.                                                       FNCM030
FNCM040
IF(ABS(ARG).LT.10.) D=EXP(-ARG*ARG)/PIROOT
FNCM050
E=0.
FNCM060
IF(ARG.GT.-10.) E=ARG*ERRF(ARG)
FNCM070
TEM=.5*(D+E)
FNCM080
FNCTM=TEM
FNCM090
RETURN
FNCM100
END

SUBROUTINE CELL(TH,LEV,A,B,K,KH,XO,I,J,DELANG,NWEDGE,XC,YC,ZC,FNB) CELLO10
DIMENSION DELANG(1),NWEDGE(1),XC(1),YC(1),ZC(1),FNB(1)          CELLO20
COMMON /THIRD/PI                                              CELLO30
CELLO40
CELLO50

THE PURPOSE OF THIS SUBROUTINE IS TO                               CELLO60
1. COMPUTE THE VOLUME OF EACH CELL (ALL 3 POSSIBLE LEVELS)    CELL070
AND STORE THE RESULT IN THE ARRAY CALLED 'FNB'.                CELL080
2. COMPUTE THE X, R, AND THETA COORDINATES OF THE CENTER OF    CELL090
EACH CELL (ALL 3 POSSIBLE LEVELS) AND STORE THE RESULTS INCELL100
ARRAYS CALLED 'XC', 'YC', AND 'ZC'.                                CELL110
CELL120
CELL130

INDEX=I+J                                                       CELL140
ZO=0.                                                       CELL150
DO 120 MT=1,LEV                                              CELL160
ICNT=NWEDGE(MT)                                              CELL170
ANGLE=DELANG(MT)                                              CELL180
FACTOR=ANGLE/180.*PI*B*B*A                                    CELL190
Z=ZO-.5*ANGLE                                                 CELL200
DO 110 L=1,ICNT                                              CELL210
Z=Z+ANGLE                                                 CELL220
Y=-.5*B                                                       CELL230

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DO 110 N=1,KH                                CELL240
X=XO-.5*A                                     CELL250
Y=Y+B                                         CELL260
DO 110 M=1,K                                     CELL270
X=X+A                                         CELL280
INDEX=INDEX+1                                    CELL290
XC(INDEX)=X                                     CELL300
YC(INDEX)=Y                                     CELL310
ZC(INDEX)=Z                                     CELL320
110 FNB(INDEX)=FACTOR*(2*N-1)                  CELL330
ZO=TH                                         CELL340
120 CONTINUE                                     CELL350
RETURN                                         CELL360
END                                            CELL370

SUBROUTINE ZERO(NWIDE,NHI,NBEG,NLONG,MUP,NAREA,ICNT,FNB)    ZERO010
DIMENSION FNB(1)                                     ZERO020
-----ZERO030
-----ZERO040
THIS SUBROUTINE SETS THE SIZES TO ZERO OF THOSE CELLS WHICH ARE ZERO050
TO BE SUBDIVIDED INTO SMALLER CELLS.                ZERO060
-----ZERO070
-----ZERO080
NGO=NBEG+1                                         ZERO090
NSTOP=NBEG+NLONG                                    ZERO100
DO 110 N=NGO,NSTOP                                ZERO110
DO 110 M=1,MUP                                     ZERO120
DO 110 L=1,ICNT                                    ZERO130
INDEX=NWIDE*(NHI*(L-1)+M-1)+N+NAREA               ZERO140
110 FNB(INDEX)=0.0                                 ZERO150
RETURN                                         ZERO160
END                                            ZERO170

SUBROUTINE SBTRCT(NGO,NTEMP,NSTOP,BWIDTH,BHITE,DELANG,XC,YC,FNB,XLSBCT010
1IM,COEFF)                                         SBCT020
DIMENSION DELANG(1),XC(1),YC(1),FNB(1),XLIM(1),COEFF(4,1)    SBCT030
COMMON /THIRD/PI,NREG                                SBCT040
-----SBCT050
-----SBCT060
THIS SUBROUTINE SUBTRACTS FROM EACH CELL SIZE ('FNB' ARRAY) THAT SBCT070
PORTION OCCUPIED BY THE BODY.                      SBCT080
-----SBCT090
-----SBCT100
FACTOR=PI*DELANG(1)/180.                           SBCT110
DO 150 N=NGO,NSTOP                                SBCT120
IF(FNB(N).LE.0.) GO TO 150                         SBCT130
DFNB=.005*FNB(N)                                    SBCT140
SLICE=.01*BWIDTH                                    SBCT150
IF(N.GT.NTEMP) FACTOR=PI*DELANG(2)/180.            SBCT160
X=XC(N)-.5*(BWIDTH+SLICE)                         SBCT170
YBOT=YC(N)-.5*BHITE                                SBCT180
YTOP=YBOT+BHITE                                    SBCT190
DO 130 M=1,100                                     SBCT200
X=X+SLICE                                         SBCT210
IF(X.LE.XLIM(2)) GO TO 130                         SBCT220

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DO 120 L=1,NREG           SBCT230
IF(X.LT.XLIM(L+2)) GO TO 125   SBCT240
120 CONTINUE                SBCT250
GO TO 130                  SBCT260
125 CALL HEIGHT(X,YBODY,L,COEFF,1)   SBCT270
IF(YBODY.LE.YBOT) GO TO 130   SBCT280
YTEMP=YTOP                 SBCT290
IF(YBODY.LT.YTOP) YTEMP=YBODY   SBCT300
FNB(N)=FNB(N)-SLICE*(YTEMP*YTEMP-YBOT*YBOT)*FACTOR   SBCT310
130 CONTINUE                SBCT320
IF(FNB(N).LT.DFNB) FNB(N)=0.   SBCT330
150 CONTINUE                SBCT340
RETURN                      SBCT350
END                         SBCT360

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SUBROUTINE IMPACT(RM,G1,G2,G3,ET,EI,PHI,CHI,ETA,XM,CIM)
COMMON/THIRD/PI
IF(PHI.EQ.0.) GO TO 20
IF(CHI.EQ.0.) GO TO 20
DP=PHI*CHI-1.
DS=PHI*(2.-.5*ETA)-1.
E=ET+EI
10 X=RAND(0)
IF(X.EQ.0.0) GO TO 10
XT=X**DF*(1.-X)**DS
IF(XT.GT.XM) GO TO 15
CIM=CIM+XT
IF(CIM.LT.XM) GO TO 10
CIM=CIM-XM
15 ET=(1.-PHI)*ET+(1.-X)*PHI*E
EI=(1.-PHI)*EI+X*PHI*E
20 GP=SQRT(ET/RM)
EP=2.*PI*RAND(0)
CSX=2.*RAND(0)-1.
SSX=SQRT(1.-CSX**2)
G1=GP*CSX
G2=GP*SSX*COS(EP)
G3=GP*SSX*SIN(EP)
RETURN
END

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SUBROUTINE GAS(NWEDGE,THETAZ,DELANG,NWEDGE,BTA,C1,DFA,NM,RLD,LWP,FNGAS0010
1B,DB,NB,NBF,LPF,PAU,PAV,PAX,PAY,PAZ,XLIM,COEFF,LM,I2,I3,LARGE,GAS0020
2MMN,MNB,DEBUG1,LCOL,NUMCEL,IP,ER,CHI,CNG,CMG,I,LB,NBM,NBN)
INTEGER*2 LM(I,1),LPP(I,1),LCOL(I,1),LB(1),NBM(I,1),NBN(1)
INTEGER*2 NB,NBF,NUMCEL
INTEGER Q
LOGICAL DUMP,DEBUG1
DIMENSION DELANG(2),NWEDGE(2),NUMCEL(1)
DIMENSION BTA(1),C1(1),DFA(1),NM(1),RLD(1),LWP(1),FNB(1),CHI(1)
DIMENSION DB(I,1),NB(I,1),NBF(I,1),PAU(I,1),PAV(I,1),PAW(I,1)
DIMENSION PAX(I,1),PAY(I,1),PAZ(I,1),ER(I,1),COEFF(4,1),XLIM(1)
DIMENSION CNG(1),CMG(1)
COMMON /FIRST/NL,NW,NH,MW,MH,LW,LH,NXA,NXB
COMMON /SECND/BW,BH,BWB,BHB,BWC,BHC,XLB,XLC

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GAS0050
GAS0060
GAS0070

GAS0110
GAS0120


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C9=C9+R/A                                GAS0580
IF(C9.LT.RWFM) GO TO 140                 GAS0590
C9=C9-RWFM                               GAS0600
IF(X.LE.XLIM(2)) GO TO 159                 GAS0610
DO 152 L=1,NREG                          GAS0620
IF(X.LT.XLIM(L+2)) GO TO 154             GAS0630
152 CONTINUE                               GAS0640
GO TO 159                                 GAS0650
154 CALL HEIGHT(X,YBODY,L,COEFF,2)        GAS0660
IF(R.LT.YBODY) GO TO 130                 GAS0670
159 PAX(MT,N)=X                           GAS0680
D=PI*RAND(0)                            GAS0690
PAY(MT,N)=R*COS(D)                      GAS0700
PAZ(MT,N)=R*SIN(D)                      GAS0710
TANG=180.* (1.-D/PI)                     GAS0720
IWDGE=TANG/DELANG(1)                     GAS0730
IF((IWDGE.GE.NWEDGE(1)).AND.(DELANG(2).NE.0.)) IWDGE=(TANG-THETAZ)
1/DELANG(2)+NWEDGE(1)                    GAS0750
IF(IWDGE.GE.NWEDG) IWDGE=NWEDG-1        GAS0760
L=X/BW+1.                                GAS0770
IF(L.GT.NW) L=NW                          GAS0780
M=R/BH                                    GAS0790
IF(M.GE.NH) M=NH-1                      GAS0800
K=(IWDGE*NH+M)*NW+L                     GAS0810
IF(K.LE.NXA) GO TO 160                  GAS0820
WRITE(6,2)L,M,K,MT,N,DELANG,NWEDGE,D,TANG,IWDGE,NH,NW
IF(DUMP) CALL ABEND(11)                  GAS0830
STOP                                     GAS0840
160 IF(NL.EQ.1) GO TO 162                 GAS0850
IF(IWDGE.GE.NWEDGE(1)) GO TO 162
IF(FNB(K).GT.0.) GO TO 162
L=(X-XLB)/BWB+1.
IF(L.GT.MW) L=MW
M=R/BHB
IF(M.GE.MH) M=MH-1
K=(IWDGE*MH+M)*MW+L+NXA
IF(K.LE.NXA+NXB) GO TO 161
WRITE(6,2)L,M,K,MT,N,DELANG,NWEDGE,D,TANG,IWDGE,MH,MW,NXA
IF(DUMP) CALL ABEND(12)
STOP
161 IF(NL.EQ.2) GO TO 162                 GAS0870
L=(X-XLC)/BWC+1.
IF(FNB(K).GT.0.) GO TO 162
IF(L.GT.LW) L=LW
M=R/BHC
IF(M.GE.LH) M=LH-1
K=(IWDGE*LH+M)*LW+L+NXA+NXB
IF(K.LE.NBX) GO TO 164
WRITE(6,2)L,M,K,MT,N,DELANG,NWEDGE,D,TANG,IWDGE,LH,LW,NXA,NXB
1B
IF(DUMP) CALL ABEND(13)
STOP
162 IF(FNB(K).GT.0.) GO TO 164
WRITE(6,3)L,M,K,MT,N,FNB(K)
IF(DUMP) CALL ABEND(14)                  GAS1050
                                         GAS1060
                                         GAS1070
                                         GAS1080
                                         GAS1090
                                         GAS1100
                                         GAS1110

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STOP	ORIGINAL PAGE IS	GAS1120
164 Q=NUMCEL(K)	OF POOR QUALITY	GAS1130
IF (Q.GT.0) GO TO 165		GAS1140
WRITE(6,2) Q,L,M,K,N,DELANG,NWEDGE,NWEDG,D,TANG,IWDGE,LH,LW,NXA,NXBGAS1150		
1,X,R		GAS1160
IF (DUMP) CALL ABEND(15)		GAS1170
STOP		GAS1180
165 J=NB(MT,Q)+1		GAS1190
KW=LWP(LA)		GAS1200
LPP(MT,N)=KW		GAS1210
LCOL(MT,N)=0		GAS1220
LLC=LLC+KW		GAS1230
IF (J.LE.MNB) GO TO 166		GAS1240
IF (DEBUG1) WRITE(6,4) MT,Q,MNB		GAS1250
GO TO 167		GAS1260
166 NB(MT,Q)=J		GAS1270
LB(N)=Q		
NBF(MT,Q)=NBF(MT,Q)+KW		GAS1290
167 IF (LLC.LT.LL) GO TO 110		GAS1300
NM(MT)=N		GAS1310
NBM(MT,1)=0		
DO 170 Q=1,NBX		GAS1330
N=NUMCEL(Q)		
IF (N.EQ.0) GO TO 170		
A=NBF(MT,N)		GAS1360
DB(MT,N)=A*DFA(MT)/FNB(Q)		GAS1370
NBM(MT,N+1)=NBM(MT,N)+NB(MT,N)		
NBN(N)=NBM(MT,N)		
170 CONTINUE		GAS1380
NG=NM(MT)		
DO 175 N=1,NG		
Q=LB(N)		
NBN(Q)=NBN(Q)+1		
NA=NBN(Q)		
175 LM(MT,NA)=N		
180 CONTINUE		GAS1390
RETURN		GAS1400
190 LARGE=1		GAS1410
RETURN		GAS1420
END		GAS1430
SUBROUTINE FLOW(NWEDG,MNM,LARGE,BTA,C1,C7,C8,D1,D2,D3,D4,DTH,NM,SNFL00010		
1,ST,THETA,LWP,RLD,PTH,ENTS,REMS,SSA,SSB,PAU,PAV,PAW,PAX,PAY,PAZ,LPPFL00020		
2P,ENT,REM,LCOL,IP,ER,CHI,CNG,CMG,I)		
INTEGER*2 LPP,LCOL		FLO0040
DIMENSION BTA(1),D1(1),D2(1),D3(1),D4(1),DTH(1),NM(1),SN(1),ST(1)		FLO0050
DIMENSION C1(1),C7(1),C8(1),LWP(1),RLD(1),THETA(1),SSA(2,1)		FLO0060
DIMENSION SSB(2,1),PAU(I,1),PAV(I,1),PAW(I,1),PAX(I,1),PAY(I,1)		
DIMENSION PAZ(I,1),LPP(I,1),ENT(2,3,6,1),REM(2,3,6,1),LCOL(I,1)		
DIMENSION ENTS(3,1),REMS(3,1),PTH(3,1),ER(I,1),CHI(1),CNG(1),		
1CMG(1)		
COMMON /THIRD/PI		FLO0100
COMMON /FORTH/NBX,RM,XR		FLO0110
-----		FLO0120
THE PURPOSE OF THIS SUBROUTINE IS TO ADD A NEW BATCH OF MOLECULESFLO0130		

TO THE SAMPLE THROUGH THE UPSTREAM BOUNDARY.

FLO0140

FLO0150

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DO 370 MT=1,IP
ARG=SN(MT)
XGO=0.
E=1.
DO 180 NT=1,2
SM=AMAX1(0.,ARG-4.)
SSM=AMAX1(0.,ARG)
TEMPC=0.
DO 170 LA=1,6
TEMPB=RLD(LA)*RLD(LA)
AY=TEMPB-TEMPC
C=TEMPC
TEMPC=TEMPB
DO 170 K=1,NWEDG
AM=ENT(NT,MT,LA,K)+REM(NT,MT,LA,K)
M=AM
AMM=M
REM(NT,MT,LA,K)=AM-AMM
IF(M.EQ.0) GO TO 170
DY=AY/AMM
DO 160 N=1,M
IF(NM(MT).GE.MNM) GO TO 380
NM(MT)=NM(MT)+1
NMX=NM(MT)
R=SQRT(C+DY*(N+RAND(0)-1.))
D=(THETA(K)+RAND(0)*DTH(K))*PI/180.
PAY(MT,NMX)=R*COS(D)
PAZ(MT,NMX)=R*SIN(D)
LPF(MT,NMX)=LWF(LA)
LCOL(MT,NMX)=0
130 V=SM+RAND(0)*(SSM+4.-SM)
C1(MT)=C1(MT)+2.*V*EXP(SSB(NT,MT)+2.*ARG*V-V*V)/SSA(NT,MT)
IF(C1(MT).LT.1.) GO TO 130
C1(MT)=C1(MT)-1.
PAU(MT,NMX)=E*V/BTA(MT)
140 V=8.*RAND(0)-4.
C7(MT)=C7(MT)+EXP(-V*V)
IF(C7(MT).LT.1.) GO TO 140
C7(MT)=C7(MT)-1.
PAV(MT,NMX)=(V+ST(MT))/BTA(MT)
150 V=8.*RAND(0)-4.
C8(MT)=C8(MT)+EXP(-V*V)
IF(C8(MT).LT.1.) GO TO 150
C8(MT)=C8(MT)-1.
PAW(MT,NMX)=V/BTA(MT)
ER(MT,NMX)=0.0
IF(CHI(MT).LE.-1.) GO TO 160
125 X=9.*RAND(0)
IF(X.LE.0.0) GO TO 125
XT=X**CHI(MT)*EXP(-X)
IF(XT.GE.CMG(MT)) GO TO 126
CNG(MT)=CNG(MT)+XT
IF(CNG(MT).LT.CMG(MT)) GO TO 125

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ORIGINAL PAGE IS
OF POOR QUALITY

FLO0170
FLO0180
FLO0190
FLO0200
FLO0210
FLO0220
FLO0230
FLO0240
FLO0250
FLO0260
FLO0270
FLO0280
FLO0290
FLO0300
FLO0310
FLO0320
FLO0330
FLO0340
FLO0350
FLO0360
FLO0370
FLO0390
FLO0400
FLO0410
FLO0420
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FLO0470
FLO0480
FLO0490
FLO0500
FLO0510
FLO0520
FLO0530
FLO0540
FLO0550
FLO0560
FLO0570
FLO0580
FLO0590
FLO0600

CNG(MT)=CNG(MT)-CMG(MT)
 126 ER(MT,NMX)=X FLO0610
 160 PAX(MT,NMX)=XGO FLO0620
 170 CONTINUE FLO0630
 ARG=-ARG FLO0640
 XGO=XR FLO0650
 E=-1. FLO0660
 180 CONTINUE FLO0670
 DO 370 K=1,NWEDG FLO0680
 AM=ENTS(MT,K)+REMS(MT,K) FLO0690
 M=AM FLO0700
 AMM=M FLO0710
 REMS(MT,K)=AM-AMM FLO0720
 IF(M.EQ.0) GO TO 370 FLO0730
 DX=XR/AMM FLO0740
 DO 365 N=1,M FLO0750
 IF(NM(MT).GE.MNM) GO TO 380 FLO0770
 NM(MT)=NM(MT)+1 FLO0780
 NMX=NM(MT) FLO0790
 330 PAX(MT,NMX)=(N-1.+RAND(0))*DX FLO0800
 TH=(THETA(K)+RAND(0)*DTH(K))*PI/180. FLO0810
 A=COS(TH) FLO0820
 B=SIN(TH) FLO0830
 SM=ST(MT)*A FLO0840
 C=0. FLO0850
 IF(ABS(SM).LT.10.) C=EXP(-SM*SM) FLO0860
 D=0. FLO0870
 IF(SM.GT.-10.) D=SQRT(PI)*SM*ERRF(SM) FLO0880
 D1(MT)=D1(MT)+(C+D)/FTH(MT,K) FLO0890
 IF(D1(MT).LT.1.) GO TO 330 FLO0900
 D1(MT)=D1(MT)-1. FLO0910
 PAZ(MT,NMX)=-RM*A FLO0920
 PAZ(MT,NMX)=RM*B FLO0930
 VNM=.5*SM+SQRT(.25*SM*SM+.5) FLO0940
 VM=AMAX1(0.,SM-4.) FLO0950
 340 V=VM+RAND(0)*(SM+4.-VM) FLO0960
 D2(MT)=D2(MT)+V*EXP(VNM*(VNM-2.*SM)-V*(V-2.*SM))/VNM FLO0970
 IF(D2(MT).LT.1.) GO TO 340 FLO0980
 D2(MT)=D2(MT)-1. FLO0990
 VN=V FLO1000
 350 V=8.*RAND(0)-4. FLO1010
 D3(MT)=D3(MT)+EXP(-V*V) FLO1020
 IF(D3(MT).LT.1.) GO TO 350 FLO1030
 D3(MT)=D3(MT)-1. FLO1040
 VT1=SN(MT)+V FLO1050
 360 V=8.*RAND(0)-4. FLO1060
 D4(MT)=D4(MT)+EXP(-V*V) FLO1070
 IF(D4(MT).LT.1.) GO TO 360 FLO1080
 D4(MT)=D4(MT)-1. FLO1090
 VT2=ST(MT)*B+V FLO1100
 PAU(MT,NMX)=VT1/BTA(MT) FLO1110
 PAV(MT,NMX)=(VN*A+VT2*B)/BTA(MT) FLO1120
 PAW(MT,NMX)=(-VN*B+VT2*A)/BTA(MT)
 LCOL(MT,NMX)=0
 ER(MT,NMX)=0.0

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225 IF (CHI(MT).LE.-1.) GO TO 365
      X=9.*RAND(0)
      IF (X.EQ.0.0) GO TO 225
      XT=X**CHI(MT)*EXP(-X)
      IF (XT.GE.CMG(MT)) GO TO 226
      CNG(MT)=CNG(MT)+XT
      IF (CNG(MT).LT.CMG(MT)) GO TO 225
      CNG(MT)=CNG(MT)-CMG(MT)
226 ER(MT,NMX)=X
365 LPP(MT,NMX)=LWF(6)
370 CONTINUE
      RETURN
380 LARGE=2
      RETURN
      END

```

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PLO1130
PLO1140
PLO1150
PLO1160
PLO1170
PLO1180

SUBROUTINE COLIDE(CN,CM,WTM,DB,DBA,NB,NCOL,LCOL,PAU,PAV,PAW,ER,T,
1LM,MT,I2,I3,NUMCEL,ETA,PHI,CHI,CN8,NP,LPP,LKW,NBP,NBM)

COL0030

```

      INTEGER TIME
      INTEGER*2 LM(NP,1), LCOL(NP,1), LPP(NP,1), LKW(1)
      INTEGER*2 NBM, NB, NBF, NUMCEL
      DIMENSION CN(3,3,1), CM(3,3,1), WTM(1), DB(NP,1), DBA(NP,1), NB(NP,1)
      DIMENSION NCOL(3,1), T(NP,NP,1), NUMCEL(1), ETA(3,1), PHI(3,1), CHI(1)
      DIMENSION PAU(NP,1), PAV(NP,1), PAW(NP,1), ER(NP,1), CN8(3,1), WA(2)
      DIMENSION NBP(NP,1), NBM(NP,1)

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COL0080
COL0090

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      COMMON /FORTH/NBX
      COMMON /FIFTH/ND,TIME,DTM
1 FORMAT(' TIME = ',F9.4,' COLL. TIMES = ',2F9.4,' NUMBERS = ',2I5/)
2 FORMAT(' COLIDE REACHED LINE 160 IN BOX NUMBER =', I5,' AT CPU TIM
1E =',F9.4/' VR= ',E12.4,' REL VEL G = ',3E12.4,' EI = ',E12.4)
3 FORMAT(' COLIDE REACHED LINE 165 IN BOX NUMBER =', I5,' AT CPU TIM
1E =',F9.4/' VR= ',E12.4,' REL VEL G = ',3E12.4,' EI = ',E12.4)

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COL0100

THE PURPOSE OF THIS SUBROUTINE IS TO ADVANCE THE ELAPSED TIMES IN
CELLS BY AN AMOUNT APPROXIMATELY EQUAL TO THE PRE-SELECTED COLLI
TIME. THERE ARE FOUR TIMES FOR EACH CELL, SAVED IN AN ARRAY CALLE
'T', CORRESPONDING TO THE FOUR TYPES OF MOLECULAR COLLISIONS WHIC
CAN OCCUR. TO ADVANCE THE VARIOUS TIMES, AN APPROPRIATE NUMBER OF
THE CORRESPONDING MOLECULAR COLLISIONS IS COMPUTED. THE ACTUAL
MOLECULES TO COLLIDE ARE SELECTED AT RANDOM, AND THEIR VELOCITY V
DIRECTIONS AFTER COLLISION ARE SELECTED AT RANDOM. -----
COL0180
COL0190

```

AIME=DTM*TIME
DO 240 MTA=1,MT
DO 230 MTB=1,MTA
D = WTM(MTA) + WTM(MTB)
WA(MTA)=WTM(MTA)/D
WA(MTB)=WTM(MTB)/D
RM=WTM(MTA)*WTM(MTB)/D
CHT=CHI(MTA)+CHI(MTB)+2.0
PHT=PHI(MTA,MTB)
ETT=ETA(MTA,MTB)
DO 220 M=1,NBX
N=NUMCEL(M)
IF (N.LE.0) GO TO 220

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AKW=LKW(N)
IF(T(MTA,MTB,N).LT.AIME) GO TO 100
IF(T(MTB,MTA,N).GE.AIME) GO TO 220
100 NA=NB(MTA,N)*NB(MTB,N)
IF(MTA.EQ.MTB) NA=(NA-NB(MTA,N))/2.
IF(NA.LT.1) GO TO 220
KS=0
120 KC=0
CPUT=ELTIME(0)
KS=KS+1
IF(KS.GT.NA) GO TO 220
130 KC=KC+1
IF(KC.GT.NA) GO TO 220
135 I=NB(MTA,N)*RAND(0)+1+NBM(MTA,N)
IF(I.GT.NBM(MTA,N+1)) I=NBM(MTA,N+1)
J=LM(MTA,I)
CR=LPP(MTA,J)/AKW
IF(CR.GT.0.99) GO TO 140
IF(RAND(0).GT.CR) GO TO 135
140 K=NB(MTB,N)*RAND(0)+1+NBM(MTB,N)
IF(K.GT.NBM(MTB,N+1)) K=NBM(MTB,N+1)
IF(MTA.EQ.MTB.AND.I.EQ.K) GO TO 140
L=LM(MTB,K)
CR=LPP(MTB,L)/AKW
IF(CR.GT.0.99) GO TO 145
IF(RAND(0).GT.CR) GO TO 140
145 CONTINUE
GM1=WA(MTA)*PAU(MTA,J)+WA(MTB)*PAU(MTB,L)
GM2=WA(MTA)*PAV(MTA,J)+WA(MTB)*PAV(MTB,L)
GM3=WA(MTA)*PAW(MTA,J)+WA(MTB)*PAW(MTB,L)
G1=PAU(MTA,J)-PAU(MTB,L)
G2=PAV(MTA,J)-PAV(MTB,L)
G3=PAW(MTA,J)-PAW(MTB,L)
GS=G1**2+G2**2+G3**2
IF(GS.LT.1.0E-8) GO TO 130
ET=RM*GS
EI=ER(MTA,J)+ER(MTB,L)
VR=GS**(-5-ETT/2.)
IF(VR.GE.CM(MTA,MTB,1)) GO TO 160
CN(MTA,MTB,1)=CN(MTA,MTB,1)+VR
IF(CN(MTA,MTB,1).LT.CM(MTA,MTB,1)) GO TO 130
CN(MTA,MTB,1)=CN(MTA,MTB,1)-CM(MTA,MTB,1)
160 CONTINUE
CPUT=ELTIME(0)
IF(M.EQ.1196) WRITE(6,2) M,CPUT,VR,G1,G2,G3,EI
CALL IMPACT(RM,G1,G2,G3,ET,EI,PHT,CHT,ETT,CM(MTA,MTB,2),CN(MTA,MTB,1,2))
165 CONTINUE
CPUT=ELTIME(0)
IF(M.EQ.1196) WRITE(6,3) M,CPUT,VR,G1,G2,G3,EI
IF(PHT.EQ.0.) GO TO 175
X1=0.0
IF(CHI(MTA).EQ.-1.) GO TO 175
X1=1.0
IF(CHI(MTB).EQ.-1.) GO TO 175

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170 X1=RAND(0)
  IF ((CHI(MTA).EQ.0.).AND.(CHI(MTB).EQ.0.)) GO TO 175
  XT=X1**CHI(MTA)*(1.-X1)**CHI(MTB)
  IF (XT.GT.CM(MTA,MTB,3)) GO TO 175
  CN(MTA,MTB,3)=CN(MTA,MTB,3)+XT
  IF (CN(MTA,MTB,3).LT.CM(MTA,MTB,3)) GO TO 170
  CN(MTA,MTB,3)=CN(MTA,MTB,3)-CM(MTA,MTB,3)
175 CONTINUE
  C=DBA(MTA,N)
  D=DBA(MTB,N)
  IF (C.EQ.0.0) C=DB(MTA,N)
  IF (D.EQ.0.0) D=DB(MTB,N)
  IF (T(MTA,MTB,N).GE.AIME) GO TO 180
  PAU(MTA,J)=GM1+WA(MTB)*G1
  PAV(MTA,J)=GM2+WA(MTB)*G2
  PAW(MTA,J)=GM3+WA(MTB)*G3
  IF (PHT.GT.0.) ER(MTA,J)=X1*EI
  LCOL(MTA,J)=1
  NCOL(MTA,MTB)=NCOL(MTA,MTB)+1
  T(MTA,MTB,N)=T(MTA,MTB,N)+CN8(MTA,MTB)*LPP(MTA,J)/NBF(MTA,N)/D/VR
  IF (MTA.EQ.MTB) GO TO 190
180 IF (T(MTB,MTA,N).GE.AIME) GO TO 210
190 PAU(MTB,L)=GM1-WA(MTA)*G1
  PAV(MTB,L)=GM2-WA(MTA)*G2
  PAW(MTB,L)=GM3-WA(MTA)*G3
  IF (PHT.GT.0.) ER(MTB,L)=(1.-X1)*EI
  LCOL(MTB,L)=1
  NCOL(MTB,MTA)=NCOL(MTB,MTA)+1
  T(MTB,MTA,N)=T(MTB,MTA,N)+CN8(MTB,MTA)*LPP(MTB,L)/NBF(MTB,N)/C/VR
210 CONTINUE
  IF (N.EQ.1196) WRITE(6,1) AIME,T(MTA,MTB,N),T(MTB,MTA,N),NBF(MTA,N)
  1,NBF(MTB,N)
  IF (T(MTA,MTB,N).LT.AIME.OR.T(MTB,MTA,N).LT.AIME) GO TO 120
220 CONTINUE
230 CONTINUE
240 CONTINUE
  RETURN
  END

  SUBROUTINE MOVE(KSWCH,AKN,MJ,NS,NWEDG,THETAZ,XSTART,I2,I3,I4,I5,DEMOV0010
  1LANG,NWEDGE,BTA,C2,C3,DFA,PL,HTI,HTR,JNT,KNM,NM,XCB,XLIM,MS,IWS,NTMOV0020
  2CF,NTCV,FV,CTI,CTR,CNI,CNR,ALPHA,SIGMA,COEFF,HTS,HTSI,NTS,NTSP,UTLMOV0030
  3,UTT,VTS,PAU,PAV,PAW,PAX,PAZ,LPF,LCOL,TB,IP,ER,CHI,CNG,CMG,I,
  4UTLI,UTTI,VTSI)
  INTEGER*2 LPF(I,1),LCOL(I,1)
  INTEGER SWTCH,TIME
  LOGICAL DUMP
  REAL LAM,MU,NU
  DIMENSION DELANG(1),NWEDGE(1),BTA(1),C2(1),C3(1),PL(1),HTI(1)
  DIMENSION HTR(1),TB(1),XCB(1),ALPHA(3,1),SIGMA(3,1),COEFF(4,1)
  DIMENSION PAU(I,1),PAV(I,1),PAW(I,1),CTI(3,1),CTR(3,1)
  DIMENSION CNI(3,1),CNR(3,1),DFA(1),JNT(1),XLIM(1),KNM(1),NM(1)
  DIMENSION HTS(3,I2,I3),HTSI(3,I2,I3),NTS(3,I2,I3),NTSP(3,I2,I3)
  DIMENSION UTLI(3,I2,I3),UTTI(3,I2,I3),VTSI(3,I2,I3)
  DIMENSION UTL(3,I2,I3),UTT(3,I2,I3),VTS(3,I2,I3),NTCP(3,I4)

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  COL0810
  COL0920
  MOV0060
  MOV0070
  MOV0080
  MOV0090

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DIMENSION PAX(1,1), PAY(1,1), PAZ(1,1), IWS(1), MS(1) MOV0170
 DIMENSION NTCV(3,14,2,15,3), FV(3,14,2,15,3) MOV0180
 DIMENSION ER(1,1), CNG(1), CHG(1), CHI(1) MOV0190
 COMMON /THIRD/PI, NREG MOV0200
 COMMON /FORTH/NBX, RM, XR, DUMP MOV0210
 COMMON /FIFTH/ND, TIME, DTM MOV0220
 COMMON /SVNTH/LAM, MU, NU, MT, N, J, XI, YI, ZI, TUSE MOV0230
 NAMELIST/CHECK/TIME, X, Y, Z, DX, DY, DZ, TLEFT, RADS, RMS, XR MOV0240

 THE PURPOSE OF THIS SUBROUTINE IS TO ADVANCE THE SPATIAL POSITION MOV0250
 OF ALL THE MOLECULES BY AN AMOUNT APPROPRIATE TO THEIR CURRENT VEN MOV0260
 LOCITIES AND THE PRE-SELECTED COLLISION TIME. MOV0270

 FORMAT(27H SOMETHING IS WRONG IN MOVE/3E20.7,4I17,E20.7) MOV0280
 NAREA=NREG+3 MOV0290
 RMS=RM**2 MOV0300
 DO 150 MT=1,IP MOV0310
 N=KNN(MT) MOV0320
 N=N+1 MOV0330
 TLEFT=DTM MOV0340
 IF (KSWCH.EQ.1) TLEFT=TLEFT*RAND(0) MOV0350
 IF (N.GT.NM(MT)) GO TO 150 MOV0360
 LAM=PAU(MT,N) MOV0370
 IF (LAM.EQ.0.) LAM=.0000001 MOV0380
 MU=PAV(MT,N) MOV0390
 NU=PAW(MT,N) MOV0400
 XI=PAX(MT,N) MOV0410
 YI=PAY(MT,N) MOV0420
 ZI=PAZ(MT,N) MOV0430
 DX=TLEFT*LAM MOV0440
 DY=TLEFT*MU MOV0450
 DZ=TLEFT*NU MOV0460
 X=XI+DX MOV0470
 Y=YI+DY MOV0480
 Z=ZI+DZ MOV0490
 RADS=Y**2+Z**2 MOV0500
 IF ((RADS.GT.2.*RMS).OR.(ABS(X).GT.2.*XR)) WRITE(6,CHECK) MOV0510
 IF (RADS.GT.RMS) GO TO 100 MOV0520
 RAD=SQRT(RMS) MOV0530
 KEY=ABS(LAM)/LAM+.5 MOV0540
 DO 60 L=1,NAREA MOV0550
 IF (XI-XLIM(L)) 65,55,60 MOV0560
 J=L+KEY-2 MOV0570
 GO TO 70 MOV0580
 CONTINUE MOV0590
 WRITE(6,2) DTM, XI, LAM, N, KEY, NAREA, L, XLIM(L) MOV0600
 IF (DUMP) CALL ABEND(16) MOV0610
 STOP MOV0620
 J=L-2 MOV0630
 K=J+KEY+1 MOV0640
 IF ((K.EQ.0).OR.(K.EQ.NAREA+1)) GO TO 100 MOV0650
 TUSE=(XLIM(K)-XI)/LAM MOV0660
 XTEMP=XLIM(K) MOV0670
 IF (TUSE.LE.TLEFT) GO TO 75 MOV0680
 TUSE=TLEFT MOV0690

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XTEMP=XI+TUSE*LAM          MOV0650
75 IF((J.EQ.0).OR.(J.EQ.NREG+1)) GO TO 85      MOV0660
CALL INTERS(AFN,MJ,NS,NWEDG,SWTCH,THETAZ,XSTART,I2,I3,I4,I5,DELANGMOV0670
1,NWEDGE,BTA,C2,C3,DFA,FL,HTI,HTR,JNT,TB,XCB,CTI,CTR,CNI,CNR,ALPHA,MOV0680
2SIGMA,COEFF,HTS,HTSI,NTS,NTSP,UTL,UTT,VTS,MS,IWS,NTCF,NTCV,FV,PAU,MOV0690
3PAV,PAW,LPP,LCOL,IP,ER,CHI,CNG,CMG,I,UTLI,UTTI,VTSI)
IF(SWTCH.EQ.1) GO TO 90          MOV0710
85 XI=XTEMP                  MOV0720
YI=YI+TUSE*NU                MOV0730
ZI=ZI+TUSE*NU                MOV0740
90 PAX(MT,N)=XI              MOV0750
PAY(MT,N)=YI                  MOV0760
PAZ(MT,N)=ABS(ZI)            MOV0770
PAW(MT,N)=ABS(ZI)/ZI*PAW(MT,N)  MOV0780
TLEFT=TLEFT-TUSE              MOV0790
IF(TLEFT.GT.0.) GO TO 15      MOV0800
GO TO 10                      MOV0810
100 NZ=NM(MT)                 ORIGINAL PAGE IS
PAX(MT,N)=PAX(MT,NZ)          OF POOR QUALITY
PAY(MT,N)=PAY(MT,NZ)          MOV0820
PAZ(MT,N)=PAZ(MT,NZ)          MOV0830
PAU(MT,N)=PAU(MT,NZ)          MOV0840
PAV(MT,N)=PAV(MT,NZ)          MOV0850
PAW(MT,N)=PAW(MT,NZ)          MOV0860
ER(MT,N)=ER(MT,NZ)            MOV0870
LPP(MT,N)=LPP(MT,NZ)          MOV0880
LCOL(MT,N)=LCOL(MT,NZ)        MOV0890
N=N-1                         MOV0900
NM(MT)=NM(MT)-1              MOV0910
GO TO 10                      MOV0920
150 CONTINUE                   MOV0930
RETURN                         MOV0940
END                            MOV0950
                                MOV0960

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SUBROUTINE ACCUM(I2,I3,FNB,NB,PAU,PAV,PAW,ER,TMP,TRP,XV,YV,ZV,LM,
1IP,I,LPP,NBF,NBM)
INTEGER*2 LM(I,1),LPP(I,1),NBM(I,1)
INTEGER*2 NB,NBF
DIMENSION FNB(1),NB(I,1),PAU(I,1),PAV(I,1),PAW(I,1),TMP(I,1)
DIMENSION XV(I,1),YV(I,1),ZV(I,1),ER(I,1),TRP(I,1),NBF(I,1)
COMMON /FORTH/NBX
----- ACUM050
----- ACUM060
----- ACUM070
----- ACUM080
----- ACUM090
----- ACUM100
----- ACUM110
----- ACUM120
----- ACUM130
----- ACUM140
----- ACUM160
----- ACUM170
----- ACUM180
----- ACUM190

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THE PURPOSE OF THIS SUBROUTINE IS TO ACCUMULATE TEMPERATURES, ACUM070
VELOCITIES, AND DENSITIES IN VARIOUS ARRAYS FOR DETERMINING THE ACUM080
AVERAGE FLOW FIELD PROPERTIES AFTER STEADY-STATE HAS BEEN REACHEDACUM090
----- ACUM100
N=0                         ACUM110
DO 180 K=1,NBX              ACUM120
IF(FNB(K).LE.0.) GO TO 180  ACUM130
N=N+1                        ACUM140
DO 110 MT=1,IP              ACUM160
XV(MT,N)=0.0                  ACUM170
YV(MT,N)=0.0                  ACUM180
ZV(MT,N)=0.0                  ACUM190
TMP(MT,N)=0.

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TRP(MT,N)=0.0
TTX=0.                                ACUM200
TTY=0.                                ACUM210
TTZ=0.                                ACUM220
TTR=0.0
M=NBM(MT,N)
IF(M.LT.1) GO TO 110
U=0.                                ACUM230
V=0.                                ACUM240
W=0.                                ACUM250
DO 100 L=1,M
NA=NBM(MT,N)+L
J=LM(MT,NA)
AKW=LPP(MT,J)
PU=PAU(MT,J)                         ACUM260
PV=PAV(MT,J)                         ACUM270
PW=PAW(MT,J)                         ACUM280
U=U+PU*AKW
V=V+PV*AKW
W=W+PW*AKW
TTR=TTR+ER(MT,J)*AKW
TTX=TTX+PU*PU*AKW
TTY=TTY+PV*PV*AKW
20 TTZ=TTZ+PW*PW*AKW
M=NBF(MT,N)
XV(MT,N)=U/M                         ACUM300
YV(MT,N)=V/M                         ACUM310
ZV(MT,N)=W/M                         ACUM320
TMP(MT,N)=(TTX+TTY+TTZ)/M
TRP(MT,N)=TTR/M
10 CONTINUE
30 CONTINUE
RETURN
END

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SUBROUTINE AVRGE(FNB,DB,DBA,NB,NBT,XV,YV,ZV,XVA,YVA,ZVA,TMP,TMPA,
1TRP,TRPA,IP,I,NBF,NBS)
INTEGER*2 NB,NBT,NBF,NBS
DIMENSION FNB(1),DB(I,1),DBA(I,1),NB(I,1),NBT(I,1),TMP(I,1)
6 DIMENSION TMPA(I,1),XV(I,1),XVA(I,1),YV(I,1),YVA(I,1),ZV(I,1)
DIMENSION ZVA(I,1),TRP(I,1),TRPA(I,1),NBF(I,1),NBS(I,1)
COMMON /FORTH/NBX

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AVG0050

AVG0060

THE PURPOSE OF THIS SUBROUTINE IS TO COMPUTE THE AVERAGE FLOW AVG0070
FIELD PROPERTIES. AVG0080

AVG0090

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N=0
DO 110 M=1,NBX
IF(FNB(M).LE.0.) GO TO 110
N=N+1
DO 100 MT=1,IP
A=NBT(MT,N)
B=NBF(MT,N)
C=A+B
NBT(MT,N)=C

```

AVG0100

AVG0110

AVG0120

AVG0130

AVG0150

AVG0170

AVG0180

NBS(MT,N)=NBS(MT,N)+NB(3T,N)	
IF(C.LT.1.) GO TO 100	AVG0190
DBA(MT,N)=(DBA(MT,N)*A+DB(MT,N)*B)/C	AVG0200
XVA(MT,N)=(XVA(MT,N)*A+XV(MT,N)*B)/C	AVG0210
YVA(MT,N)=(YVA(MT,N)*A+YV(MT,N)*B)/C	AVG0220
ZVA(MT,N)=(ZVA(MT,N)*A+ZV(MT,N)*B)/C	AVG0230
TMPA(MT,N)=(TMPA(MT,N)*A+TMP(MT,N)*B)/C	AVG0240
TRPA(MT,N)=(TRPA(MT,N)*A+TRP(MT,N)*B)/C	AVG0250
100 CONTINUE	AVG0260
110 CONTINUE	AVG0270
RETURN	AVG0280
END	
ORIGINAL PAGE IS OF POOR QUALITY	
SUBROUTINE DRAG(AKN,MJ,NS,NWEDGE,THETAZ,XSTART,I2,I3,I4,I5,DELANG, 1NWEDGE,BTA,C2,C3,DFA,FL,HTI,HTR,TB,XCB,CTI,CTR,CNI,CNR,ALPHA,SIGMA 2,COEFF,HTS,HTSI,NTS,NTSP,UTL,UTT,VTS,MS,IWS,NTCP,NTCV,FV,PAU,PAV,P 3AW,LPP,LCOL,IP,ER,CHI,CNG,CMG,I,UTLI,UTTI,VTSI)	
INTEGER*2 LPP,LCOL	DRG0040
INTEGER TIME,TST	DRG0050
REAL LAM,MU,NU,JAY,KAY	DRG0060
DIMENSION DELANG(1),NWEDGE(1),BTA(1),C2(1),C3(1),FL(1),HTI(1)	DRG0070
DIMENSION HTR(1),TB(1),XCB(1),ALPHA(3,1),SIGMA(3,1),COEFF(4,1)	
DIMENSION LPP(I,1),PAU(I,1),PAV(I,1),PAW(I,1),CTI(3,1),CTR(3,1)	
DIMENSION HTS(3,I2,I3),HTSI(3,I2,I3),NTS(3,I2,I3),NTSP(3,I2,I3)	
DIMENSION UTL(3,I2,I3),UTT(3,I2,I3),VTS(3,I2,I3),LCOL(I,1)	
DIMENSION UTLI(3,I2,I3),UTTI(3,I2,I3),VTSI(3,I2,I3)	
DIMENSION CNI(3,1),CNR(3,1),DFA(1),IWS(1),MS(1),NTCP(3,I4)	
DIMENSION NTCV(3,I4,2,I5,3),FV(3,I4,2,I5,3)	
DIMENSION ER(I,1),CNG(1),CMG(1),CHI(1)	
COMMON /THIRD/PI	DRG0140
COMMON /FIFTH/ND,TIME,DTM,TI,ITS,ITP,TST	DRG0150
COMMON /SVNTH/LAM,MU,NU,MT,N,J,XCL,YCL,ZCL	DRG0160
-----	DRG0170
THE PURPOSE OF THIS SUBROUTINE IS TO ACCUMULATE THE DRAG AND HEAT DRG0180 TRANSFER INCREMENTS ON THE BODY CONTRIBUTED BY EACH MOLECULE WHICH DRG0190 COLLIDES WITH THE BODY. IN ADDITION, EACH MOLECULE WHICH DRG0200 WITH THE BODY IS ASSIGNED AN APPROPRIATE NEW VELOCITY (OF REFLECT DRG0210 WHICH IS USED TO CONTINUE ITS SPATIAL TRANSLATION (IN SUBROUTINE DRG0220	
-----	DRG0230
CALL NORMAL(EYE,JAY,KAY,ONE,COEFF)	DRG0240
AKW=LPP(MT,N)	DRG0250
RAD=SQRT(YCL*YCL+ZCL*ZCL)	DRG0260
ARG=YCL/RAD	DRG0270
TANG=180.*(1.-ARCCOS(ARG)/PI)	DRG0280
IWDG=TANG/DELANG(1)+1.	DRG0290
IF((IWDG.GT.NWEDGE(1)).AND.(DELANG(2).NE.0.)) IWDG=(TANG-THETAZ)	
1/DELANG(2)+NWEDGE(1)+1	
IF(IWDG.LT.1) IWDG=1	
IF(IWDG.GT.NWEDG) IWDG=NWEDG	DRG0310
D=(LAM*LAM+MU*MU+NU*NU)*AKW	DRG0320
G=ER(MT,N)*AKW	
H=G	
DO 100 M=1,ND	DRG0330
IF(XCL.LT.XCB(M)) GO TO 110	DRG0340
100 CONTINUE	DRG0350

110	UI=LAM	DRG0360
	WI=(NU*JAY-MU*KAY)/ONE	DRG0370
	VID=LAM*EYE+MU*JAY+NU*KAY	DRG0380
	UID=LAM*ONE-EYE*(MU*JAY+NU*KAY)/ONE	DRG0390
	E=RAND(0)	DRG0400
	I=(E.LT.SIGMA(MT,M)) GO TO 115	DRG0410
	VRD=-VID	DRG0420
	URD=UID	DRG0430
	WR=WI	DRG0440
	GO TO 125	DRG0450
115	V=4.*RAND(0)	DRG0460
	C2(MT)=C2(MT)+.544331*V*V*V*EXP(1.5-V*V)	DRG0470
	IF(C2(MT).LT.1.) GO TO 115	DRG0480
	C2(MT)=C2(MT)-1.	DRG0490
	IF(NTSP(MT,M,IWDG).NE.0) GO TO 117	DRG0500
	ATR=ALPHA(MT,M)*TB(M)/SIGMA(MT,M)	DRG0510
	GO TO 118	DRG0520
117	ATE=ALPHA(MT,M)*TB(M)/SIGMA(MT,M)+(1.-ALPHA(MT,M)/SIGMA(MT,M))*HTSDRG0530	
	I(MT,M,IWDG)/NTSP(MT,M,IWDG)/(3.+CHI(MT))	
118	ABR=SQRT(ATR)	DRG0550
	V=V*ABR/BTA(MT)	DRG0560
120	A=RAND(0)	DRG0570
	C3(MT)=C3(MT)+A	DRG0580
	IF(C3(MT).LT.1.) GO TO 120	DRG0590
	C3(MT)=C3(MT)-1.	DRG0600
	B=SQRT(1.-A*A)	DRG0610
	C=2.*PI*RAND(0)	DRG0620
	VRD=V*A	DRG0630
	URD=V*B*COS(C)	DRG0640
	WR=V*B*SIN(C)	DRG0650
	IF(CHI(MT).EQ.-1.) GO TO 125	
122	X=9.*RAND(0)	
	IF(X.EQ.0.0) GO TO 122	
	XTEMP=1.0	
	IF(CHI(MT).NE.0.0) XTEMP=1.**CHI(MT)	
	CNG(MT)=CNG(MT)+XTEMP*EXP(-X)	
	IF(CNG(MT).LT.CMG(MT)) GO TO 122	
124	CONTINUE	
	CNG(MT)=CNG(MT)-CMG(MT)	
	IF(CNG(MT).GE.CMG(MT)) GO TO 124	
	ER(MT,N)=X*ATR	
	H=ER(MT,N)*AKW	
125	UR=EYE*VRD+ONE*URD	DRG0660
	PAU(MT,N)=UR	DRG0670
	PAV(MT,N)=JAY*VRD-(KAY*WR+EYE*JAY*URD)/ONE	DRG0680
	PAW(MT,N)=KAY*VRD+(JAY*WR-EYE*KAY*URD)/ONE	DRG0690
	IF(TIME.LE.TST) RETURN	
	IF(TI.GT.0.) GO TO 130	DRG0710
	TI=TST*DTM	
130	XMZ=(XCL-XSTART)*AKN	DRG0730
	YMZ=RAD*AKN	DRG0740
	B=(URD*URD+VRD*VRD+WR*WR)*AKW	DRG0750
	UTI=UID*UID+WI*WI	DRG0760
	UYI=-(UID*EYE*JAY+WI*KAY)/ONE	DRG0770
	UYR=-(URD*EYE*JAY+WR*KAY)/ONE	DRG0780

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PL (MT) = PL (MT) + AKW*DFA (MT)	DRG0790	
HTI (MT) = HTI (MT) + D+G		
HTR (MT) = HTR (MT) - B-H		
CTI (MT, 1) = CTI (MT, 1) + UID*ONE*AKW	DRG0820	
CTI (MT, 2) = CTI (MT, 2) + UYI*AKW	DRG0830	
CTI (MT, 3) = CTI (MT, 3) + (XMZ*UYI-YMZ*UID*ONE)*AKW	DRG0840	
CNI (MT, 1) = CNI (MT, 1) + VID*EYE*AKW	DRG0850	
CNI (MT, 2) = CNI (MT, 2) + VID*JAY*AKW	DRG0860	
CNI (MT, 3) = CNI (MT, 3) + (XMZ*JAY-YMZ*EYE)*VID*AKW	DRG0870	
CTR (MT, 1) = CTR (MT, 1) - URD*ONE*AKW	DRG0880	
CTR (MT, 2) = CTR (MT, 2) - UYR*AKW	DRG0890	
CTR (MT, 3) = CTR (MT, 3) - (XMZ*UYR-YMZ*URD*ONE)*AKW	DRG0900	
CNR (MT, 1) = CNR (MT, 1) - VRD*EYE*AKW	DRG0910	
CNR (MT, 2) = CNR (MT, 2) - VRD*JAY*AKW	DRG0920	
CNR (MT, 3) = CNR (MT, 3) - (XMZ*JAY-YMZ*EYE)*VRD*AKW	DRG0930	
NTS (MT, M, IWDG) = NTS (MT, M, IWDG) + 1	DRG0940	
NTSF (MT, M, IWDG) = NTSF (MT, M, IWDG) + AKW	DRG0950	
UTLI (MT, M, IWDG) = UTLI (MT, M, IWDG) + UID*AKW		
UTL (MT, M, IWDG) = UTL (MT, M, IWDG) + (UID-URD)*AKW	DRG0960	
UTTI (MT, M, IWDG) = UTTI (MT, M, IWDG) + WI*AKW		
UTT (MT, M, IWDG) = UTT (MT, M, IWDG) + (WI-WR)*AKW	DRG0970	
VTSI (MT, M, IWDG) = VTSI (MT, M, IWDG) - VID*AKW		
VTS (MT, M, IWDG) = VTS (MT, M, IWDG) + (VRD-VID)*AKW	DRG0980	
HTSI (MT, M, IWDG) = HTSI (MT, M, IWDG) + D+G		
HTS (MT, M, IWDG) = HTS (MT, M, IWDG) + D-B+G-H		
IF (NS.EQ.0) RETURN	DRG1010	
NC = (2*LCOL (MT, N)) / (1+LCOL (MT, N)) + 1	DRG1020	
DO 160 L=1, NS		
MTEST=MS (L)		
IWT=IWS (L)		
IF ((M.NE.MTEST).OR.(IWDG.NE.IWT)) GO TO 160	DRG1060	
IP (NC.EQ.1) NTCF (MT, L) = NTCF (MT, L) + AKW		
DO 145 JJ=1, MJ	DRG1080	
IP (ABS(VID).GT.FV (MT, L, NC, JJ, 1)) GO TO 135		
NTCV (MT, L, NC, JJ, 1) = NTCV (MT, L, NC, JJ, 1) + AKW		
135 IP (UID.GT.FV (MT, L, NC, JJ, 2)) GO TO 140		
NTCV (MT, L, NC, JJ, 2) = NTCV (MT, L, NC, JJ, 2) + AKW		
140 IP (WI*ABS(ZCL)/ZCL.GT.FV (MT, L, NC, JJ, 3)) GO TO 145		
NTCV (MT, L, NC, JJ, 3) = NTCV (MT, L, NC, JJ, 3) + AKW		
145 CONTINUE	DRG1150	
160 CONTINUE	DRG1160	
RETURN	DRG1170	
END	DRG1180	
SUBROUTINE INTERS (AKN, MJ, NS, NWEDGE, SWTCH, THETAZ, XSTART, I2, I3, I4, I5, INT0010		
1DELANG, NWEDGE, BTA, C2, C3, DFA, PL, HTI, HTR, JNT, TB, XCB, CTI, CTR, CNI, CNR, INT0020		
2ALPHA, SIGMA, COEFF, HTS, HTSI, NTS, NTSF, UTL, UTT, VTS, MS, IWS, NTCF, NTCV, PINT0030		
3V, PAU, PAV, PAW, LPP, LCOL, IP, ER, CHI, CNG, CMG, I, UTLI, UTTI, VTSI)		
INTEGER*2 LPP, LCOL		INT0050
INTEGER SWTCH		INT0060
REAL LAM, MU, NU, LINEAR		INT0070
DIMENSION DELANG (1), NWEDGE (1), BTA (1), C2 (1), C3 (1), PL (1), HTI (1)		INT0080
DIMENSION HTR (1), TB (1), XCB (1), ALPHA (3, 1), SIGMA (3, 1), COEFF (4, 1)		
DIMENSION LPP (1, 1), PAU (1, 1), PAV (1, 1), PAW (1, 1), CTI (3, 1), CTR (3, 1)		
DIMENSION HTS (3, I2, I3), HTSI (3, I2, I3), NTS (3, I2, I3), NTSF (3, I2, I3)		

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DIMENSION UTL(3,I2,I3), UTT(3,I2,I3), VTS(3,I2,I3), LCOL(I,1)
DIMENSION UTLI(3,I2,I3), UTTI(3,I2,I3), VTSI(3,I2,I3)
DIMENSION CNI(3,1), CNR(3,1), DPA(1), IWS(1), MS(1), NTCP(3,I4)
DIMENSION JNT(1), NTCV(3,I4,2,I5,3), FV(3,I4,2,I5,3)
DIMENSION ER(I,1), CNG(1), CMG(1), CHI(1)
COMMON /SVNTH/LAM,MU,NU,MT,N,J,XI,YI,ZI,TUSE           INTO 150
-----INTO 160
THE PURPOSE OF THIS SUBROUTINE IS TO DETERMINE FOR EACH MOLECULAR INTO 170
TRAJECTORY IF THERE IS AN INTERSECTION OF THE TRAJECTORY WITH THE INTO 180
BODY SURFACE.                                         INTO 190
-----INTO 200
SWTCH=0                                         INTO 210
A=COEFF(1,J)                                     INTO 220
B=COEFF(2,J)                                     INTO 230
C=COEFF(3,J)                                     INTO 240
D=COEFF(4,J)                                     INTO 250
ONE=A*LAM                                         ORIGINAL PAGE IS
TWO=B*MU                                         OF POOR QUALITY
TRE=B*NU                                         INTO 260
PUR=.5*C*LAM                                     INTO 270
SQUARE=ONE*LAM+TWO*MU+TRE*NU                      INTO 280
LINEAR=ONE*XI+TWO*YI+TRE*ZI+PUR                  INTO 290
CONST=(A*XI+C)*XI+B*(YI*YI+ZI*ZI)+D           INTO 300
IF(SQUARE.EQ.0.) GO TO 150                         INTO 310
DISCR=LINEAR*LINEAR-SQUARE*CONST                  INTO 320
IF(DISCR.LT.0.) RETURN                            INTO 330
SDISC = SQRT(DISCR)                             INTO 340
TYME = (-LINEAR-SDISC)/SQUARE                     INTO 350
GO TO 250                                         INTO 360
150 IF(LINEAR.EQ.0.) RETURN                         INTO 370
TYME=-.5*CONST/LINEAR                            INTO 380
250 IF(TYME.GT.TUSE) RETURN                         INTO 390
IF(TYME.LE.0.) RETURN                            INTO 400
JNT(MT)=JNT(MT)+1                                INTO 410
XI=XI+LAM*TYME                                     INTO 420
YI=YI+MU*TYME                                     INTO 430
ZI=ZI+NU*TYME                                     INTO 440
CALL DRAG(AKN,MJ,NS,NWEDG,THETAZ,XSTART,I2,I3,I4,I5,DELANG,NWEDGE,INTO 450
1BTA,C2,C3,DPA,PL,HTI,HTR,TB,XCB,CTI,CTR,CNI,CNR,ALPHA,SIGMA,COEFF,INTO 460
2HTS,HTSI,NTS,NTSP,UTL,UTT,VTS,MS,IWS,NTCP,NTCV,PV,PAU,PAV,PAW,LPP,INTO 470
3LCOL,IP,ER,CHI,CNG,CMG,I,UTLI,UTTI,VTSI)
LCOL(MT,N)=2                                         INTO 480
TUSE=TYME                                         INTO 490
SWTCH=1                                         INTO 500
RETURN                                         INTO 510
END                                             INTO 520
-----INTO 530
SUBROUTINE NORMAL(EYE,JAY,KAY,ONE,COEFF)          NORM010
REAL LAM,MU,NU,JAY,KAY
DIMENSION COEFF(4,1)                                NORM020
COMMON /SVNTH/LAM,MU,NU,MT,N,J,XCL,YCL,ZCL      NORM030
DFDX=2.*COEFF(1,J)*XCL+COEFF(3,J)                NORM040
DFDY=2.*COEFF(2,J)*YCL                            NORM050
DFDZ=2.*COEFF(2,J)*ZCL                            NORM060
DENOM=SQRT(DFDX*DFDX+DFDY*DFDY+DFDZ*DFDZ)      NORM070
                                              NORM080

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EYE=DPDX/DENOM          NORM090
JAY=DPDY/DENOM          NORM100
KAY=DPDZ/DENOM          NORM110
ONE=SQRT (JAY*JAY+KAY*KAY) ORIGINAL PAGE IS
RETURN                   OF POOR QUALITY
END                      NORM120
                          NORM130
                          NORM140

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FUNCTION ARCCOS(ARG)          ARCS010
COMMON /THIRD/PI             ARCS020
IF (ARG) 30,10,20             ARCS030
10 A=.5*PI                   ARCS040
GO TO 40                      ARCS050
20 A=ATAN (SQRT (1.-ARG*ARG)/ARG) ARCS060
GO TO 40                      ARCS070
30 A=PI+ATAN (SQRT (1.-ARG*ARG)/ARG) ARCS080
40 ARCCOS=A                   ARCS090
RETURN                         ARCS100
END                            ARCS110

```

```

FUNCTION EERRF(SS)           EERRF010
EERRF=ERFC(-SS)              EERRF050
RETURN                         EERRF060
END

```

```

SUBROUTINE PRINT1(DT,COSANG,SINANG,RMA,RNU,DRF,PCF,HTF,PL,HTI,HTR,
1CTI,CTR,CNI,CNR)
DIMENSION DD(3),WD(2,5),PP(4,4),QQ(4,4),RR(4,4),SS(4,4),TT(4,4)
DIMENSION UU(4,4),P1(4,4),Q1(4,4),R1(4,4),PA(4),PB(4),PC(4)
DIMENSION PL(1),HTI(1),HTR(1),CTI(3,1),CTR(3,1),CNI(3,1),CNR(3,1)
DIMENSION RMA(1),RNU(1)
DATA WD// 'X-PO','RCE ','Y-PO','RCE ','Z-MO','MENT','DRAG','  ','LPT10070
1IPT','  '/
----- PT10080
----- PT10090

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```

THE PURPOSE OF THIS SUBROUTINE IS TO PRINT OUT THE GROSS SURFACE PT10100
COEFFICIENTS OF THE BODY. PT10110
----- PT10120
----- PT10130
----- PT10140
----- PT10150
----- PT10160
----- PT10170
----- PT10180

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FORMATS

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1 FORMAT(//1X,50('*'),' GROSS SURFACE COEFFICIENTS ',50('*')// MOLEC
1ULAR WEIGHT',12X,F8.3,3(19X,F8.3)/25X,
2 'INC.      REF.      TOT.      INC.      REF.      TOT.      INC.
3REF.      TOT.      INC.      REF.      TOT.')
10 FORMAT(' NUMBER FLUX      ',4(F9.3,18X))
12 FORMAT(1X,2A4,2X,'SHEAR      ',4(3F8.3,3X))
14 FORMAT(11X,'PRESSURE      ',4(3F8.3,3X))
16 FORMAT(11X,'TOTAL      ',4(3F8.3,3X)/)
18 FORMAT(' HEAT TRANSFER',7X,4(3F8.3,3X)/)
----- PT10280
***** PT10290
----- PT10300
----- PT10310

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RMR=0.0
DO 50 MT=1,3
  DD (MT)=RMA (MT) *RNU (MT) *DRP/DT
* 50 RMR=RMR+RMA (MT) *RNU (MT)
  WRITE (6, 1) (RMA (MT), MT=1, 3), RMR
  PF=PL (1) *PCP/DT
  QP=PL (2) *PCP/DT
  RF=PL (3) *PCP/DT
  SF=PF+QP+RF
*   WRITE (6, 10) PF, QP, RF, SF
  DO 200 I=1,3
    PP (4, I)=0.0
    QQ (4, I)=0.0
    RR (4, I)=0.0
    SS (4, I)=0.0
    TT (4, I)=0.0
    UU (4, I)=0.0
    P1 (4, I)=0.0
    Q1 (4, I)=0.0
    R1 (4, I)=0.0
  DO 150 MT=1,3
    PP (MT, I)=CTI (MT, I) *DD (MT) /RMR
    QQ (MT, I)=CTR (MT, I) *DD (MT) /RMR
    SS (MT, I)=CNI (MT, I) *DD (MT) /RMR
    TT (MT, I)=CNR (MT, I) *DD (MT) /RMR
    P1 (MT, I)=PP (MT, I) +SS (MT, I)
    Q1 (MT, I)=QQ (MT, I) +TT (MT, I)
    RR (MT, I)=PP (MT, I) +QQ (MT, I)
    UU (MT, I)=SS (MT, I) +TT (MT, I)
    R1 (MT, I)=P1 (MT, I) +Q1 (MT, I)
    PP (4, I)=PP (4, I) +PP (MT, I)
    QQ (4, I)=QQ (4, I) +QQ (MT, I)
    RR (4, I)=RR (4, I) +RR (MT, I)
    SS (4, I)=SS (4, I) +SS (MT, I)
    TT (4, I)=TT (4, I) +TT (MT, I)
    UU (4, I)=UU (4, I) +UU (MT, I)
    P1 (4, I)=P1 (4, I) +P1 (MT, I)
    Q1 (4, I)=Q1 (4, I) +Q1 (MT, I)
    R1 (4, I)=R1 (4, I) +R1 (MT, I)
*150 CONTINUE
  WRITE (6, 12) (WD (J, I), J=1, 2), (PP (K, I), QQ (K, I), RR (K, I), K=1, 4)
  WRITE (6, 14) (SS (K, I), TT (K, I), UU (K, I), K=1, 4)
  WRITE (6, 16) (P1 (K, I), Q1 (K, I), R1 (K, I), K=1, 4)
200 CONTINUE
  AA=COSANG
  BB=SINANG
  DO 300 I=4,5
    DO 250 K=1,4
      PP (K, 4)=AA*PP (K, 1) +BB*PP (K, 2)
      QQ (K, 4)=AA*QQ (K, 1) +BB*QQ (K, 2)
      RR (K, 4)=AA*RR (K, 1) +BB*RR (K, 2)
      SS (K, 4)=AA*SS (K, 1) +BB*SS (K, 2)
      TT (K, 4)=AA*TT (K, 1) +BB*TT (K, 2)
      UU (K, 4)=AA*UU (K, 1) +BB*UU (K, 2)
      P1 (K, 4)=AA*P1 (K, 1) +BB*P1 (K, 2)

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PT10400

PT10460

PT10470

PT10480

PT10490

PT10630

PT10640

PT10650

PT10660

PT10680

PT10690

PT10700

PT10710

PT10720

PT10730

PT10740

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Q1(K,4)=AA*Q1(K,1)+BB*Q1(K,2) PT10750
250 R1(K,4)=AA*R1(K,1)+BB*R1(K,2) PT10760
  WRITE(6,12) (WD(J,I),J=1,2), (PP(K,4),QQ(K,4),RR(K,4),K=1,4)
  WRITE(6,14) (SS(K,4),TT(K,4),UU(K,4),K=1,4)
  WRITE(6,16) (P1(K,4),Q1(K,4),R1(K,4),K=1,4)
  AA=-SINANG PT10800
  BB=COSANG PT10810
300 CONTINUE PT10820
  HD=HTP/DT PT10830
  PA(4)=0.0
  PB(4)=0.0
  PC(4)=0.0
  DO 400 MT=1,3
    PA(MT)=HTI(MT)*RMA(MT)*RNU(MT)*HD/RMR
    PB(MT)=HTR(MT)*RMA(MT)*RNU(MT)*HD/RMR
    PC(MT)=PA(MT)+PB(MT)
    PA(4)=PA(4)+PA(MT)
    PB(4)=PB(4)+PB(MT)
    PC(4)=PC(4)+PC(MT)
400 CONTINUE
  WRITE(6,18) (PA(I),PB(I),PC(I),I=1,4)
  RETURN PT10950
  END PT10960

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SUBROUTINE PRINT2(AKN,XSTART,DT,RNU,RMA,DRF,FCP,HTP,UTLI,UTTI,VTSI
1,HTSI,DELANG,NWEDGE,XS,XCB,YCB,HTS,NTS,NTSP,UTL,UTT,VTS,I2,I3,IP)
DIMENSION RMA(1),RNU(1),DELANG(1),NWEDGE(1),XS(1),XCB(1),YCB(1)
DIMENSION HTS(3,I2,I3),NTS(3,I2,I3),NTSP(3,I2,I3),UTL(3,I2,I3)
DIMENSION UTT(3,I2,I3),VTS(3,I2,I3),UTLI(3,I2,I3),UTTI(3,I2,I3)
DIMENSION VTSI(3,I2,I3),HTSI(3,I2,I3)
COMMON /FIFTH/ND PT20060
----- PT20070

```

```

THE PURPOSE OF THIS SUBROUTINE IS TO PRINT OUT THE DISTRIBUTION PT20080
ON SURFACE OF THE SURFACE COEFFICIENTS PT20090
----- PT20100
----- PT20110
----- PT20120
----- PT20130
----- PT20140
----- PT20150
----- PT20160

```

FORMATS

```

8 FORMAT(//1X,45(**),' DISTRIBUTION ON SURFACE ',45(**)/1X,'INC.
1 TOT. INC. TOT. INC. TOT.'/1X,'SEGMENT GEOMETRY',
2 14X,'MOL. MOLE SAMP NUM. SKIN SKIN PRES- PRES-
3 HEAT HEAT'/* NO. CENTER DELX CENTER DELANG',4X,'WGHT. P
4 RACT.',10X,
5 'FLUX FRCTN FRCTN SURE SURE TRNSF TRNSF')
10 FORMAT(1X,I3,F8.3,F7.3,F9.3,F8.3,1X,2F8.4,I6,7F8.4)
11 FORMAT(37X,F8.4,' 1.0000',I6,7F8.4)

```

PT20230
***** PT20240
PT20250
PT20260

```

RMR=0.0
DO 50 MT=1,IP

```

50 RMR=RMR+RMA(MT)*RNU(MT)
 WRITE(6,8)
 I=0
 DO 110 N=1,ND
 DTY=DT*YCB(N)/180.
 P=XS(N)
 Q=2.*((ICB(N)-XSTART)*AKN-XS(N))
 ANGLE=0.
 R=0.
 J=0
 DO 105 L=1,2
 R=R+.5*ANGLE
 ANGLE=DELANG(L)
 R=R-.5*ANGLE
 ICNT=NWEDGE(L)
 IF(ICNT.EQ.0) GO TO 105
 PMLT=PCP/(DTY*ANGLE)
 QMLT=DRP/(DTY*ANGLE)
 SMLT=HTP/(DTY*ANGLE)
 DO 100 K=1,ICNT
 R=R+ANGLE
 I=I+1
 J=J+1
 M3=0
 P3=0.0
 Q3=0.0
 Q4=0.0
 R3=0.0
 R4=0.0
 S3=0.0
 S4=0.0
 DO 90 MT=1,IP
 M1=NTS(MT,N,J)
 M3=M3+M1
 P1=NTSP(MT,N,J)*PMLT*RNU(MT)
 P3=P3+P1
 Q1=SQRT(UTLI(MT,N,J)**2+UTTI(MT,N,J)**2)*RNU(MT)*RMA(MT)*QMLT/RMR
 Q2=SQRT(UTL(MT,N,J)**2+UTT(MT,N,J)**2)*RNU(MT)*RMA(MT)*QMLT/RMR
 Q3=Q3+Q1
 Q4=Q4+Q2
 R1=VTSI(MT,N,J)*RNU(MT)*RMA(MT)*QMLT/RMR
 R2=VTS(MT,N,J)*RNU(MT)*RMA(MT)*QMLT/RMR
 S1=HTSI(MT,N,J)*RNU(MT)*RMA(MT)*SMLI/RMR
 S2=HTS(MT,N,J)*RNU(MT)*RMA(MT)*SMLT/RMR
 R3=R3+R1
 R4=R4+R2
 S3=S3+S1
 S4=S4+S2
 90 WRITE(6,10)I,P,Q,R,ANGLE,RMA(MT),RNU(MT),M1,P1,Q1,Q2,R1,R2,S1,S2
 WRITE(6,11)RMR,M3,P3,Q3,Q4,R3,R4,S3,S4
 100 CONTINUE PT20660
 105 CONTINUE PT20670
 110 CONTINUE PT20680
 RETURN PT20690
 END PT20380

ORIGINAL DATA IS
 OF POOR QUALITY

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SUBROUTINE PRINT3(           IP, MJ, NS, NWEDG, I2, I3, I4, I5, RMA, XS, IWS,
1MS, TANGN, NTSP, NTCF, NTCV, PV) PT30020
DIMENSION RMA(1), XS(1), IWS(1), MS(1), TANGN(1), NTSP(3, I2, I3)
DIMENSION NTCF(3, I4), NTCV(3, I4, 2, I5, 3), PV(3, I4, 2, I5, 3), QUO(3, 3) PT3
----- PT30050
----- PT30060
----- PT30070
----- PT30080
----- PT30090
FORMATS PT30100
2 FORMAT(//1X, 40('*'), * MOMENTS OF INCIDENT DISTRIBUTION FUNCTIONS * PT30110
1, 40('*')) PT30110
6 FORMAT(/21X, I5, * UNCOLLIDED MOLECULES*, F8.4, 27X, I5, * COLLIDED MOLE
1CULES*, F8.4/) PT30170
8 FORMAT(I5, 4X, A3, 12(1X, F9.4)) PT30180
10 FORMAT(12X, 12(1X, F9.4)) PT30190
12 FORMAT(1H ) PT30200
***** PT30210
***** PT30220
***** PT30230
***** PT30240
WRITE(6, 2)
DO 155 I=1, NS
MR=MS(I) PT30260
ITT=IWS(I)
N=(MR-1)*NWEDG+ITT PT30270
DO 150 MT=1, IP PT30280
A=0.
B=1. PT30370
IC=NTCF(MT, I)
ID=NTSP(MT, MR, ITT)-IC PT30380
E=NTSP(MT, MR, ITT)
IF(E.LE.0.) GO TO 110 PT30400
A=IC/E
B=1.-A
110 WRITE(6, 6) IC, A, ID, B
E=NTCF(MT, I) PT30450
DO 121 NC=1, 2
DO 120 K=1, 3
QUO(NC, K)=0. PT30460
IF(E.EQ.0.) GO TO 120 PT30470
QUO(NC, K)=NTCV(MT, I, NC, 1, K)/E PT30480
120 CONTINUE PT30490
E=NTSP(MT, MR, ITT)-NTCF(MT, I) PT30500
121 CONTINUE PT30510
E=NTSP(MT, MR, ITT)-NTCF(MT, I) PT30520
122 CONTINUE PT30530
WRITE(6, 8) N, RMA(MT), PV(MT, I, 1, 1, 1), QUO(1, 1), PV(MT, I, 1, 1, 2), QUO
1(1, 2), PV(MT, I, 1, 1, 3), QUO(1, 3), PV(MT, I, 2, 1, 1), QUO(2, 1), PV(MT, I, 2, 1, PT30550
22), QUO(2, 2), PV(MT, I, 2, 1, 3), QUO(2, 3) PT30560
IF(MJ.EQ.1) GO TO 150 PT30570
DO 140 J=2, MJ PT30580
E=NTCF(MT, I)
DO 131 NC=1, 2 PT30590
DO 130 K=1, 3 PT30600
QUO(NC, K)=0. PT30610
IF(E.EQ.0.) GO TO 130 PT30620
PT30630

```

130	QUO(NC,K)=NTCV(MT,I,NC,J,K)/E	ORIGINAL PAGE IS	PT30640
131	CONTINUE	OF POOR QUALITY	PT30650
140	E=NTSP(MT,MR,ITT)-NTCP(MT,I)		PT30660
141	CONTINUE		PT30670
142	WRITE(6,10) FV(MT,I,1,J,1),QUO(1,1),FV(MT,I,1,J,2),QUO(1,2),FV(MT,I,1,J,3),QUO(1,3),FV(MT,I,2,J,1),QUO(2,1),FV(MT,I,2,J,2),QUO(2,2),FV(MT,I,2,J,3),QUO(2,3)	PT30680	
150	WRITE(6,12)		PT30690
155	CONTINUE		PT30700
	RETURN		PT30710
	END		PT30720
			PT30730
	SUBROUTINE PRINT4(MSP,CHI,RNU,I,TRP,NUMCEL,FDN,WTM,DB,NS,TMP,XV, 1 YV,ZV,NB,IC,YC,ZC,LEV,LKW)		
	INTEGER*2 LKW(1)		PT50050
	INTEGER*2 NB,NUMCEL,NS		PT50060
	DIMENSION FDN(1),RNU(1),CHI(1),WTM(1),NUMCEL(1),TMP(I,1),TRP(I,1)		PT50070
	DIMENSION DB(I,1),NB(I,1),XV(I,1),YV(I,1),ZV(I,1),DBT(3),NS(I,1)		PT50080
	DIMENSION IC(1),YC(1),ZC(1),LEV(1)		
	COMMON /FORTH/NBX		
	-----		PT50090
	-----		PT50100
	-----		PT50110
	-----		PT50120
	-----		PT50130
	-----		PT50140
	-----		PT50150
	1 FORMAT(//1X,45('**'),' INSTANTANEOUS FLOW FIELD INFORMATION ',45('**'))		PT50160
	2 FORMAT(/2X,'LEVEL=',I3,3X,'WEIGHTING FACTOR (MAX)=',I3,3X,'WEDGE AN		PT50170
	GLE =',F7.2,' DEGREES',3X,' RADIAL POSITION =',E11.3/2X,'BOX# X PO		
	2SITION SAMP DENSITY MACH NO X VEL. Y VEL. Z VEL. T(KIN) T(RO		
	3T) TEMP.',14X,'MOLE FRACTIONS')		
	3 FORMAT(//1X,46('**'),' ACCUMULATED FLOW FIELD INFORMATION ',46('**'))		
	4 FORMAT(1X,I4,E11.3,I6,8F8.3,3X,3E11.3)		PT50220
	*****		PT50230
			PT50240
			PT50250
	IF(KS.EQ.0) WRITE(6,1)		
	IF(KS.NE.0) WRITE(6,3)		
	DO 40 MT=1,3		
40	DBT(MT)=0.0		
	FDA=0.		
	CHT=0.		
	DO 50 MT=1,MSP		
	CHT=CHT+CHI(MT)*RNU(MT)		
50	FDA=FDA+FDN(MT)*WTM(MT)		
	YCT=0.0		
	ZCT=0.0		
	LEVEL=1		

```

DO 110 M=1,NBX
N=NUMCEL(M)
IF (N.LE.0) GO TO 110
IF ((ZC(M).EQ.ZCT).AND.(YC(M).EQ.YCT)) GO TO 52
ZCT=ZC(M)
YCT=YC(M)
IF (M.GE.LEV(1)) LEVEL=2
IF (M.GE.LEV(2)) LEVEL=3
WRITE(6,2) LEVEL,LKW(N),ZCT,YCT
52 XCT=XC(M)
NSAMP=0
DBA=0.
XVM=0.
YVM=0.
ZVM=0.
TMRPM=0.
TRPM=0.
E=0.
F=0.
DO 100 MT=1,MSP
NSAMP=NSAMP+NS(MT,N)
XVM=XVM+XV(MT,N)*RNU(MT)*WTM(MT)*NB(MT,N)
YVM=YVM+YV(MT,N)*RNU(MT)*WTM(MT)*NB(MT,N)
ZVM=ZVM+ZV(MT,N)*RNU(MT)*WTM(MT)*NB(MT,N)
DBA=DBA+DB(MT,N)*WTM(MT)
TMRPM=TMRPM+WTM(MT)*RNU(MT)*TMRP(MT,N)*NB(MT,N)
TRPM=TRPM+RNU(MT)*NB(MT,N)*TRP(MT,N)
E=E+WTM(MT)*RNU(MT)*NB(MT,N)
100 F=F+RNU(MT)*NB(MT,N)
DBA=DBA/FDA
IF (E.EQ.0.0) GO TO 55
XVM=XVM/E
YVM=YVM/E
ZVM=ZVM/E
VS=XVM**2+YVM**2+ZVM**2
TMRPM=TMRPM/E-VS
TRPM=TRPM/F
55 CONTINUE
TTM=(TMRPM+TRPM)/(2.5+CHT)
TMRPM=TMRPM/1.5
IF (CHT.NE.-1.) TRPM=TRPM/(1.+CHT)
AMS=SQRT(VS)
IF (TTM.GT.0.) AMS=SQRT((5.+2.*CHT)*VS/(TTM*(3.5+CHT)))
CCZ=COS(ZCT/57.29578)
SCZ=SQRT(1.-CCZ**2)
RVM=ZVM*SCZ-YVM*CCZ
TVM=YVM*SCZ+ZVM*CCZ
DO 60 MT=1,MSP
DBT(MT)=RNU(MT)*NB(MT,N)
IF (F.NE.0.) DBT(MT)=DBT(MT)/F
60 CONTINUE
WRITE(6,4) M,XCT,NSAMP,DBA,AMS,XVM,RVM,TVM,TMRPM,TRPM,TTM,(DBT(J),
1J=1,3)
110 CONTINUE
RETURN

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PT50290

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PT50480

FILE: GKBEXT DECK A

PRINCETON UNIVERSITY TIME-SHARING SYSTEM

END

PT50490

/LKED.SYSLMOD DD DSN=U.GKBSPACE.ROTATION(FORTH),
DISP=OLD, UNIT=3350, VOL=SER=RES101, SPACE=

/LKED.SYSIN DD *

INCLUDE SYSLMOD(FORTH)

ENTRY MAIN

/ EXEC COMPRESS,DSN='U.GKBSPACE.ROTATION',RLSE=RLSE

OPTIONAL
OR PRIMARY

APPENDIX B

MASTER'S THESIS OF Y.P. TSAI

COLLISION INDUCED VIBRATIONAL TRANSITION
PROBABILITIES IN DIATOMIC MOLECULES

ABSTRACT

In order to improve the Monte-Carlo Direct Simulation calculations for hypersonic flow in the transition regime, we have to incorporate the effects due to vibrational non-equilibrium and potential dissociation for a diatomic gas. The state-to-state transition probabilities are desired. In this paper, we model the diatomic molecule as a harmonic oscillator which collides with another molecule collinearly. Two different methods have been developed, the first one is the semi-classical treatment and the second one is the fully quantum mechanical approach. The interaction potential between two molecules is assumed to be the Lennard-Jones 12-6 interaction law which is a small perturbation to the colliding system. Some numerical results of the state-to-state transition probabilities and comparisons are presented in Chapter 4. Discussions, which present the important aspects of this kind of problem for further study, are made in Chapter 5.

Chapter 1 Introduction

The characteristic flow in a highly rarefied gas is called "free molecular flow". In this regime the mean free path is large compared to the characteristic dimensions of an aerodynamic body in the flow; and molecules that impinge on the body, and are then reemitted from it will, in general, be far away from the body before they strike another molecule. The characteristic flow in a moderately rarefied gas is called "slip flow". The flow regime intermediate between slip and molecular flow is known as the "transition flow regime". It corresponds to densities for which the mean free path has the same general order of magnitude as the characteristic dimension of the flow field. There is a dimensionless parameter called the Knudsen number Kn , introduced to serve as a criterion for determining the relative importance of these rarefaction effects.

where

$$\text{Kn} = \frac{\lambda}{d} \sim \frac{M}{Re}$$

and

λ = molecular mean free path

d = characteristic dimension of vehicle

M = Mach Number

Re = Reynolds Number

Free molecular flow is usually defined as that flow for which $\text{Kn} > 10$. Slip flow is characterized by a Knudsen number of a few per cent, $0.01 < \text{Kn} < 0.1$, and the intermediate transition regime corresponds to Knudsen number in the range $0.1 < \text{Kn} < 10$.

These values are, of course, arbitrary. Since the Knudsen number is defined as the ratio of particle mean free path to body size, it therefore increases with altitude for a fixed size body travelling through the atmosphere.

Generally, for a body of the order of a meter at altitudes 150 Kilometers above the earth's surface, the mean free path of the particles is much larger than vehicle size (here the free stream mean free path $\lambda \approx 40\text{m}$). Hence the particle-particle collision process in the vicinity of the vehicle need not be considered and the relation of measurement on the surface to atmospheric properties is explained in a relatively straightforward manner through free molecular theory. Below 90Km, (here the free stream mean free path $\lambda \approx 2.5\text{ cm}$), the Knudsen number is so small that the fluid can be treated as a continuum with limited influence of transport properties and slip-flow boundary conditions. However, between these two zones, i.e. in the lower region of the thermosphere, is what we called the "transition flow regime" in which neither of the limiting theories is applicable.

Flight at very high altitudes often involves extremely high velocities and resulting high gas temperature. At velocities which correspond to effective temperature of the order of a few thousand degrees Kelvin, the so called "real gas" effects associated with vibration, dissociation, and ionization of the gas molecules can begin to be of importance. We are interested in atmospheric entry of a

vehicle such as the space shuttle. The Mach number is generally above 20 in passing through "transition flow regime" and the flow is hypersonic there. At such high speeds, any significant number of collisions between incoming particles and those reflected from the body can produce extreme changes in the environment near the surface of the vehicle. Measurements on the surface are thus strongly affected making inference of conditions in the ambient atmosphere extremely difficult. For the purpose of data interpretation, previous calculation based on direct simulation Monte-Carlo computer technique have been developed by G.A. Bird(1). By using the Monte-Carlo method, the real gas molecules are replaced by their statistical models, and the motion of one or more of the chosen particles is traced by the computer. In the "Bird" method the real gas molecules are simulated by several thousand modeled molecules, rigid spheres in the simplest version(1). Theoretical calculations of the heat transfer and aerodynamic characteristics of a body submerged in the transition flow regime may be carried out in this way. Modifications involving more realistic interaction laws have been also carried out. Molecules, however, contain internal structure, which is important primarily for its effect on the energy content of the flowing gas. Being composed of nuclei and electrons that have motion relative to the center of mass of the molecule, the molecules can possess rotational and vibrational as well as electronic

internal states. Therefore, translational energy is not necessarily conserved in all collisions. The hypersonic flow past the sharp leading edge of a flat plate incorporating the effects of rotational non-equilibrium for a diatomic gas was studied by D.I. Pullin, J.K. Harvey, and G.K. Bienkowski(2), and subsequently applied to other blunt body problems(3,4). A different model was used to incorporate the same effect in other works such as references 5 and 6. At the present stage, we want to improve the Monte-Carlo calculations by including the effects due to the vibrational degree of freedom of diatomic molecules.

It is convenient to divide vibrational energy exchange into two cases: the V-V process, in which the total vibrational quantum number of the system is unchanged, and the V-T process in which energy is exchanged between translation and vibration without conserving the vibrational quantum. For harmonic oscillators in the V-V process, the amount of vibrational energy lost by one molecule is gained by the other and no vibrational-translational energy transfer occurs. Considerable interest has been shown in the details of inelastic molecular collisions. The treatment of scattering between particles with internal structure is capable of producing differential and total cross section for state-to-state transitions. Our purpose is to calculate scattering cross sections for different transitions with known initial and final states as a function of collision energy or initial relative velocity. We emphasize here that

we need state-to-state transition cross sections rather than some overall "rates of transition" (rate coefficients) which are averaged over an equilibrium velocity distribution for the relative translational motion (e.g. Maxwellian distribution). The reasons are twofold:

1. We are dealing with a hypersonic flow system in a highly non-equilibrium state with a non-Maxwellian distribution of relative velocities of collision. The relative contributions of different energy molecules to the overall rates may therefore be drastically different than in the equilibrium state. This effect is accentuated by the steep rise of cross-sections with energy coupled to the generally decreasing magnitude of the distribution function with energy. This unifies that small changes in the fraction of molecules with high energies due to the non-equilibrium aspect of the flow can have extreme effects on inelastic processes without corresponding effects on mean properties such as density or fluid momentum.
2. The Direct Simulation Monte Carlo Code (DSMCC) consists of tracing a set of "test" molecules through a designated volume surrounding the body. The velocities and internal states of the simulated molecules are altered on the basis of collisions computed (as determined by local

collision probabilities) at fixed time steps. The positions are then advanced to new values on the basis of motion through the time step increments. This detailed computation of individual molecular collisions requires, in principle, state to state transition probabilities in order to incorporate the inelastic energy exchange into the computation of molecular velocities and internal states after collision. While the results we desire must ultimately be consistent with overall measured rates the level of detail necessary within the program is well beyond the level of availability of experimental data.

Early theoretical studies of vibrational, rotational, and translational energy transfer in collisions were based on approximate analytical solutions to the quantum mechanical and classical equations of motion. The method of Zener (7), later to become known as the distorted wave method, and the Born (8) approximation are leading examples of approximate solutions to quantum mechanical collision problems based on first order perturbation theory. A more detailed literature review on previous work in the field of vibrational collisions is given in the next Chapter.

Instead of doing an approximate treatment of a three-dimensional realistic system, in this work, we do an exact numerical treatment of a simpler model one-dimensional system

which has some important features in common with the real one. A calculation of transition probabilities for vibrational-vibrational-translational energy transfer in a collision of two diatomic molecules is to be presented. This simple collision model is approximate, utilizing a collinear collision of harmonic oscillators with an exponential repulsion between center atoms and no chemical reaction between the molecules. It may be argued that the configuration allowing the most efficient transfer of energy between translation and vibration is that in which the atoms are collinear. Collinear or head-on collisions make the most significant contribution to the transition probability. The averaged probability is equal to the probability of excitation in a head-on collision times a "steric factor" smaller than unity which takes account of unfavorable trajectories. For homonuclear molecules it is usually taken as $\frac{1}{3}$ (the average of $\cos^2\theta$ taken over a sphere). A more detailed theory for the steric factor of linear molecules has been propounded by Herzfeld (9). It is our belief that an accurate treatment of a collinear model is of more worth than an approximate result for the three-dimensional problem. The latter approach frequently contains errors which are difficult to estimate. The general magnitudes and trends of the transition probability obtained by this restricted treatment can show us qualitatively, or semi-quantitatively, some characteristic features of the problem.

Two different methods for calculating transition

probabilities are discussed in this paper. One is a semi-classical approach, and the other one is a fully quantum mechanical treatment. In the semi-classical calculation, it is assumed that the vibrational amplitude of the harmonic oscillators are small and so the molecular oscillations do not greatly affect the external classical collision trajectory. The trajectory can be calculated from the classical equation of motion. The classical trajectory is assumed to define a time-dependent perturbation potential for the colliding system and quantum theory is used to derive the transition probability. Essentially, the Schrodinger equation is solved subject to certain initial conditions according to time-dependent perturbation theory. A detailed discussion about this theory can be found in the book Quantum Mechanic by Schiff (10).

In the quantum mechanical calculation of transition probability, the wavefunction of the whole system is expanded in terms of the complete set of eigenfunctions of vibrational states of the diatomic molecule. With the aid of the orthonormality property of these eigenfunctions, a set of coupled second order differential equations is obtained for the translational wavefunctions. The transition probability is given by the solution to this set of equations in the asymptotic region, subject to appropriate boundary condition. Several different numerical methods for solving the set of coupled equations have been developed by Diestler and McKoy (11).

Riley and Kuppermann (12), and Gutshick et al (13). Solving a problem of quantum scattering between two diatomic molecules is then reduced to the task of finding a good numerical scheme for integrating a system of coupled differential equations accurately. In this work, we use IBM IMSL ROUINE DGEAR to solve for the scattered wavefunction in the asymptotic region and then calculate the transition probability.

In Chapter 2, a brief review of previous semiclassical and quantum mechanical methods on vibrational collisions will be given so that we can identify the new points in the current work in Chapter 3. Chapter 3 is devoted to a discussion of the general theory. We derive some equations and expressions there which make a numerical algorithm feasible. In Chapter 4, theoretical results of state-to-state transition probability specifically for N_2-N_2 collisions at different relative velocities are presented. Comparison is made with published results. Finally, we discuss some important problems related to inelastic molecular collisions which deserve further study because they make extension to a more realistic treatment of molecular scattering possible.

Chapter 2 Literature Review

In 1931, Oldenberg (14) discussed molecular collision processes qualitatively to show the persistence of the rotational and vibrational motion. Zener (7) was the first to give a detailed mathematical treatment for collisions in which molecular vibrations are excited or de-excited. He restricted himself to collinear collisions between a diatomic molecule and an atom. His theory was based on the distorted wave method which includes direct transition from the initial state to the final state and assumes that the probability of transition is small. It is a perturbation method and cannot treat strongly coupled system. Takayanagi (15) then extended Zener's one-dimensional treatments to three-dimensional collisions. In order to save computational labor, the modified wave number approximation was introduced. Meanwhile, Schwartz, Slawsky and Herzfeld (16) gave a mathematical formulation for the vibrational transitions based on the distorted wave approximation due to Jackson and Mott (17), simplified by the modified wave number approximation due to Takayanagi, in diatom-diatom collisions. Their formulation is referred to as the SSH theory now and is a quantum mechanical result. For purpose of future reference, we describe SSH theory in more detail.

Consider the head-on collision between two diatomic molecules AB and CD (assume harmonic oscillators). For exponential intermolecular interaction between nearest

atoms B and C, one has,

$$V = V_0 \exp(-\alpha r_{BC}) \\ = V_0 \exp[-\alpha(R - \lambda_1 r_1 - \lambda_2 r_2)] \quad (2-1)$$

where R is the distance between centers of mass of two molecules, r_1 and r_2 are vibrational coordinates, $\frac{1}{\alpha}$ is the range of the potential and V_0 is a constant.

$$\lambda_1 = \frac{m_A}{m_A + m_B}, \quad \lambda_2 = \frac{m_D}{m_C + m_D}$$

m_i is the mass of the i -th atom. Solving the Schrodinger equation, the transition probability is given in the closed analytic form:

$$P(n_1 n_2 \rightarrow n'_1 n'_2) = \frac{\pi^2}{4} |v_1(n'_1 n_1)|^2 |v_2(n'_2 n_2)|^2 \frac{\Delta q^2}{\cosh \pi q' - \cosh \pi q} \quad (2-2)$$

where

$$v_1(n'_1 n_1) = \int_{-\infty}^{+\infty} z_1(n'_1, r) \exp[\alpha \lambda_1(r - r_{e1})] z_1(n_1, r) dr \quad (2-3)$$

$$v_2(n'_2 n_2) = \int_{-\infty}^{+\infty} z_2(n'_2, r) \exp[\alpha \lambda_2(r - r_{e2})] z_2(n_2, r) dr \quad (2-4)$$

$z_i(n_j, r)$ is the vibrational wave function for harmonic oscillator "i" in quantum state n_j ; r_{e1} and r_{e2} are equilibrium separations for AB and CD respectively, and

$$q = q_{n_1 n_2} = \frac{2kn_1 n_2}{\alpha} \quad (2-5)$$

$$q' = q_{n'_1 n'_2} = \frac{2kn'_1 n'_2}{\alpha} \quad (2-6)$$

$$\Delta q^2 = q'^2 - q^2 = - \left(\frac{8\mu}{h^2 \alpha^2} \right) [\epsilon_1(n'_1) + \epsilon_2(n'_2) - \epsilon_1(n_1) - \epsilon_2(n_2)] \quad (2-7)$$

The wave number of relative motion before and after collision are $k_{n_1 n_2}$ and $k_{n'_1 n'_2}$, which satisfy the energy conservation law:

$$\epsilon_1(n_1) + \epsilon_2(n_2) + \frac{h^2 k^2 n_1^2 n_2^2}{2\mu} = \epsilon_1(n'_1) + \epsilon_2(n'_2) + \frac{h^2 k^2 n'_1 n'_2}{2\mu} \quad (2-8)$$

μ is the reduced mass of the whole system; $\epsilon_i(n)$ is the energy of molecule "i" in the n -th vibration state. For the special case, V-V transition, in which $n_1+n_2=n'_1+n'_2$ and $\Delta q=0$. Applying L'Hospital's rule, we get:

$$\lim_{q' \rightarrow q} \frac{q'^2 - q^2}{\cosh \pi q' - \cosh \pi q} = \frac{2q}{\pi \sinh \pi q}$$

Then the transition probability for V-V process becomes

$$p(n_1 n_2 \rightarrow n'_1 n'_2) = |v_1(n'_1 n_1)|^2 |v_2(n'_2 n_2)|^2 q^2 \quad (2-9)$$

SSH theory has been most widely used for quantitative comparison with experimental measurements of vibrational relaxation, but the coupling between rotation has been ignored.

Zelechow, Rapp and Sharp (ZRS) (18) have developed a semi-classical method for calculating transition probabilities for V-V and V-T energy transfer in a collision of two diatomic molecules. Their basic

assumption are:

- (1) The perturbation potential is linearized in the oscillator coordinates.
- (2) The collision velocity is not too high (e.g. the upper limit for N_2-N_2 collision is 10Km/sec) and the collision induced time-varying force constant $k'(t)$ is small compared to k , the characteristic force constant of the molecule.

Under these two conditions, Kerner (19) method can be applied to solve the Schrodinger equation and closed form analytical results are obtained. However, this approach restricts itself to the transitions of processes of symmetric type only. The general formula is:



where n, m are vibration quantum numbers before collision and n', m' are that after collision. The collision is symmetric in the sense that the two B atoms are in the center.

The development of the high speed electronic computer has made it possible to solve the collision problem by direct numerical techniques. T.E. Sharp and D. Rapp (20) have calculated the vibrational transition probabilities for collisions between a diatomic molecule and an atom. In their semi-classical treatment, an N -state approximation method is used, in which the total wave function is expanded in terms of N eigenfunctions of stationary states of the system including the initial, final and all

energetically intervening states. A Runge-Kutta single-step integration method is employed in the computation program. Generally, the value of N needed in expanding the total wave function increases with collision velocity. An "exact" solution for any transition probability $P_{j \rightarrow k}$ is reached when the addition of more states to the computation results in no significant change in $P_{j \rightarrow k}$. We extend this method to collisions between two diatomic molecules (Chapter 3-A).

In quantum mechanical treatment of collisions between two diatomic molecules AB and CD, taking B and C as the inner atoms of the system, the total wave function is expanded in terms of normalized vibrational wave functions $z_{AB}(n_1, r_1)$ and $z_{CD}(n_2, r_2)$. That is:

$$\psi = \sum_{n_1} \sum_{n_2} f_{n_1 n_2}(R) z_{AB}(n_1, r_1) z_{CD}(n_2, r_2)$$

Inevitably, we have to solve a system of coupled differential equations, which are equivalent to the Schrodinger equation, of the following form (detailed discussions will be given in Chapter 3-C).

$$\frac{\hbar^2}{2\mu} \frac{d^2}{dR^2} + k_{n_1 n_2}^2 f_{n_1 n_2}(R) = \sum_{n_1'} \sum_{n_2'} \langle n_1' n_2' | V | n_1 n_2 \rangle f_{n_1' n_2'}(R) \quad (2-10)$$

where $\langle n_1' n_2' | V | n_1 n_2 \rangle$ is the matrix element. In principle if we can obtain the solution to equation (2-10) with the asymptotic form

$$\begin{aligned}
 f_{n'_1 n'_2}(R) &\rightarrow 0 & R \rightarrow -\infty \\
 f_{n'_1 n'_2}(R) &\rightarrow \delta_{n'_1 n'_2} \delta_{n'_2 n'_2} \exp(-ik_{n'_1 n'_2} R) \\
 &\quad + A_{n'_1 n'_2} n'_1 n'_2 \exp(+ik_{n'_1 n'_2} R) & R \rightarrow \infty
 \end{aligned} \tag{2-11}$$

the probability per collision for transition $(n_1, n_2) \rightarrow (n'_1, n'_2)$ will be given by:

$$P(n_1 n_2 \rightarrow n'_1 n'_2) = \frac{k_{n'_1 n'_2}}{k_{n_1 n_2}} |A_{n'_1 n'_2} n'_1 n'_2|^2 \tag{2-12}$$

A number of numerical methods (7-9) have been proposed for solving the systems of equations (2-10). However, due to the rapidly oscillating wavelike solutions to the Schrodinger equation, the numerical technique is not straightforward. Riley (12) developed the initial-value technique with periodic "reorthogonalization". Gadschick et al. (13), on the other hand introduced a technique of integration using Dirichlet boundary condition and simple one-step Euler integration. A new method for constructing wave function for bound states and scattering has been proposed by Roy G. Gordon (21), perhaps this procedure can save much computer time. Our quantum mechanical treatment of this molecule scattering problem is similar to the method due to Riley and Kuppermann (12). It is relatively simple and straightforward, but in our procedures, the virtual states (energetically inaccessible) are not included in the total wavefunction expansion.

Chapter 3 Theory

A. General Formalism

The collision model is shown in Fig. 3-1.

Fig. 3-1 Collision Coordinates

This figure is the collinear collision configuration between two diatoms AB and CD. Assuming that CD is the target, and AB is the incident projectile from right. The laboratory coordinates of A, B, C and D are x_A , x_B , x_C and x_D ; their masses are m_A , m_B , m_C and m_D respectively. Let V_{AB} and V_{CD} be the binding potential of molecules AB and CD. The short range interaction is assumed to be a sum of interatomic interactions,

$$V_{INT} = V_{AC}(x_A - x_C) + V_{AD}(x_A - x_D) + V_{BD}(x_B - x_D) + V_{BC}(x_B - x_C)$$

The interatomic potentials are exponentially decreasing functions, so that for the collinear configuration under consideration, only the term $V_{BC}(x_B - x_C)$ is important and thus

$$V_{INT} \approx V_{BC}(x_B - x_C)$$

The Schrodinger equation for the system is:

$$\left\{ -\frac{\hbar^2}{2M_A} \frac{\partial^2}{\partial x_A^2} - \frac{\hbar^2}{2M_B} \frac{\partial^2}{\partial x_B^2} - \frac{\hbar^2}{2M_C} \frac{\partial^2}{\partial x_C^2} - \frac{\hbar^2}{2M_D} \frac{\partial^2}{\partial x_D^2} + V_{AB}(x_A - x_B) + V_{CD}(x_C - x_D) + V_{INT}(x_B - x_C) \right\} \psi(x_A, x_B, x_C, x_D) = E_{TOT} \psi(x_A, x_B, x_C, x_D) \quad (3-1)$$

We designate the distance between the centers of mass of two molecules as R . In molecule AB the distance between the atoms is x ; in CD it is y . x and y are internal coordinates. i.e.

$$R = \frac{m_A x_A + m_B x_B}{m_A + m_B} - \frac{m_C x_C + m_D x_D}{m_C + m_D}$$

$$x = x_A - x_B$$

$$y = x_C - x_D$$

Let $R_{CM} = \frac{m_A x_A + m_B x_B + m_C x_C + m_D x_D}{m_A + m_B + m_C + m_D}$ which is the coordinate of center of mass of the whole system. In terms of the new coordinates (x, y, R, R_{CM}) the Schrodinger equation becomes:

$$\left\{ -\frac{\hbar^2}{2\mu_{AB}} \frac{\partial^2}{\partial x^2} - \frac{\hbar^2}{2\mu_{CD}} \frac{\partial^2}{\partial y^2} - \frac{\hbar^2}{2\mu} \frac{\partial^2}{\partial R^2} - \frac{\hbar^2}{2M} \frac{\partial^2}{\partial R_{CM}^2} + V_{AB}(x) + V_{CD}(y) + V_{INT}(R - \gamma_{AB}x - \gamma_{CD}y) \right\} \psi(x, y, R, R_{CM}) = E_{TOT} \psi(x, y, R, R_{CM}) \quad (3-2)$$

where $M = m_A + m_B + m_C + m_D$ = total mass of the system.

$\mu = \frac{(m_A + m_B)(m_C + m_D)}{M}$ = reduced mass of the system

$$\mu_{AB} = \frac{m_A m_B}{m_A + m_B} \quad = \text{reduced mass of the molecule AB}$$

$$\mu_{CD} = \frac{m_C m_D}{m_C + m_D} \quad = \text{reduced mass of the molecule CD}$$

$$\gamma_{AB} = \frac{m_A}{m_A + m_B}, \quad \gamma_{CD} = \frac{m_D}{m_C + m_D}$$

Since there is no external force applied to the system, the center of mass of the system moves like a free particle and its motion can be described by a plane wave $e^{iK_{CM}^R R_{CM}}$, and the energy of the center of mass, $T_{CM} = \frac{h^2 K_{CM}^2}{2M}$, is also a constant of motion. This does not affect the energy transfer and need not be considered further. We can remove the R_{CM} -dependent part of the wave function (x, y, R, R_{CM}) by separation of variables.

Let

$$\psi(x, y, R, R_{CM}) = \psi(x, y, R) e^{iK_{CM}^R R_{CM}}, \quad E_{TOT} = E + T_{CM}$$

Substituting these into equation (3-2), we arrive at a Schrodinger equation concerning the internal coordinates x, y and the relative motion R of the two colliding molecules as follows:

$$\left\{ -\frac{h^2}{2\mu_{AB}} \frac{\partial^2}{\partial x^2} - \frac{h^2}{2\mu_{CD}} \frac{\partial^2}{\partial y^2} - \frac{h^2}{2\mu} \frac{\partial^2}{\partial R^2} + V_{AB}(x) + V_{CD}(y) + V_{INT}(R - \gamma_{AB}x - \gamma_{CD}y) \right\}$$

$$\psi(x, y, R) = E\psi(x, y, R) \quad (3-3)$$

Define

$$X = x - x_{eq}$$

$$Y = y - y_{eq}$$

where x_{eq} and y_{eq} are equilibrium separations of AB and CD. Then X and Y are displacements from equilibrium of each oscillator. Now we introduce harmonic bonds (intramolecular potential) into AB and CD with force constants k_{AB} and k_{CD} , hence

$$V_{AB} = 1/2 k_{AB} X^2,$$

$$V_{CD} = 1/2 k_{CD} Y^2,$$

A conventional representation of the intermolecular potential energy curve is given by the Lennard-Jones 12-6 equation. Since the elementary models for energy transfer are based on exponential potential, the exponential function $V_{INT}(x_B - x_C) \sim e^{-(x_B - x_C)/L} - \epsilon$ must be fitted to the Lennard-Jones potential (Appendix 1), where L is a parameter characterizing the range of the interaction. Landau and Teller (22) assumed that only the short range repulsive part of the intermolecular potential is steep enough to influence energy transfer, so that the long-range attractive potential ϵ can be neglected. The molecular interaction is then assumed to be an exponential repulsion between atoms C and B.

Let v = initial relative velocity

$E = 1/2 \mu v_0^2$ = initial relative kinetic energy

$$\bar{R} = R - R_T$$

where R_T is the distance at the classical turning point. The potential energy in equation (3-3) may be expressed as

(16) :

$$v_{\text{INT}}(x, y, \bar{R}) = E_0 \exp \left[\frac{-1}{L} (\bar{R} - \gamma_{AB}x - \gamma_{CD}y) \right] \quad (3-4)$$

Equation (3-3) can be written in the form:

$$\left\{ \frac{-h^2}{2\mu_{AB}} \frac{\partial^2}{\partial x^2} - \frac{-h^2}{2\mu_{CD}} \frac{\partial^2}{\partial y^2} - \frac{h^2}{2\mu} \frac{\partial^2}{\partial \bar{R}^2} + \frac{1}{2} k_{AB}x^2 + \frac{1}{2} k_{CD}y^2 + E_0 \exp \left[\frac{-1}{L} (\bar{R} - \gamma_{AB}x - \gamma_{CD}y) \right] \right\} \psi(x, y, \bar{R}) = E\psi(x, y, \bar{R}) \quad (3-5A)$$

or

$$\left\{ -\frac{h^2}{2\mu_{AB}} \frac{\partial^2}{\partial x^2} - \frac{h^2}{2\mu_{CD}} \frac{\partial^2}{\partial y^2} + \frac{1}{2} k_{AB}x^2 + \frac{1}{2} k_{CD}y^2 + E_0 \exp \left[\frac{-1}{L} (\bar{R} - \gamma_{AB}x - \gamma_{CD}y) \right] \right\} \psi(x, y, \bar{R}) = ih \frac{\partial}{\partial t} \psi(x, y, \bar{R}) \quad (3-5B)$$

We shall solve the time-independent Schrodinger equation (3-5A) by purely quantum mechanical method and the time-dependent Schrodinger equation (3-5B) by semi-classical method.

B. Semi-classical Calculation

Since the deBroglie wavelength of the relative motion of two molecules is usually very small compared with atomic dimensions (e.g. for N_2-N_2 collisions at velocity 5Km/sec, the deBroglie wavelength is of the order 10^{-15} cm, however the dimension of N_2 molecules is of the order of a few Å), it is a fairly good approximation to use the classical trajectory for relative motion. The classical equations of motion are:

$$\mu \frac{d^2 \bar{R}}{dt^2} = - \frac{\partial}{\partial \bar{R}} V_{INT}(X, Y, \bar{R}) \quad (3-6A)$$

$$\mu_{AB} \frac{d^2 X}{dt^2} = - k_{AB} X - \frac{\partial}{\partial X} V_{INT}(X, Y, \bar{R}) \quad (3-6B)$$

$$\mu_{CD} \frac{d^2 Y}{dt^2} = - k_{CD} Y - \frac{\partial}{\partial Y} V_{INT}(X, Y, \bar{R}) \quad (3-6C)$$

Generally, the incident energy is much larger than the change in the vibrational energy, or in other words, due to small transition probabilities, only a small fraction of the translational energy is transferred to vibrational energy. We may then assume that during the collision, the vibrational amplitudes of the oscillators are not driven to large values, that means

$$\begin{aligned} X &\ll L, \\ Y &\ll L. \end{aligned} \quad (3-7)$$

In a series calculation given by Wolfberg and Kelley (23), we can see that conditions (3-7) are justifiable. Wolfberg and Kelley have calculated the energy transfer for collisions involving two harmonic oscillators via an exponential collision with $L=0.22\text{\AA}$. Other parameters and data are: $m_A=m_B=m_C=m_D=12\text{a.m.u.}$, angular frequency $=2.3 \times 10^{14} \text{ sec}^{-1}$ (cf: for N_2-N_2 collisions $m_A=m_B=m_C=m_D=14 \text{ a.m.u.}, =4.45 \times 10^{14} \text{ sec}^{-1}$), the initial energy $E_0 = 5.078 \text{ ev}$ (corresponds to $v_0 = 9 \text{ Km/sec}$), then the vibrational energy transferred to each diatomic molecule is $\Delta E_{CD} = 1.78 \times 10^{-3} \text{ ev}$. Obviously, both ΔE_{AB} and ΔE_{CD} are much less than E_0 .

The vibrational amplitude never exceeds 0.007 Å, and since $L=0.22$ Å the conditions (3-7) are fairly well satisfied. At low velocities (still high enough so that the deBroglie wavelength is much less than the atomic dimension), the energy transfer becomes smaller because the oscillator can readjust adiabatically to the perturbation caused by the incident particle. Conditions (3-7) are satisfied even better. Under conditions (3-7), the molecular oscillations do not greatly affect the external classical collision trajectory. Therefore, one can neglect the motion in X and Y in treating the motion in R. Equations (3-6A), (3-6B), and (3-6C) may then be replaced by:

$$\mu \frac{d^2\bar{R}}{dt^2} = - \frac{\partial}{\partial \bar{R}} V_{INT}(\bar{R}, X=0, Y=0) \quad (3-8)$$

Solving equations (3-8) with V_{INT} given by (3-4), one finds that the trajectory R(t) satisfies the relation

$$\exp(-\frac{\bar{R}(t)}{L}) = \operatorname{sech}^2(\frac{u_0 t}{2L}) \quad (3-9)$$

R(t) is then inserted into the interaction potential function $V_{INT}(X, Y, \bar{R})$ to obtain $V_{INT}(X, Y, t)$. In this semi-classical treatment, $V_{INT}(X, Y, t)$ is used as a transition inducing perturbation acting upon a quantum mechanical harmonic oscillator. Finally, from equation (3-5B) we get the time-dependent Schrodinger equation to be

solved numerically by first order time-dependent perturbation theory.

$$\left\{ -\frac{\hbar^2}{2\mu_{AB}} \frac{\partial^2}{\partial x^2} + \frac{1}{2} k_{AB} x^2 - \frac{\hbar^2}{2\mu_{CD}} \frac{\partial^2}{\partial y^2} + \frac{1}{2} k_{CD} y^2 + E_0 \exp\left(\frac{\gamma_{AB}x + \gamma_{CD}y}{L}\right) \right. \quad (3-10)$$

$$\left. \operatorname{sech}^2\left(\frac{u_C t}{2L}\right) \right\} \quad \psi(x, y, t) = i\hbar \frac{\partial}{\partial t} \psi(x, y, t)$$

There is one more comment about the assumptions of this semi-classical model. In principle, if the deBroglie wavelength associated with the relative motion is much less than the atomic dimensions, it is a good approximation to use classical mechanics for solving the relative trajectory. As just mentioned before, if the criteria (3-7) are fulfilled, in other words, each of the oscillation amplitudes of the two colliding diatomic molecules is very small compared to the characteristic range of the interaction potential, a further simplification can be made in solving for $\bar{R}(t)$, the relative trajectory. These procedures correspond to the "approximate" classical method, because the trajectory $\bar{R}(t)$ is determined with x and y set equal to zero. As a result, these calculations do not include the conservation of energy, and E is assumed to remain as the energy in coordinate \bar{R} , regardless of how much excitation occurs in the oscillators. For more accurate semi-classical calculations, one must solve equations (3-6A), (3-6B), and (3-6C) for "exact" classical trajectory $\bar{R}(t)$ to be used as a time-dependent perturbation. Prior to the work of Wolfsberg and

Kelley (23), it has been thought that the criteria (3-7) are automatically satisfied for all low-velocity collisions. Actually, this is not the case. Wolfsberg and Kelley proved that the approximate classical method should be limited to collisions between a light particle and a heavy oscillator.

For our case of N_2-N_2 collision, Wolfsberg and Kelley's requirement is satisfied. The energy transferred to each diatomic species is very small compared to the kinetic energy of relative motion. Both molecules are negligibly distorted during collision approach and the molecular oscillations never deviate substantially from their equilibrium configurations. The semi-classical approach is a good approximation even within the high velocity range in which we are interested. The reason is that as v_0 increases, the deBroglie wavelength characterizing the relative motion in R is small compared to the distance over which the interaction potential varies significantly in R , i.e. $L.k_{n1n2} \gg 1$. This implies that both q and q' defined by equations (2-5) and (2-6) are much greater than unity. In this limit, $q \gg 1$ and $q' \gg 1$, the quantum mechanical results of equation (2-2) are reduced to the classical results of Landau and Teller (22).

Let ω_{AB} and ω_{CD} be the angular frequencies for the AB and CD molecules, their isolated Hamiltonians are $H^{(AB)}(X)$ and $H^{(CD)}(Y)$ respectively, then

$$\omega_{AB} = \frac{k_{AB}}{\mu_{AB}}$$

$$\omega_{CD} = \frac{k_{CD}}{\mu_{CD}}$$

$$H^{(AB)}(X) = -\frac{\hbar^2}{2\mu_{AB}} \frac{\partial^2}{\partial X^2} + \frac{1}{2} \mu_{AB} \omega_{AB}^2 X^2$$

$$H^{(CD)}(Y) = -\frac{\hbar^2}{2\mu_{CD}} \frac{\partial^2}{\partial Y^2} + \frac{1}{2} \mu_{CD} \omega_{CD}^2 Y^2$$

Suppose the individual eigenfunctions of $H^{(AB)}(X)$ and $H^{(CD)}(Y)$ are $\phi_n^{(AB)}(X)$ and $\phi_j^{(CD)}(Y)$, we have:

$$H^{(AB)}(X) \phi_n^{(AB)}(X) = (n + \frac{1}{2}) \hbar \omega_{AB} \phi_n^{(AB)}(X) \quad n = 0, 1, 2, \dots$$

$$H^{(CD)}(Y) \phi_j^{(CD)}(Y) = (j + \frac{1}{2}) \hbar \omega_{CD} \phi_j^{(CD)}(Y) \quad j = 0, 1, 2, \dots$$

Set

$$\phi_{nj}(X, Y) = \phi_n^{(AB)}(X) \phi_j^{(CD)}(Y) \quad (3-11)$$

$$W_{ij} = (n + \frac{1}{2}) \hbar \omega_{AB} + (j + \frac{1}{2}) \hbar \omega_{CD} \quad (3-12)$$

$$H_0(X, Y) = H^{(AB)}(X) + H^{(CD)}(Y) \quad (3-13)$$

$H_0(X, Y)$ is the unperturbed Hamiltonian of the system. It is obvious that $H_0(X, Y) \phi_{nj}(X, Y) = W_{nj} \phi_{nj}(X, Y)$. (3-14)

W_{ij} are eigenvalues of the unperturbed Hamiltonian $H_0(X, Y)$. The solution to the equation (3-10), i.e. the total wavefunction for the system of two oscillators $\psi(X, T, t)$ can be expanded in terms of the individual harmonic oscillator wavefunctions $\phi_n^{(AB)}(X)$ and $\phi_j^{(CD)}(Y)$.

$$\psi(x, y, t) = \sum_n \sum_j a_{nj}(t) \phi_n^{(AB)}(x) \phi_j^{(CD)}(y) e^{-i(n+\frac{1}{2})\omega_{AB}t} e^{-i(j+\frac{1}{2})\omega_{CD}t} \quad (3-15)$$

where the expansion coefficients a_{nj} depend on time.

If the oscillators AB and CD are initially in state N and J respectively, then the initial conditions are:

$$n = 0, 1, 2, \dots$$

$$a_{nj}(-\infty) = \delta_{nN} \delta_{jJ}$$

(Appendix 2)

$$j = 0, 1, 2, \dots$$

The probability of the system ending up with AB in state Q and CD in state K is

$$P_{NJ+QK} = |a_{QK}(+\infty)|^2$$

We are now in the position to solve for the expansion coefficients $a(t)$. Substituting equation (3-15) into Schrodinger equation (3-10), and using the relations (3-11) \rightarrow (3-14), we find:

$$(ih) \sum_{nj} \frac{d}{dt} a_{nj}(t) \phi_n^{(AB)}(x) \phi_j^{(CD)}(y) e^{-i(n+\frac{1}{2})\omega_{AB}t} e^{-i(j+\frac{1}{2})\omega_{CD}t} \quad (3-16)$$

$$= E_0 \exp\left(\frac{\gamma_{AB}x + \gamma_{CD}y}{L}\right) \operatorname{sech}^2\left(\frac{u_0 t}{2L}\right) \sum_{nj} a_{nj}(t) \phi_n^{(AB)}(x) \phi_j^{(CD)}(y) e^{-i(n+\frac{1}{2})\omega_{AB}t} e^{-i(j+\frac{1}{2})\omega_{CD}t}$$

Equation (3-16) is multiplied by $\phi_n^{(AB)*}(x) \phi_j^{(CD)*}(y)$ on both sides, where $\phi_n^{(AB)*}$ is the complex conjugate function of $\phi_n^{(AB)}$ and $\phi_j^{(CD)*}(y)$ is the complex conjugate of $\phi_j^{(CD)}(y)$, and then integrated. Equation (3-16) then becomes:

$$(ih) \sum_{nj} \frac{d}{dt} a_{nj}(t) \delta_{nn}, \delta_{jj}, e^{-i(n+\frac{1}{2})\omega_{AB}t} e^{-i(j+\frac{1}{2})\omega_{CD}t} \\ = E_0 \operatorname{sech}^2\left(\frac{u_0 t}{2L}\right) \sum_{nj} a_{nj}(t) \left[\int_{-\infty}^{+\infty} \phi_{n'}^{(AB)*}(x) e^{\gamma_{AB} x} \phi_n^{(AB)}(x) dx \right] (3-17)$$

$$[\int_{-\infty}^{+\infty} \phi_j^{(CD)*}(y) e^{\gamma_{CD} y} \phi_j^{(CD)}(y) dy] e^{-i(n+\frac{1}{2})\omega_{AB}t} e^{-i(j+\frac{1}{2})\omega_{CD}t}$$

where use has been made of the orthonormal property for the set of harmonic oscillator energy eigenfunctions $\phi_n^{(AB)}(x)$ and $\phi_j^{(CD)}(y)$.

$$\int_{-\infty}^{+\infty} \phi_{n'}^{(AB)*}(x) \phi_n^{(AB)}(x) dx = \delta_{nn},$$

$$\int_{-\infty}^{+\infty} \phi_{j'}^{(CD)*}(y) \phi_j^{(CD)}(y) dy = \delta_{jj},$$

Let $U_{n'n} = \int_{-\infty}^{+\infty} \phi_{n'}^{(AB)*}(x) e^{\gamma_{AB} x} \phi_n^{(AB)}(x) dx$

$$V_{j'j} = \int_{-\infty}^{+\infty} \phi_{j'}^{(CD)*}(y) e^{\gamma_{CD} y} \phi_j^{(CD)}(y) dy$$

$U_{n'n}$ and $V_{j'j}$ are the coupling terms between states, called matrix elements. Equations (3-17) can be written now as:

$$\frac{d}{dt} a_{n'j'}(t) = \left(\frac{E_0}{ih} \operatorname{sech}^2\left(\frac{u_0 t}{2L}\right) \sum_{nj} a_{nj}(t) U_{n'n} V_{j'j} e^{-i(n+\frac{1}{2})\omega_{AB}t} e^{-i(j+\frac{1}{2})\omega_{CD}t} \right) (3-18)$$

Equation (3-18) is a set of coupled first order differential equations subject to certain initial conditions. Actually, this set of differential equations is equivalent to the Schrodinger equation (3-10). Next,

let

$$a_{nj}(t) = A_{nj}(t) + iB_{nj}(t)$$

(3-19)

where $A_{nj}(t)$ is the real part of $a_{nj}(t)$, and $B_{nj}(t)$ is the imaginary part of $a_{nj}(t)$. Both $A_{nj}(t)$ and $B_{nj}(t)$ are real functions of time t . We also know that $U_{m'n}$, V_{jj} are real, and

$$e^{i(n+\frac{1}{2})\omega_{AB}t} e^{i(j+\frac{1}{2})\omega_{CD}t} = \cos [(n+\frac{1}{2})\omega_{AB}t + (j+\frac{1}{2})\omega_{CD}t] + i \sin [(n+\frac{1}{2})\omega_{AB}t + (j+\frac{1}{2})\omega_{CD}t] \quad (3-20)$$

Inserting equations (3-19) and (3-20) into equation (3-18), with some algebra, the real part and imaginary part on each side of the equation must be equal separately. So we obtain a set of coupled differential equations for A's and B's:

$$\begin{aligned} \frac{d}{dt} A_{n'nj}(t) &= \frac{E_0}{h} \operatorname{sech}^2 \left(\frac{u_0 t}{2L} \right) \sum_{nj} \{ A_{nj}(t) \sin [(n'-n)\omega_{AB}t + (j'-j)\omega_{CD}t] \\ &\quad + B_{nj}(t) \cos [(n'-n)\omega_{AB}t + (j'-j)\omega_{CD}t] \} U_{n'n} V_{j'j} \end{aligned} \quad (3-21)$$

$$\begin{aligned} \frac{d}{dt} B_{n'nj}(t) &= \frac{E_0}{h} \operatorname{sech}^2 \left(\frac{v_0 t}{2L} \right) \sum_{nj} \{ B_{nj}(t) \sin [(n'-n)\omega_{AB}t + (j'-j)\omega_{CD}t] \\ &\quad - A_{nj}(t) \cos [(n'-n)\omega_{AB}t + (j'-j)\omega_{CD}t] \} U_{n'n} V_{n'j} \end{aligned}$$

$A_{rs}(+\infty)$ and $B_{rs}(+\infty)$ are to be found, hence

$$|a_{rs}(+\infty)|^2 = |A_{rs}(+\infty)|^2 + |B_{rs}(+\infty)|^2,$$

the desired transition probability.

In Chapter 2, we mentioned that in the work of ZRS (18), they treated collinear collisions between two

diatomic molecules with symmetric configuration only.

Therefore, γ_{AB} is equal to γ_{CD} i.e.:

$$\gamma_{AB} = \gamma_{CD} = \gamma$$

The intermolecular potential function in equation (3-10), $E_0 \operatorname{sech}^2 \left(\frac{v_0 t}{2L} \right) \exp \left(\frac{\gamma_{AB} X + \gamma_{CD} Y}{L} \right)$ is expanded in a Taylor series to the second order of X and Y ;

$$E_0 \operatorname{sech}^2 \left(\frac{v_0 t}{2L} \right) \exp \left[\frac{\gamma}{L} (X+Y) \right] = E_0 \operatorname{sech}^2 \left(\frac{v_0 t}{2L} \right) \left\{ 1 + \frac{\gamma}{L} (X+Y) + \frac{1}{2} \left(\frac{\gamma}{L} \right)^2 (X+Y)^2 \right\} \quad (3-22)$$

neglecting high order terms in X, Y . Equation (3-10), with interaction potential given by equation (3-22), can be solved analytically by applying Kerner (19) method. Our semi-classical approach differs from the method of ZRS in that we do not expand the potential function $E_0 \operatorname{sech}^2 \left(\frac{v_0 t}{2L} \right) \exp \left(\frac{\gamma_{AB} X + \gamma_{CD} Y}{L} \right)$ into a Taylor series. We make a direct numerical integration of the coupled equations (3-21), which are equivalent to equation (3-10). It is appropriate to carry out classical calculations in which the approximate equation (3-8) is used, but the potential expansion (3-22) is not carried out. Numerical comparison between these two methods for N_2-N_2 collision will be given in Chapter 4.

C. Quantum Mechanical Calculation

To clarify the presentation of the theory, we shall use a scaled Schrodinger equation. Define (13):

$$y^* = \left(\frac{\mu_{CD} k_{CD}}{h^2} \right)^{\frac{1}{4}} y \quad (3-23A)$$

$$x^* = \left(\frac{\mu_{AB} k_{AB}}{h^2} \right)^{\frac{1}{4}} x \quad (3-23B)$$

Inserting equations (3-23A) and (3-23B) into equation (3-5A), it becomes:

$$\begin{aligned} & \left\{ \frac{1}{2} h \omega_{CD} \left(-\frac{\partial^2}{\partial y^*^2} + y^{*2} \right) + \frac{1}{2} h \omega_{AB} \left(-\frac{\partial^2}{\partial x^*^2} + x^{*2} \right) - \frac{h^2}{2\mu} \frac{\partial^2}{\partial \bar{R}^2} \right. \\ & \left. + E_0 \exp \left[\frac{-1}{L} \gamma_{CD} \left(\frac{h^2}{\mu_{CD} k_{CD}} \right)^{\frac{1}{4}} \left(\frac{1}{\gamma_{CD}} \left(\frac{\mu_{CD} k_{CD}}{h^2} \right)^{\frac{1}{4}} \bar{R} - y^* \right) - \frac{\gamma_{AB}}{\gamma_{CD}} \left(\frac{\mu_{CD} k_{CD}}{h^2} \right)^{\frac{1}{4}} x^* \right) \right\} \\ & \psi(x^*, y^*, \bar{R}) = E \psi(x^*, y^*, \bar{R}) \quad (3-24A) \end{aligned}$$

Dividing both sides by $1/2h \omega_{CD}$, equation (3-24) looks simpler. If we define:

$$\omega = \frac{\omega_{AB}}{\omega_{CD}}, \quad E^* = \frac{1}{2} \frac{E}{h \omega_{CD}}, \quad r^* = \frac{1}{\gamma_{CD}} \left(\frac{\mu_{CD} k_{CD}}{h^2} \right)^{\frac{1}{4}} \bar{R}$$

$$L^* = \frac{1}{L} \gamma_{CD} \left(\frac{h^2}{\mu_{CD} k_{CD}} \right)^{\frac{1}{4}} = \frac{1}{L} \frac{m_D}{m_C + m_D} \left(\frac{h^2}{\mu_{CD} k_{CD}} \right)^{\frac{1}{4}}$$

$$\frac{1}{m} = \frac{h^2}{\mu \omega_{CD}} \left(\frac{\mu_{CD} k_{CD}}{h^2} \right)^{\frac{1}{2}} \frac{1}{\gamma_{CD}^2} = \frac{\mu_{CD}}{\mu} \left(\frac{m_C + m_D}{m_D} \right)^2 = \frac{m_C m_D}{(m_A + m_B) m_D}$$

$$\beta = \frac{\gamma_{AB}}{\gamma_{CD}} \left(\frac{k_{CD} \mu_{CD}}{k_{AB} \mu_{AB}} \right)^{\frac{1}{4}} = \frac{m_C}{m_B} \left(\frac{\mu_{AB}}{\mu_{CD}} \right)^{\frac{1}{2}} \frac{1}{\omega^{1/2}}$$

The new equation (3-24) is dimensionless, namely;

$$\left\{ \left(-\frac{\partial^2}{\partial y^2} + y^{*2} \right) + \left(-\frac{\partial^2}{\partial x^2} + x^{*2} \right) - \frac{1}{\bar{m}} \frac{\partial^2}{\partial r^2} + \left(\frac{E_0}{1/2\hbar\omega_{CD}} \right) \exp(-L^*(r^*-y^*-\beta x^*)) \right\} \psi(x^*, y^*, r^*) = E^* \psi(x^*, y^*, r^*) \quad (3.24B)$$

Since the constant coefficient $\frac{E_0}{1/2\hbar\omega_{CD}}$ can be absorbed into the argument of the exponential function $\exp(-L^*(r^*-y^*-\beta x^*))$, and the operator $\frac{\partial}{\partial r^2}$ is invariant under the transformation

$r^* \longrightarrow r^* + \text{any constant}$,
 we can choose $\frac{E_0}{1/2\hbar\omega_{CD}} = 1$ with no loss of generality.
 The notation x^* , y^* and r^* refer to dimensionless variables, we now drop the ** for convenience. Hence the scaled Schrodinger equation is:

$$\left\{ \left(-\frac{\partial^2}{\partial y^2} + y^2 \right) + \omega \left(-\frac{\partial^2}{\partial x^2} + x^2 \right) - \frac{1}{\bar{m}} \frac{\partial^2}{\partial r^2} + \exp[-L^*(r-y-\beta x)] \right\} \psi(x, y, r) = E^* \psi(x, y, r) \quad (3.25)$$

The system is specified by the five dimensionless parameters. They are ω , \bar{m} , β , L^* and E^* , E^* is energy of the system in term of ground state energy of molecule CD. The dimensionless form of the unperturbated Hamiltonian H_0 is:

$$H_0(x, y) = \left(-\frac{\partial^2}{\partial y^2} + y^2 \right) + \omega \left(-\frac{\partial^2}{\partial x^2} + x^2 \right)$$

with eigenfunctions $\psi_{ij}(x, y)$ and eigenvalues W_{ij} , then

$$H_0(x, y) \psi_{ij}(x, y) = W_{ij} \psi_{ij}(x, y) \quad i = 0, 1, 2, \dots$$

$$j = 0, 1, 2, \dots$$

where

$$\phi_{ij}(x,y) = \phi_i(y) \phi_j(x)$$

$$W_{ij} = (2i+1) + (2j+1) \omega$$

$\phi_{ij}(x,y)$ is the product of the individual harmonic oscillator wavefunction which indicates that molecule AB with internal coordinate x is the vibrational state j , and molecule CD with internal coordinate y is in the state i . Let the system be in a particular initial state (n_0, m_0) . We can expand the total stationary scattering wavefunction $\psi_{n_0 m_0}(x, y, r)$, i.e. the solution to equation (3-25) in terms of $\phi_{nm}(x, y)$ because they form a complete set.

$$\psi_{n_0 m_0}(x, y, r) = \sum_{n=0}^{N-1} \sum_{m=0}^{M-1} f_{nm, n_0 m_0}(r) \phi_n(y) \phi_m(x) \quad (3-26)$$

Substituting equation (3-26) into equation (3-25) get:

$$\begin{aligned} & \sum_{n=0}^{N-1} \sum_{m=0}^{M-1} \{W_{nm} f_{nm, n_0 m_0}(r) \phi_n(y) \phi_m(x) - \frac{1}{m} \phi_n(y) \phi_m(x) \frac{d^2}{dr^2} f_{nm, n_0 m_0}(r) \\ & + \exp[-L^*(r-y-\beta x)] f_{nm, n_0 m_0}(r) \phi_n(y) \phi_m(x)\} \\ & = E^* \sum_{n=0}^{N-1} \sum_{m=0}^{M-1} f_{nm, n_0 m_0}(r) \phi_n(y) \phi_m(x) \end{aligned} \quad (3-27)$$

where we have used the relation

$$H_0(x, y) \phi_{ij}(x, y) = W_{ij} \phi_{ij}(x, y)$$

When equation (3-27) is multiplied by $\phi_{n'm}^*(y) \phi_{m'}^*(x)$ on both sides and integrated over x and y , we have:

$$\begin{aligned} & \sum_{n=0}^{N-1} \sum_{m=0}^{M-1} \{W_{nm} f_{nm, n_0 m_0}(r) \langle \phi_{n'm}, (x, y) \rangle - \frac{1}{m} \frac{d^2}{dy^2} f_{nm, n_0 m_0}(r) \langle \phi_{n'm}, (x, y) \rangle | \phi_{nm}(x, y) \rangle \\ & + f_{n_0 m_0, nm}(r) \langle \phi_{n'm}, (x, y) | \exp[-L^*(r-y-\beta x)] | \phi_{nm}(x, y) \rangle\} \end{aligned} \quad (3-28)$$

$$= E^* \sum_{n=0}^{N-1} \sum_{m=0}^{M-1} f_{nm, n_0 m_0}(r) \langle \phi_{n'm}, (x, y) | \phi_{nm}(x, y) \rangle$$

$$\begin{aligned}
 \text{where } \langle \phi_{n'm'}(x, y) | \phi_{nm}(x, y) \rangle &\equiv \int_{-\infty}^{+\infty} \phi_{n'm'}(y) \phi_{nm}(y) dy \int_{-\infty}^{+\infty} \phi_{m'}(x) \phi_m(x) dx \\
 &= \delta_{n'm'} \delta_{m'm} \tag{3-29}
 \end{aligned}$$

$$\begin{aligned}
 &\langle \phi_{n'm'}(x, y) | \exp[-L^*(x-y-\beta x)] | \phi_{nm}(x, y) \rangle \\
 &= \exp(-L^* x) \int_{-\infty}^{+\infty} \phi_{n'm'}(y) \exp(L^* y) \phi_{nm}(y) dy \int_{-\infty}^{+\infty} \phi_{m'}(x) \exp(L^* \beta x) \phi_m(x) dx
 \end{aligned} \tag{3-30}$$

$$\text{Defining: } v_{n'm',nm}(x) = \bar{m} \langle \phi_{n'm'}(x, y) | \exp[-L^*(x-y-\beta x)] | \phi_{nm}(x, y) \rangle \tag{3-30}$$

and using equation (3-29) and (3-30), equation (3-28) becomes

$$\begin{aligned}
 \bar{m} w_{n'm',n_0m_0}(x) - \frac{d^2}{dx^2} f_{n'm',n_0m_0}(x) + \sum_{n=0}^{N-1} \sum_{m=0}^{M-1} v_{n'm',nm}(x) f_{nm,n_0m_0}(x) \\
 = \bar{m} E^* f_{n'm',n_0m_0}(x)
 \end{aligned}$$

Rearranging this equation, we get

$$\begin{aligned}
 \frac{d^2}{dx^2} f_{n'm',n_0m_0}(x) &= \sum_{n=0}^{N-1} \sum_{m=0}^{M-1} v_{n'm',nm}(x) f_{nm,n_0m_0}(x) \\
 &\quad - \bar{m} (E^* - w_{n'm'}) f_{n'm',n_0m_0}(x) \tag{3-31}
 \end{aligned}$$

Define

$$\begin{aligned}
 k_{ij} &= \sqrt{\bar{m} (E^* - w_{ij})} \\
 &= \sqrt{\bar{m} (E - (2i+1) - (2j+1))} \tag{3-32}
 \end{aligned}$$

N and M are the number of states of CD and AB included in the expansion. We introduce (9)

$$\begin{aligned}
 i &= n + m.N & n &= 0, 1, 2, \dots, N-1 & m &= 0, 1, \dots, M-1 \\
 j &= n + m.N & n &= 0, 1, 2, \dots, N-1 & m &= 0, 1, \dots, M-1 \\
 k &= n + m.N & n &= 0, 1, 2, \dots, N-1 & m &= 0, 1, \dots, M-1
 \end{aligned}$$

to indicate the states of the system. For example, $i = (n, m)$, $j = (n, m)$ and $k = (n, m)$ etc. By incorporating the

definition of k , equation (3-32), we obtain a system of coupled second order ordinary differential equations in matrix form which is equivalent to the Schrodinger equation (3-32).

$$\frac{d^2}{dr^2} \vec{F}(r) = (\vec{V} - \vec{K}^2) \vec{F} \quad (3-33)$$

where $(\vec{F}(r))_{ij} = (\vec{F}(r))_{i(n,m)j(n',m')} = f_{nm,n'm'}(r)$

$$(\vec{V}(r))_{ij} = (\vec{V}(r))_{i(n,m)j(n',m')} = V_{nm,n'm'}(r)$$

and $(\vec{K})_{ij} = (\vec{K})_{i(n,m)j(n',m')} = k_{nm} \delta_{mn} \delta_{mm'}$

In the asymptotic region where r is very large, $V(r)$ tends to zero, so integrating equation (3-33), we get

$$\lim_{r \rightarrow \infty} \vec{F}(r) = e^{-i\vec{K}r} \vec{G} + e^{i\vec{K}r} \vec{J} \quad (3-34)$$

Equation (3-34) is the asymptotic form of $\vec{F}(r)$ at large r . In principle if \vec{G} and \vec{J} are determined, then the transition probability from state $i = i(n_0, m_0)$ to state $j = j(n, m)$ is given by:

$$p_{ij} = |(\vec{J} \vec{G}^{-1})_{ji}| \frac{k_j}{k_i} \quad (3-35)$$

where \vec{G}^{-1} is the inverse matrix of \vec{G} . However, it is rather difficult to find matrices \vec{G} and \vec{J} in a straightforward manner. We go to the following alternative way:

$$\frac{d}{dr} \vec{F}(r) = \vec{E}(r) \quad (3-36)$$

Set

$$\frac{d}{dr} \vec{E}(r) = (\vec{V}(r) - \vec{K}^2) \vec{F}(r)$$

Letting r_0 be some point in the asymptotic regime, the asymptotic form of $\vec{F}(r)$ may be also written as

$$\lim_{r \rightarrow \infty} \vec{F}(r) = e^{-i\vec{K}(r-r_0)} \vec{G} + e^{i\vec{K}(r-r_0)} \vec{J} \quad (3-37)$$

and then,

$$\lim_{r \rightarrow \infty} \vec{E}(r) = \lim_{r \rightarrow \infty} \frac{d\vec{F}}{dr} = -i\vec{K}e^{-i\vec{K}(r-r_0)} \vec{G} + i\vec{K}e^{i\vec{K}(r-r_0)} \vec{J}.$$

So, as the point $r = r_0$ in the asymptotic region, it is obvious that:

$$\vec{F}(r_0) = \vec{G} + \vec{J} \quad (3-38A)$$

$$\vec{E}(r_0) = -i\vec{K} \cdot \vec{G} + i\vec{K} \cdot \vec{J} \quad (3-38B)$$

These two relations will be used later. From equation (3-37) we have:

$$\lim_{r \rightarrow \infty} \vec{F}(r) \vec{G}^{-1} e^{-i\vec{K}r_0} = e^{-i\vec{K}r} + e^{i\vec{K}r} (e^{-i\vec{K}r_0} \vec{J} \vec{G}^{-1} e^{-i\vec{K}r_0}),$$

then the transition probabilities are:

$$P_{ij} = |(e^{-i\vec{K}r_0} \vec{J} \vec{G}^{-1} e^{-i\vec{K}r_0})_{ji}|^2 \frac{k_j}{k_i} \quad (3-39)$$

Since k_j is real for open channels, that is to say the incoming particle has sufficient energy to excite the bound particle to any of its lowest N eigenstates, matrix \vec{K} is real too, and these are the only observed ones. In equation (3-39) since \vec{K} is real, we end up with

$$P_{ij} = (\vec{J} \vec{G}^{-1})_{ji} \frac{k_j}{k_i}$$

This result is exactly the same as the probabilities obtained based on the asymptotic form of $\vec{F}(r)$ in equation (3-34). We can find \vec{G} and \vec{J} in terms of $\vec{F}(r_0)$, and $\vec{E}(r_0)$ with no difficult, since \vec{G} and \vec{J} are related to $\vec{F}(r_0)$, $\vec{E}(r_0)$ by equation (3-38A) and (3-38B). Solving for \vec{G} and

we obtain:

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$$2\vec{G} = \vec{F}(r_0) + i\vec{K}^{-1}\vec{E}(r_0)$$

$$2\vec{J} = \vec{F}(r_0) - i\vec{K}^{-1}\vec{E}(r_0)$$

Defining $\vec{S} = \text{Re } \vec{S} + i \cdot \text{Im } \vec{S} = \vec{J}\vec{G}^{-1}$.

where $\text{Re } \vec{S}$ is the real part of \vec{S} and $\text{Im } \vec{S}$ is the imaginary part of \vec{S} , (both are real) then

$$|\vec{J}\vec{G}^{-1}|^2 = |\vec{S}|^2 = (\text{Re } \vec{S})^2 + (\text{Im } \vec{S})^2$$

We must find $\text{Re } \vec{S}$ and $\text{Im } \vec{S}$ in terms of $\vec{F}(r_0)$, $\vec{E}(r_0)$ and \vec{K} .

We know that (19) given three matrices \vec{A} , \vec{X} , and \vec{Y} such that

$$\vec{A} = \vec{X} + i\vec{Y}$$

where \vec{X} and \vec{Y} are real, if

$$\begin{pmatrix} x & y \\ -y & x \end{pmatrix}^{-1} = \begin{pmatrix} z & w \\ -w & z \end{pmatrix}$$

then the inverse matrix of \vec{A} is

$$\vec{A}^{-1} = \vec{Z} + i\vec{W}$$

In our case $\vec{F}(r_0)$, $\vec{E}(r_0)$ and \vec{K} are all real, we may define

$$\vec{D} = \vec{K}^{-1}\vec{E}(r_0)$$

so that equation (3-40) becomes:

$$2\vec{G} = \vec{F}(r_0) + i\vec{D}$$

$$2\vec{J} = \vec{F}(r_0) - i\vec{D}$$

Here we want to find the inverse of \vec{G} , then set:

$$\begin{pmatrix} \vec{F}(r_0) & \vec{D} \\ -\vec{D} & \vec{F}(r_0) \end{pmatrix}^{-1} = \begin{pmatrix} \vec{Z} & \vec{W} \\ -\vec{W} & \vec{Z} \end{pmatrix} \quad (3-41)$$

and solve for \vec{Z} , \vec{W} in terms of $\vec{F}(r_0)$ and \vec{D} . Equation (3-41) means that

$$\begin{pmatrix} \vec{F}(r_0) & \vec{D} \\ -\vec{D} & \vec{F}(r_0) \end{pmatrix} \begin{pmatrix} \vec{Z} & \vec{W} \\ -\vec{W} & \vec{Z} \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

Or equivalently $\vec{F}(r_0)\vec{Z} - \vec{D}\vec{W} = 1$

$$\vec{F}(r_0)\vec{W} + \vec{D}\vec{Z} = 0$$

From the second equation of (3-42),

$$\vec{W} = -(\vec{F}(r_0))^{-1} \vec{D} \vec{z}$$

Substituting \vec{W} into the first equation of (3-42), get:

$$(\vec{F}(r_0) + \vec{D}(\vec{F}(r_0))^{-1} \vec{D}) \vec{z} = 1$$

hence $\vec{z} = (\vec{F}(r_0) + \vec{D}(\vec{F}(r_0))^{-1} \vec{D})^{-1}$

and $\vec{W} = -(\vec{F}(r_0))^{-1} \vec{D}(\vec{F}(r_0) + \vec{D}(\vec{F}(r_0))^{-1} \vec{D})^{-1}$

in turn, $\vec{G} = 2((\vec{F}(r_0) + \vec{D}(\vec{F}(r_0))^{-1} \vec{D})^{-1} - i(\vec{F}(r_0))^{-1} \vec{D}(\vec{F}(r_0) + \vec{D}(\vec{F}(r_0))^{-1} \vec{D})^{-1})$

The derived result is clear now:

$$\begin{aligned} \vec{G}^{-1} &= (\vec{F}(r_0) - i\vec{D})((\vec{F}(r_0) + \vec{D}(\vec{F}(r_0))^{-1} \vec{D})^{-1} - i(\vec{F}(r_0))^{-1} \vec{D}(\vec{F}(r_0) + \vec{D}(\vec{F}(r_0))^{-1} \vec{D})^{-1}) \\ &= \vec{D}(\vec{F}(r_0) + \vec{D}(\vec{F}(r_0))^{-1} \vec{D})^{-1} = \text{Re } \vec{S} + i \text{Im } \vec{S} \end{aligned} \quad (3-43)$$

Since the real part and imaginary part on each side of equation (3-43) are equal to each other respectively, we have:

$$\begin{aligned} \text{Re } \vec{S} &= (\vec{F}(r_0) - \vec{D}(\vec{F}(r_0))^{-1} \vec{D})(\vec{F}(r_0) + \vec{D}(\vec{F}(r_0))^{-1} \vec{D})^{-1} \\ \text{Im } \vec{S} &= - (2\vec{D})(\vec{F}(r_0) + \vec{D}(\vec{F}(r_0))^{-1} \vec{D})^{-1} \end{aligned} \quad (3-44)$$

where $\vec{D} = \vec{K}^{-1} \vec{E}(r_0)$.

The transition probability from state $i = (n_0, m_0)$ to state $j = (n, m)$ is:

$$P_{i(n_0, m_0) \rightarrow j(n, m)} = ((\text{Re } \vec{S})_{ji}^2 + (\text{Im } \vec{S})_{ji}^2) \frac{k_j}{k_i}$$

where $\text{Re } \vec{S}$ and $\text{Im } \vec{S}$ are given by equation (3-44). This completes the basic principle in calculating transition probability for one-dimensional scattering problem within the quantum mechanical approach. Discussions of boundary conditions in integrating equation (3-36) will be made in the next Chapter.

The original Schrodinger equation has rapidly oscillating wavelike solutions which are difficult to represent numerically. The integration of equation (3-33)

is numerically unstable, unless special algorithms are used. Secrest and Johnson, () in their exact quantum mechanical treatment of the one-dimensional scattering problem, convert the coupled differential equations into equivalent integral equations. The integral involved is then replaced by a quadrature sum. The resulting matrix equation is then solved indirectly by numerical method to obtain the transition probabilities. Chan et al. (24), propose a different numerical approach to this problem. It involves converting the set of coupled second-order equations for the translational wavefunctions into first-order equations in matrix form and then solving it by an exponential method developed by W. Magnus (25). The idea was first conceived by Light et al. (26). The method we just discussed for quantum mechanical calculation of transition probabilities is a direct integration of the state expanded Schrodinger equation. This treatment is similar to the method due to Riley and Kupermann (12). It is relatively simple and straightforward, but in our procedures, the virtual states are not included in the total wavefunction expansion. Roy G. Gordon (21) developed another method for integrating coupled differential equations arising in bound state and scattering problems in quantum mechanics. The wavefunctions are constructed in piecewise analytic form, to any prescribed accuracy. The chief advantage of this method is that it avoids searching

for the correct initial derivatives of the wavefunction.
It is claimed to be numerically very stable.

Chapter 4 Numerical Results

A. Semi-classical Results

For the specific N^2-N^2 molecular collision, the two molecules are identical, then $\omega_0 = \omega_{AB} = \omega_{CD} = 4.45 \times 10^{-14}$ sec, $\gamma_{AB} = \gamma_{CD}$, and $m_A = m_B = m_C = m_D = 14.0$ a.m.u... This implies that the matrix elements $U_{n'n}$ and $V_{n'n}$ are equal. The matrix elements $U_{n'n}$ are given in Table 4-1.

As mentioned in Chapter 3-B, we may set:

$$k = j + (n-1) J, \quad j = 1, 2, 3, \dots$$
$$n = 1, 2, 3, \dots$$

where J is an integer which is the number of states of molecules CD and AB included in the expansion of the total wavefunction ψ . The integer k is used to represent the state (j, n) which means that molecule CD is in the state j and molecule AB is in the state n . For example, if we choose $J = 4$, there are $4 \times 4 = 16$ states involved in the expansion of total wavefunction ψ in terms of the individual harmonic oscillator wavefunctions. In other words, there are $16 \times 2 = 32$ coupled first order differential equations to be solved in equation (3-21). (for J greater than 4, the extension is straightforward). In general, $J=N$, let $K=k+N.N$, where $k=1, 2, 3, \dots (N.N)$. In this way, A_{nj} and B_{nj} can be designated as:

$$Y(k) = A_{nj} \quad k = 1, 2, \dots (N.N)$$

$$Y(K) = B_{nj} \quad K = (N.N)+1, (N.N)+2, \dots 2(N.N)$$

\mathbf{Y} is then a vector of length $2(N \cdot N)$. This is the suitable form for doing numerical integration of equation (3-21). In our computer program, the IBM IMSL ROUTINE DGEAR is called. On input, $\mathbf{Y}(1), \mathbf{Y}(2), \dots, \mathbf{Y}(2N \cdot N)$ supply initial values which are initial conditions for the system. One of the arguments in the subroutine DGEAR, TOL, must be chosen suitably. Otherwise the computer time is unnecessarily long. This parameter TOL, is an estimate of the local truncation error. In a series test calculations, we choose $N=11$, $v_0 = 8$ km/sec and initial state = (1,2), for three different values of TOL, 10^{-7} , 10^{-8} , and 10^{-9} . We obtain the data as shown in Table 4-2. We then choose $TOL = 10^{-8}$. The initial value of the step size H , is chosen small enough at the beginning of integration so that it can pass the error test (based on TOL). In the subsequent procedures H is adjusted by the routine itself, but changing in the step size always satisfies the error test. The number of states used in the total wavefunction expansion, $N \cdot N$ plays a very important role in integration of equation (3-21).

In principle, we have to increase N until the final transition probabilities converge to values independent of N . Table 4-3 to Table 4-5 show the transition probability as a function of N for initial state (1,2) at low, medium, and high initial relative velocity; i.e. $v_0 = 3$ Km/sec, 6 Km/sec, and 9 Km/sec respectively. From these tables, it is obvious that for high value of velocity v_0 , we need more

states in the expansion of total wavefunction ψ . Generally speaking, for v_0 less than 6 Km/sec, $N = 7$ i.e. $7 \times 7 = 49$ states expansion is good enough for initial states (1,1), (1,2), (2,2), (3,1), and (3,2). For $v_0 = 7$ Km/sec, N should be no less than 9; and for $v_0 = 8$ Km/sec, N should be no less than 11. For $v_0 = 9$ Km/sec, N must be larger than 12. We also find that the value of N depends on the initial state. For example, for initial state (4,1) at velocity 6 Km/sec, only $N = 9$, i.e. 81 states expansion makes the transition probability converge. The integration limits are adjusted until the constraint equation (4-1) is satisfied: (for suitable N)

$$\left| \sum_{\text{final state}} P_{\text{initial state} + \text{final state}}^{-1.0} \right| < \delta \quad (4-1)$$

where δ is an arbitrary small number, we choose $\delta = 10^{-4}$ here. The integration limits depend on the initial relative velocity v_0 , e.g. for $v_0 = 3$ Km/sec, the lower limit $T = -1.6 \times 10^{-13}$ and the upper limit $T_{\text{END}} = 1.6 \times 10^{-13}$; for $v_0 = 6$ Km/sec, $T = -1.0 \times 10^{-13}$, and $T_{\text{END}} = 1.0 \times 10^{-14}$; for $v_0 = 9$ Km/sec, $T = -5.0 \times 10^{-14}$, and $T_{\text{END}} = 5.0 \times 10^{-14}$. Condition (4-1) serves as a useful criterion on numerical calculations.

The numerical results of transition probabilities for different initial and final states as a function of v_0 are obtained by the semi-classical method and shown in Tables

4-6 to 4-10. Figure 4-1 through Figure 4-5 plot the transition probabilities as a function of v_o for five initial states (1,1), (1,2), (2,2), (3,1) and (3,2). A scaling relationship can be very useful for both the theoretical and experimental analysis of molecular scattering problem. With this relationship, it is easier to write the direct Monte Carlo computation program and save computer storage space. For a systematic study of the scaling relationship, we need more data. It is too expensive to be done at this time. In Figure 4-6a and 4-6b, we plot only the scaling relationship for V-T processes (1,i) (1,i+1) at different v_o . Since the collision is symmetric, there is nothing new in the results by changing the initial state (i,j) to (j,i). We examine Figure 4-1 through Figure 4-5 and find that at low energies the probability of transferring a given number of quanta by a V-V process is much greater than the probability of converting them by a V-T process into translational energy. For V-T process, the probability increases rapidly with increasing collision energies in the low energy regime. However, the probability of V-V transfer rises less sharply with increasing collision energy. For V-V transitions involving two quanta jumps such as (3,1) (1,3) and (3,2) (1,4), the transition probability is less than that of V-V process (1,2) (2,1), which involve only one quantum jump. Generally, at low energies, the transition probabilities are very small and multiple quantum transitions are assumed

primarily due to stepwise transitions via single collision $j + j+1 + \dots + K-1 + K$. At high collision energies the direct transition $j \rightarrow K$ has significant contribution to the transition probability of multiple quantum jump. For V-V-T transfer to an adjacent level, processes involving transfer of a single quantum (such as $(2,1) \rightarrow (3,1)$ and $(2,1) \rightarrow (2,1)$) are much more probable than processes of several quanta such as $(1,2) \rightarrow (3,1)$. Likewise, transitions $(3,2) \rightarrow (4,2)$, $(3,2) \rightarrow (2,2)$ and $(3,2) \rightarrow (3,1)$ are much more probable than the transition $(3,2) \rightarrow (2,1)$. In Figure 4-6a and 4-6b we notice that the general trend of the scaling relationship for transitions

$(1,i) \rightarrow (1,i+1)$, seems to be a weak v_o -dependent function. We need more data for further analysis.

A useful check on numerical results is provided by time-reversal invariance (which leads to the principle of detailed balance). Stated classically, the principle implies that a system executes its motion in reverse if time is allowed to run backward. In quantum scattering processes this means that $P_{ij} = P_{ji}$, i.e., the probability of a transition for state i to state j is equal to that for transition for state j to i . For example, check table 4-8 and table 4-9, we have $P_{(2,2) \rightarrow (3,1)} = 0.845 \times 10^{-3}$, $P_{(3,1) \rightarrow (2,2)} = 0.843 \times 10^{-3}$, this gives

$$\frac{\left| P_{(3,1)+(2,2)} - P_{(2,2)+(3,1)} \right|}{P_{(3,1)+(2,2)}} = 0.237\%, \text{ at}$$

$$v_o = 4 \text{ Km/sec.}$$

At $v_o = 6 \text{ Km/sec}$, both $P_{(3,1)+(2,2)}$ and $P_{(2,2)+(3,1)}$ are equal to 0.121. The principle of time reversal invariance is satisfied quite well.

B. Quantum Mechanical Results and Comparison

Referring to Chapter 3-C, it is clear that the numerical procedures for calculating transition probability by the quantum mechanical method are as follows:

- (1) Integrate equation (3-36) and solve for $\vec{F}(r_o)$ and $\vec{E}(r_o)$.
- (2) Form \vec{D} and the expression $(\vec{F}(r_o) + \vec{D}(\vec{F}(r_o)))^{-1}$
 $\vec{D})^{-1}$.
- (3) Construct $\text{Re } \vec{S}$ and $\text{Im } \vec{S}$.
- (4) $P_{ij} = ((\text{Re } \vec{S})^2 + (\text{Im } \vec{S})^2)_{ji} \frac{k_j}{k_i}$

For N^2-N^2 collisions the parameters of the system are $\omega = 0.113$ (corresponding to $L = 0.2 \text{ \AA}$), $\omega = 1.0$, $\beta = 1.0$, and $\bar{m} = 0.5$. The total energy E can be assigned a suitable value which corresponds to some value of v_o . Having all these parameters, the IBM IMSL routine DGEAR is called to integrate equation (3-36) and find $\vec{F}(r_o)$, $\vec{E}(r_o)$. There are four factors in this problem that can affect the numerical integration.

- (1) Integration error.
- (2) Number of states retained in the state expansions.
- (3) Starting point of integration r_s .
- (4) End point of integration r_o .

The local truncation error is controlled by TOL , which is one of the arguments of the subroutine DGEAR. We choose $\text{TOL} = 1.0 \times 10^{-8}$ here. The starting point is chosen as the point $r = r_s$ at which the largest diagonal element of $\vec{V}(r_o)$

is equal to twice of the total energy there. The starting point is just beyond the classical turning point and in the classical forbidden region. Now, it is appropriate to discuss the initial conditions for this system. We set $N = M = 2$, so that there are four states involved in the total wavefunction expansion (for high energy collisions we need larger N and M to get more accurate results). The initial states of the molecule AB and CD are n_0 and m_0 definitely at the beginning, where $n_0 = 1, 2$; $m_0 = 1, 2$. Therefore the initial value of the matrix $\vec{F}(r)$ is a unit matrix:

$$\vec{F}_0 = \begin{matrix} 1.0 & 0 & 0 & 0 \\ 0 & 1.0 & 0 & 0 \\ 0 & 0 & 1.0 & 0 \\ 0 & 0 & 0 & 1.0 \end{matrix}$$

Since the point $r = r_s$ is in the classical forbidden region, the diagonal element of $\vec{V}(r_s)$ is much larger than the element of K , equation (3-33) becomes:

$$\frac{d^2\vec{F}}{dr^2} = \vec{V}(r_s) \cdot \vec{F}$$

The diagonal element of \vec{F} , f , satisfies the equation

$$\frac{d^2f}{dr^2} = me^{-\alpha rs} f(r) \quad \text{as } r \rightarrow \infty$$

The asymptotic solution is then $\lim_{r \rightarrow \infty} f(r) = e^{\sqrt{me^{-\alpha rs} r}}$. Recall that $E = \frac{d\vec{F}}{dr}$, this gives the initial value of matrix \vec{E} at starting point r_s as:

$$\begin{matrix}
 & g & 0 & 0 & 0 \\
 & 0 & g & 0 & 0 \\
 \vec{E}_0 = & 0 & 0 & g & 0 \\
 & 0 & 0 & 0 & g
 \end{matrix}$$

where $g = \sqrt{me^{-r_s}}$. The stopping point r_o is chosen as the point where the diagonal element of $V(r_o)$ are less than $\frac{1.0}{5000}$ of the total energy. In principle, we have to increase the total number of states N, M , until the probabilities do not change significantly. We keep $N=2$, and $M=2$ in this work and calculate transition probabilities among states $(1,1)$, $(1,2)$ and $(2,2)$. Obviously, this four-state expansion is good for low energy collisions only. Table 4-11 and Table 4-12 display the quantum transition probabilities as a function of V_o for initial states $(2,1)$ and $(1,1)$. In Table 4-13 through Table 4-16, we list the probabilities for transitions $(1,2) \rightarrow (1,1)$, $(1,2) \rightarrow (2,2)$, $(1,2) \rightarrow (2,1)$ and $(1,1) \rightarrow (1,2)$ by the four different methods. Since the computer expense is prohibitively large for fully quantum mechanical method, we do not have enough data for plotting purpose. Here, we show that the method just explored does work. The plots of these tables are thus given by Figure 4-7, Figure 4-8, Figure 4-9 and Figure 4-10 respectively.

We check these Figures (4-7) through (4-10) and find that for the three V-T transitions the probability obtained by the semi-classical method is almost one order of magnitude smaller than that obtained from ZRS analytic

method. For V-1 process $(1,2) \rightarrow (2,1)$, the semi-classical and ZRS results are very close to each other, however, the ZRS results are always slightly larger than the semi-classical results. The SSH theory is only good for low energy collisions, if the energy is too high the probability is greater than unity. This theory breaks down there. Since we are interested in high energy collision processes of two molecules. We concentrate on the semi-classical method. We believe the semi-classical treatment can supply a reasonable estimate in calculating transition probabilities. This is very helpful because the semi-classical method can save much computer time and the numerical algorithm is relatively simple.

Chapter 5 Discussions

In this final chapter, we discuss some important problems requiring further study for vibrational energy transfer.

A. Three-Dimensional Collisions

One of the assumptions which we have made in the collinear molecular collision model is that the target molecule is struck in the direction of its axis. To avoid this assumption in the collinear treatment, we have to average over the relative orientation of the molecule at the proper stage of calculations. However, the period of rotation is usually comparable with the duration of the collision, there is no simple way to take the average. A constant steric factor is generally used.

Since the rotational energy spacing is much smaller than the vibrational spacing, appreciable rotational scattering occurs over a range of molecule-molecule separations that is considerably longer than that for which vibrational transitions are important. The coupling between rotational and translation is usually strong too, so that the rotational state generally changes before the vibrational transition occurs. When a vibrational transition takes place the corresponding energy change will appear in either translational motion, or rotational motion of molecules, or both. It is obvious that if we calculate the vibrational transition probabilities,

effects of rotational motion have to be considered. In the collinear treatment however, we have assumed that the simultaneous rotational and vibrational transitions are not important and the impact parameter is zero. The realistic three-dimensional analyses that take rotational transitions into account should include the correction to the vibrational transition probability that results from the finite size of the rotational energy spacing in future work. Also, the incident particle is described by a plane wave which contains the partial waves of different orbital angular momentum (the one-dimensional model corresponding to an s-wave scattering problem). Usually, many partial waves have to be considered and this makes the problem very difficult.

B. Effect of Anharmonicity

It has been found by experience that the potential energy function of actual diatomic molecules can be represented quite accurately by a simple analytical function called Morse potential, which contains three adjustable parameters. If Morse potentials are used to describe the intramolecular forces, the diagonal matrix elements of the interaction potential which enter into the quantum theory of vibrational energy transfer are approximately but not identically equal. In the calculation given by F.H. Mies, (27) the consideration was restricted to the head-on collision between a structless incident particle and a diatomic molecule. The transition

probability is found to decrease markedly when the ratio of the diagonal elements of the initial and final oscillator states is allowed to deviate even slightly from one. The deviation in turn, increases with the anharmonicity of the molecular vibrations, and an anharmonic correction factor of the order of 10^{-1} to 10^{-2} should be applied to the generally used probability expression for atom-molecule collision. There must exist a correction factor of this kind for molecule-molecule collision.

C. Interaction Potential

Choice of a potential function to be used in calculating the transition probability is a very important task since it affects the results considerably. In the theory of inelastic molecular collisions, the scattering potential to be adopted should be simple enough to make the calculations feasible as long as the essential features of the physics of collision is not lost. This requirement is relaxed if we deal with numerical calculations. The chosen interaction potential for some pair of molecules must be relatively accurate and can be used to represent the real situation. If the intermolecular interaction is strongly orientation dependent, as in the polar gas, the molecules may take a particular orientation during the encounter. For this problem of preferential orientation, a somewhat different treatment is required.

D. Exact Classical Trajectory

In our semi-classical approximation, the classical equations are first solved to obtain the relative motion of the molecules as a function of the time. The time-dependent Schrodinger equation for the internal motion under the external perturbation is then solved to obtain the probabilities of various transitions. However, the occurrence of inelastic processes are not taken into account in solving the classical equation of motion in that the effective intermolecular potential and the effective translational energy depend on the internal state. If the incident energy is much greater than the internal energy, the influence of inelastic process on the relative motion is unimportant. It is a good approximation to ignore the internal state in calculating classical trajectory. For more rigorous calculations, an exact classical trajectory must be found in which the energy conservation law is satisfied. For high energy collisions, the semi-classical treatment is, however, a fairly good approximation. It requires less computational effort and saves much computer time.

Appendix 1 Exponential Interaction Potential

A conventional representation of the intermolecular potential energy curve is given by the Lennard-Jones 12-6 power law (28);

$$V(r) = 4\epsilon \left(\left(\frac{\sigma}{r} \right)^{12} - \left(\frac{\sigma}{r} \right)^6 \right) \quad (A1-1)$$

where $V(r)$ is the potential energy at separation r , and r is the distance between atom B and C. This is shown graphically in Figure A1-1.

is the depth of the potential well at intermolecular distance r_m , where the repulsive force $\left(\frac{\sigma}{r} \right)^{12}$ takes over the long range attractive force $\left(\frac{\sigma}{r} \right)^6$, and $V(r_m)$ is the minimum of the potential function $V(r)$. σ is the

separation at zero energy, when $V(r) = 0$, sometimes loosely called the "collision diameter". The exponential function

$$V_{INT}(r) = \text{constant} \cdot e^{-r/L} - \epsilon \quad (Al-2)$$

must be fitted to the Lennard-Jones potential $V(r)$, equation (Al-1). Here, the choice is made that the magnitudes and slopes of the potentials are set equal at $r = r_c$. r_c is the minimum value of r . These two potentials are illustrated in Figure Al-2.

Figure Al-2. The exponential potential $V_{INT}(r)$ fitted to the Lennard Jones potential $V(r)$. The magnitudes and slopes of the two potentials are set equal to each other at the classical turning point $r = r_c$.

We deduce an approximate formula:

$$L \approx \frac{\sigma}{17.5}$$

For N_2 molecule, $\sigma = 3.749 \text{ \AA}$, so $L = 0.21 \text{ \AA}$.

Appendix 2 The Choice of Initial Time Reference

Coordinate in Quantum Scattering Process

Consider the general scattering in one-dimensional space, A flux of incoming particles with mean momentum p_0 are incident from left and scattered by an arbitrary potential distribution $V(x)$ as shown in Figure A2-1, where $V(x)$ is finite and $V(x) \rightarrow 0$ as $x \rightarrow \pm\infty$

Figure A2-1. Particles scattered by an arbitrary potential.

For large and negative x , the wave packet with mean momentum p_0 can be superposed as:

$$\psi(x, t) = \int_{-\infty}^{+\infty} dp \exp(-\alpha(p-p_0)^2) \exp\left(\frac{ipx}{\hbar}\right) \exp\left(-i \frac{p^2}{2m\hbar} t\right) \quad (A2-1)$$

$$+ \int_{-\infty}^{+\infty} dp R(p) \exp(-\alpha(p-p_0)^2) \exp\left(-\frac{ipx}{\hbar}\right) \exp\left(-i \frac{p^2}{2m\hbar} t\right)$$

where $R(p)$ is the reflection coefficient which is a constant over a region $\Delta p \sim \frac{1}{\sqrt{\alpha}}$. For large and positive x , the transmitted wave packet is:

$$\psi(x, t) = \int_{-\infty}^{+\infty} dp T(p) \exp(-\alpha(p-p_0)^2) \exp\left(\frac{ipx}{\hbar}\right) \exp\left(-i\frac{p^2 t}{2m\hbar}\right) \quad (A2-2)$$

where $T(p)$ is the transmission coefficient which is constant over a region $\Delta p \sim \frac{1}{\sqrt{\alpha}}$. Since for large $|x|$ and $|t|$, the term $\exp(i(\frac{px}{\hbar} - \frac{p^2 t}{2m\hbar}))$ in equation (A2-1) and (A2-2) is a very rapidly varying function of momentum p , the integrals are essentially zero unless p , x , and t are such that the stationary phase conditions are satisfied:

$$\begin{aligned} p &\approx p_0 \\ \frac{\partial}{\partial p} \left(\frac{px}{\hbar} - \frac{p^2 t}{2m\hbar} \right) &\approx 0 \end{aligned} \quad (A2-3)$$

From equation (A2-3) we get:

$$x \approx \frac{p_0}{m} t \quad (A2-4)$$

where $\frac{p_0}{m}$ is the classical velocity. Equation (A2-4) gives the result that $t \ll 0$, if $x \ll 0$. Let's check the second term in the right hand side of equation (A2-1). Since $R(p)$ is approximately a constant within a width $\Delta p \sim \frac{1}{\sqrt{\alpha}}$ centered at $p = p_0$, then,

$$\int_{-\infty}^{+\infty} dp R(p) \exp(-(p-p_0)^2) \exp\left(-i\frac{px}{\hbar}\right) \exp\left(-i\frac{p^2 t}{2m\hbar}\right)$$

$$= R(p_0) \int_{-\infty}^{+\infty} dp \exp(-(p-p_0)^2) \exp\left(-i\left(\frac{px}{\hbar} + \frac{p^2 t}{2m\hbar}\right)\right)$$

Stationary phase conditions require that:

$$p \approx p_0$$

and

$$\frac{\partial}{\partial p} \left(\frac{px}{\hbar} + \frac{p^2 t}{2m\hbar} \right) \approx 0 \quad (A2-5)$$

Equation (A2-5) implies that $x \approx \frac{-p_0 t}{m}$, this means that reflection occurs only when $t > 0$ because the reflected wave exists only at $x \rightarrow -\infty$. Combining the discussions just made, we conclude that the incident particles hit the potential at $x \approx 0$ and $t \approx 0$, and the initial conditions of the system are described at $t \rightarrow -\infty$, the final conditions are the states at $t \rightarrow +\infty$.

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REFERENCES

1. G.A. Bird, "Shock Wave Structure in a Rigid Sphere Gas"
Academic Press, Vol. 1, P. 216, New York 1965.
2. D.I. Pullin, J.K. Harvey, and G.K. Bienkowski
"Hypersonic Leading Edge Flow of a Diatomic Gas by the
Direct Simulation Method".

APPENDIX C

COMPUTER CODE INTERNAL

```

// JOB GKB 0367425.GKBSPACE N=WATRUN2 REG=560 T=1.0 P=100 C=0
// EXEC WATFIV
//WATFIV.FT09F001 DD DISP=OLD,DSN=U.GKBSPACE.KM115
//WATFIV.SYSIN DD DATA
$JOB      BIENKOWSKI, T=59, P=100, NOLIST
C      MAIN PROGRAM FOR MONTE CARLO 3-D ENTRANCE PROBLEM CALCULATIONS
C      OBJECTIVE OF THIS MAIN PROGRAM IS TO SET THE DIMENSIONS
C      MAIN RUNNING PROGRAM IS *** RUN ***
C
C      FOLLOWING TWO CARDS HAVE TO BE ELIMINATED FOR NON IBM MACHINES
C***** ****
C      INTEGER*2 LB,NBM,NBN,NB,NBT
C      INTEGER*2 LM,LCOL
C***** ****
C
C      THE NEXT CARD IS ASSOCIATED WITH PRINCETON RANDOM NUMBER GENERATOR
C***** ****
C      COMMON/RANCOM/NRAN(4)
C***** ****
C
C      THE FOLLOWING DIMENSION STATEMENTS SET THE MAJOR ARRAY DIMENSIONS
C      AND MUST BE CONSISTENT WITH THE FOLLOWING DATA CARD -----
C
C      NSP=NUMBER OF SPECIES - EXAMPLE BELOW NSP=1
C      NJV=NUMBER OF SUBDIVISIONS OF INPUT DISTRIBUTION FUNCTION
C          EXAMPLE BELOW NJV=22
C      NMC=NUMBER OF FINAL CELLS - EXAMPLE BELOW NMC=150
C      NMP=MAX NUMBER OF MOLECULES OF EACH SPECIES ALLOWED IN PROGRAM.
C          IF EXCEEDED, PROGRAM EITHER FAILS OR RESTARTS AT BEGINNING
C          WITH NUMBER REDUCED BY 10% - EXAMPLE BELOW NMP=5000
C      NPB=MAXIMUM NUMBER IN EACH CELL - EXAMPLE NPB=150
C
C      DIMENSION DBA(1,150),NB(1,150),NBT(1,150)
C      DIMENSION TMP(1,150),TMPA(1,150),XV(1,150),XVA(1,150)
C      DIMENSION YV(1,150),YVA(1,150),ZV(1,150),ZVA(1,150),DB(1,150)
C      DIMENSION TRP(1,150),TRPA(1,150),NBM(1,150)
C      DIMENSION NBN(150),T(1,1,150)
C      DIMENSION LB(5000),LM(1,5000),ER(1,5000)
C      DIMENSION PAU(1,5000),PAV(1,5000),PAW(1,5000)
C      DIMENSION PAX(1,5000),PAY(1,5000),PAZ(1,5000),LCOL(1,5000)
C      DIMENSION FNB(150),XC(150),YC(150),ZC(150)
C      DIMENSION VEL(22,4,1),PPV(22,4,1)
C      DATA NSP/1/,NJV/22/,NMC/150/,NMP/5000/,NPB/150/

```

2 FORMAT (/17X,'NORMAL TERMINATION OF THE PROGRAM')
 NAMELIST/DIM/NSP,NJV,NMC,NMP,NPB,NRAN

INITIALIZATION OF RANDOM NUMBER GENERATOR - PRINCETON ROUTINE

```

NRAN(1)=0
NRAN(2)=0
NRAN(3)=0
NRAN(4)=0

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PRINTOUT OF MAJOR ARRAY DIMENSIONS USED ABOVE

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WRITE(6,DIM)

CALL OF MAIN OPERATING PROGRAM WHICH REQUIRES INPUTS:
 &CONTRO, &TIMES, &FLOREF, &MOLEC, &SHAPES, &GEOM, &INCUPL, &INOUT
 THESE INPUTS ARE ALL CURRENTLY IN THE NAMELIST FORMAT
 AND MAY HAVE TO BE CHANGED IF THAT CONVENTION IS NOT AVAILABLE
 BRIEF DESCRIPTION OF THE PARAMETERS FOLLOWS

&CONTRO - ONE OCCURRENCE (NEW OR RESTART)

PARAMETER	DEFAULT	DEFINITION OR EXPLANATION
NAME	8 BLANKS	ANY ALPHANUMERIC NAME UP TO 8 CHARACTERS
TITLE	24 BLNKS	ANY ALPHANUMERIC TITLE UP TO 24 CHARACTERS
PERCNT	.001	ACCURACY IN INTEGRATION PROCEDURES
ICOPY	1	NUMBER OF ADDITIONAL COPIES OF OUTPUT
DUMP	.TRUE.	IF TRUE WILL CAUSE SYSTEM DUMP FOR ANY OF 12 PROGRAMMER DESIGNED ERROR HALTS.
DEBUG (1)	.FALSE.	IF TRUE WILL PRINT MESSAGE WHEN CELL POP. EXCEEDS MNB
DEBUG (2)	.FALSE.	IF TRUE WILL PRINT CPU TIME AROUND EACH PART OF LOOP
DEBUG (3)	.TRUE.	IF TRUE WILL PRINT CPU TIME REMAINING AT END OF LOOP
NEW	.TRUE.	IF TRUE - NEW RUN - IF FALSE - RESTART OF RUN
SAVE	.FALSE.	IF TRUE - SNAPSHOT SAVED ON TAPE(9) FOR RESTART
REDO	.FALSE.	IF TRUE PROGRAM WILL AUTOMATICALLY RESTART WITH 90% OF TOTAL IF TOTAL CELL POPULATION EXCEEDS MNM

&TIMES - ONE OCCURRENCE (NEW OR RESTART)

PARAMETER	DEFAULT	DEFINITION OR EXPLANATION
DTM	---	REAL NUMBER - FRACTION OF MEAN FREE TIME PER CYCLE
ITS	---	INTEGER - NUMBER OF CYCLES PER SAMPLE
ITP	---	INTEGER - NUMBER OF CYCLES BETWEEN PRINTOUTS
TST	---	INTEGER - ESTIMATE OF NUMBER OF CYCLES TO STEADY STATE
TLIM	---	INTEGER - TOTAL NUMBER OF CYCLES TO END OF RUN - WILL TERMINATE SOONER IF CPU TIME IS TO BE EXCEEDED

&FLOREF - ONE OCCURRENCE (NEW RUN ONLY)

PARAMETER	DEFAULT	DEFINITION
LLM	---	INITIAL NUMBER OF MOLECULES LLM<NNM< OR = NMP
MNM	---	MAXIMUM NUMBER OF MOLECULES PER SPECIES
MNB	---	MAXIMUM NUMBER PER CELL - DIAGNOSTIC ONLY
MSP	---	NUMBER OF MOLECULAR SPECIES (MAX. IS 3)
MET	0	IF 0 - DATA IS IN SI (METRIC) UNITS IF >0 - DATA IS IN ENGLISH UNITS
U	---	FLOW VELOCITY (M/SEC) OR (FT/SEC)
ANGLE	---	ANGLE OF ATTACK (DEGREES)
RNU	0.0	ARRAY GIVING MOLE FRACTIONS OF SPECIES IN FREE STREAM
RMA	0.0	ARRAY GIVING MOLECULAR WEIGHTS OF SPECIES ABOVE
TF	---	FREE STREAM TEMPERATURE (K OR R)
DENF	---	FREE STREAM NUMBER DENSITY (NUM/M**3 OR NUM/FT**3)

&MOLEC - ONE OCCURRENCE (NEW RUN ONLY)

PARAMETER	DEFAULT	DEFINITION
TRF	---	REFERENCE TEMPERATURE FOR MOLECULAR DATA
DIR	0.0	CROSS-SECTIONS AT REFERENCE TEMP. (MSPXMSP)
ETA	0.0	PARAMETERS IN DIFFUSION AND VISCOSITY LAW (MSPXMSP)
PHI	0.0	PARAMETERS FOR ROTATIONAL RELAXATION (MSPXMSP)

CHI .0.0 ROTATIONAL DEGREE OF FREEDOM PARAMETER (NROT/2 - 1)
 ACR .001 ACCURACY IN MOLECULAR COLLISION CALCULATIONS

&SHAPES - ND+1 OCCURRENCES WHERE ND=NUMBER OF BODY SEGMENTS (NEW RUN)
 PARAMETER DEFAULT DEFINITION
 FIRST OCCURRENCE

BODY(I) I>3 - - NEED NOT BE SPECIFIED
 BODY(1) - - - START OF BODY (XSTART) IN ARBITRARY COORDINATE
 BODY(2) 0.0 TEMPERATURE AT FRONT OF TUBE IN K OR R
 BODY(3) 0.0 DIAMETER OF TUBE IN METERS OR FT.
 SUBSEQUENT OCCURRENCES (ND)
 BODY(1) - - - X COORDINATE FROM FRONT OF BODY OF THE DOWNSTREAM
 EDGE OF THE CURRENT BODY SEGMENT
 BODY(2) - - - TEMPERATURE AT THE BACK OF THIS BODY SEGMENT
 BODY(3) - - - SWITCH - IF NOT 0.0 THIS IS THE LAST SHAPES CARD
 BODY(I) I EVEN ALPHA - ENERGY ACCOMODATION COEFFICIENT FOR SPECIES
 BODY(J) J ODD SIGMA - TANGENTIAL ACCOMODATION COEFF. FOR SPECIES
 I AND J < (4+2*MSP)

&GEOM - ONE OCCURRENCE (NEW RUN ONLY)

PARAMETER DEFAULT DEFINITION
 NWEDG - - - INTEGER GIVING THE NUMBER OF WEDGES WITHIN 180 DEGREES
 NW - - - NUMBER OF FIRST LEVEL CELLS IN X DIRECTION
 NH - - - NUMBER OF FIRST LEVEL CELLS IN RADIAL DIRECTION

&INCPL - ONE OCCURRENCE (NEW RUN ONLY) - INPUT DISTRIBUTION

PARAMETER DEFAULT DEFINITION
 FLUXIN 1.0 FLUX INPUT IN TERMS OF FREE STREAM FLUX - ONE
 NUMBER FOR EACH SPECIES
 PCOL 0.0 FRACTION OF ARRIVING MOLECULES THAT HAVE
 PREVIOUSLY COLLIDED
 BMP - - - RATIO OF "CAVITY" PRESSURE TO THE EFFECTIVE
 PRESSURE OF THE INCOMING STREAM AT ENTRANCE
 JV - - - THE NUMBER OF VELOCITY INTERVALS FOR DISTRIBUTION
 FUNCTION INFORMATION
 KMX 4 NUMBER OF COMPONENTS OF DISTRIBUTION
 KMX=3 IF NO ROTATIONAL ENERGY (CHI=-1)
 KMX=4 IF ROTATIONAL ENERGY IS INCLUDED (CHI>-1)

&INOUT - MSP*KMX OCCURRENCES

PARAMETER DEFAULT DEFINITION
 MT DESIGNATES SPECIES
 K=1 DESIGNATES NORMAL VELOCITY
 K=2 DESIGNATES TANGENTIAL VELOCITY IN FLOW DIRECTION
 K=3 DESIGNATES TRANSVERSE TANGENTIAL VELOCITY
 K=4 DESIGNATES ROTATIONAL ENERGY
 VARG (J) - - - VELOCITY BOUNDARIES FOR DISTRIBUTION FUNCTION
 1<J<JV VELOCITY BOUNDARIES
 CURV (J) - - - PROBABILITY OF VELOCITY (OR ROTATIONAL ENERGY)
 INCIDENT AT ENTRANCE BEING BELOW VARG (J)

A SAMPLE INPUT DECK IS GIVEN BELOW:

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```

ECONTROL NAME='INTE', 'RNAL', TITLE=' PAR', 'ABOL', 'A AT', ' 95K', 'N N', 'ON. ',
  DEBUG=.F., .F., .T., NEW=.T., SAVE=.F., ICOPY=0, REDO=.T. &END
  TIMES DTM=.010, ITS=5, ITP=1000, TST=400, TLIM=1000 &END
  &FLOREP LLM=2000, MNM=5000, MNB=150, MSP=1, MET=0, U=7485.9, ANGLE=28., RNU=1., 2*0.,
    RMA=28.94, 0., 0., TF=195.51, DENP=2.52E+19 &END
  &MOLEC TRF=1000, DIR=3.5E-19, ETA=.104, PHI=0.0, CHI=-1., ACR=.001 &END
  &SHAPES BODY=0.0, 1000., .00235 &END
  &SHAPES BODY=.0025, 555., 0.0, 2*1.0 &END
  &SHAPES BODY=.0050, 345., 0.0, 2*1.0 &END
  &SHAPES BODY=.0100, 300., 0.0, 2*1.0 &END
  &SHAPES BODY=.0200, 300., 0.0, 2*1.0 &END
  &SHAPES BODY=.0300, 300., 0.0, 2*1.0 &END
  &SHAPES BODY=.0400, 300., 0.0, 2*1.0 &END
  &SHAPES BODY=.0500, 300., 0.0, 2*1.0 &END
  &SHAPES BODY=.0600, 300., 0.0, 2*1.0 &END
  &SHAPES BODY=.0700, 300., 0.0, 2*1.0 &END
  &SHAPES BODY=.0800, 300., 0.0, 2*1.0 &END
  &SHAPES BODY=.0870, 300., 1.0, 2*1.0 &END
  &GEOM NWEDG=2, NW=20, NH=3, &END

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  &INCUPL FLUXIN=2.1429, PCOL=1.0, RMP=0.0, JV=22, KMX=3 &END
  &INOUT VARG=0., 1., 2., 3., 4., 5., 6., 7., 8., 9., 10., 11., 12., 13., 14., 15., 16., 17., 18.,
    19., 20., 21., CURV=0.0, .070, .170, .282, .369, .459, .537, .599, .656, .710, .750, .785,
    .815, .845, .872, .900, .922, .951, .975, .988, .996, 1.00, &END
  &INOUT VARG=-20., -19., -17., -15., -13., -11., -9., -7., -5., -3., -1., 1., 3., 5., 7., 9.,
    11., 13., 15., 17., 19., CURV=4*0., .003, .013, .036, .084, .149, .250, .406, .611, .762,
    .871, .932, .962, .984, .995, .999, 3*1.0, &END
  &INOUT VARG=-20., -19., -17., -15., -13., -11., -9., -7., -5., -3., -1., 1., 3., 5., 7., 9.,
    11., 13., 15., 17., 19., CURV=4*0.0, .003, .013, .036, .084, .149, .250, .406, .611, .762,
    .871, .932, .962, .984, .995, .999, 3*1.0, &END

```

```

  CALL RUN (NSP, NJV, NMC, NMP, NPB, DBA, NB, NBT, TMP, TMPA, XV, XVA, XV,
  1 YVA, ZV, ZVA, T, DB, FNB, XC, YC, ZC, LM, PAU, PAV, PAW, PAX, PAY,
  2 PAZ, LCOL, TRP, TRPA, ER, LB, NBM, NBN, VEL, PFV)
  WRITE (6,2)
  STOP
  END

```

```

  SUBROUTINE RUN (NSP, NJV, NMC, NMP, NPB, DBA, NB, NBT, TMP, TMPA, XV, XVA,
  1 YV, YVA, ZV, ZVA, T, DB, FNB, XC, YC, ZC, LM, PAU, PAV, PAW, PAX, PAY, PAZ,
  2 LCOL, TRP, TRPA, ER, LB, NBM, NBN, VEL, PFV)
  MAIN RUNNING PROGRAM ** RUN *** CALLS ALL OTHER SUBROUTINES

```

```

***** ****
  INTEGER*2 LM, LCOL
  INTEGER*2 LB, NBM, NBN, NB, NBT
  INTEGER PRT, SAMP, TST, TLIM, TIME, Q
  LOGICAL DUMP, DEBUG (3), SAVE, NEW, REDO
  REAL INTGRL, LAM, MU, NU, JAY, KAY
  DIMENSION BTA(3), C1(3), C2(3), C3(3), C7(3), C8(3), DPA(3), FL(3)
  DIMENSION FDN(3), HTI(3), HTR(3), JNT(3), KNM(3), NM(3), SR(3)
  DIMENSION NAME(2), TITLE(6)
  DIMENSION RNU(3), RMA(3), WTM(3), CHI(3), DIR(3,3), DAM(3,3), PHI(3,3)
  DIMENSION ETA(3,3), CN8(3,3), CNG(3), CMG(3), CN(3,3,3), CM(3,3,3)

```

```

DIMENSION CTI(3,3),CTR(3,3),CNI(3,3),CNR(3,3),SN(3),ST(3)
DIMENSION D1(3),D2(3),D3(3),D4(3),BODY(15),DBG1(3,3),LIMIT(10)
DIMENSION COEFF(4)
DIMENSION XLIM(2),NCOL(3,3)
DIMENSION XCB(18),XS(18),YCB(18),TB(18),ALPHA(3,18),SIGMA(3,18)
DIMENSION WTS(3,18,12),UTL(3,18,12),UTT(3,18,12)
DIMENSION VTS(3,18,12),HTSI(3,18,12),HTS(3,18,12)
DIMENSION UTLI(3,18,12),UTTI(3,18,12),VTSI(3,18,12)
DIMENSION ENT(3,2),REM(3,2),FLUXIN(3),FCOL(3)
DIMENSION VARG(42),CURV(42),IPLUX(3,2)
DIMENSION LB(NMP),NBN(NMC),NBM(NSP,NMC),LM(NSP,NMP)
DIMENSION ER(NSP,NMP),TRP(NSP,NMC),TRPA(NSP,NMC)
DIMENSION DBA(NSP,NMC),NB(NSP,NMC),NBT(NSP,NMC)
DIMENSION TMP(NSP,NMC),TMPA(NSP,NMC),XV(NSP,NMC),XVA(NSP,NMC)
DIMENSION YV(NSP,NMC),YVA(NSP,NMC),ZV(NSP,NMC),ZVA(NSP,NMC)
DIMENSION T(NSP,NSP,NMC),DB(NSP,NMC)
DIMENSION FNB(NMC),XC(NMC),YC(NMC),ZC(NMC)
DIMENSION VEL(NJV,4,NSP),PPV(NJV,4,NSP)

DIMENSION PAU(NSP,NMP),PAV(NSP,NMP),PAW(NSP,NMP)
DIMENSION PAX(NSP,NMP),PAY(NSP,NMP),PAZ(NSP,NMP),LCOL(NSP,NMP)

```

RUN0460

```

COMMON /RANCOM/NRAN(4),KAWLS
COMMON /FIRST/NL,NW,NH
COMMON /SECND/BW,BH,RMP,RN,RMP
COMMON /THIRD/PI,NREG,S,SINANG,COSANG,AKN,AKT,AKN1,AKN2,AKT1,AKT2
COMMON /FORTH/NBX,RM,XR,DUMP,C9,LL(3),LLM
COMMON /PIFTH/ND,TIME,DTM,TI,ITS,ITP,TST,TLIM,RMA,RNU,DIR
COMMON /SIXTH/RMB,XSTART,INM,MNM,MNB,NEW,SAVE,PERCNT,NSR,TR
COMMON /SVNTH/LAM,MU,NU,MT,N,J,X,Y,Z,TUSE
COMMON/EIGHT/DENF,U,TF,ANGLE,TRP,CHI,PHI,ETA,WTM,DAM,VELR,XREP
NAMELIST/CONTRL/NAME,TITLE,PERCNT,ICOPY,DUMP,DEBUG,NEW,SAVE,REDO
NAMELIST/TIMES/DTM,ITS,ITP,TST,TLIM
NAMELIST/FLOREF/LLM,MNM,MNB,MSP,MET,U,ANGLE,RNU,RMA,TF,DENF
NAMELIST/MOLEC/TRP,DIR,ETA,PHI,CHI,ACR
NAMELIST/SHAPES/BODY
NAMELIST/GEOM/NWEDG,NF,NH
NAMELIST/INCUPL/FLUXIN,FCOL,RMP,JV,KMX
NAMELIST/INOUT/VARG,CURV
DATA IC/0/,ICOPY/1/
DATA DBG1/' GAS',' AT ','110 ','FLOW',' AT ','130 ',' RUN',' AT ',RUN0630
1*303 '/RUN0640
DATA LIMIT/12,9,18,500,3600,70,900,3,20,3/
DATA TITLE/' ',' ',' ',' ',' ',' ',' ',' ',' '/RUN0660
DATA NAME/' ',' ',' '/RUN0670
DATA CPC/0.0/,CPM/0.0/,CPB/0.0/,CPJ/15.0/

```

RUN0680

*****RUN0690

RUN0700

RUN0710

RUN0720

RUN0730

RUN0740

RUN0750

FORMATS

1 FORMAT(1H1)

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2 FORMAT(1H1/17X,'RARIFIED SUPERSONIC FLOW OF BINARY GAS',T74,'I') RUN0760
3 FORMAT('+',103X,'COPY ',I2) RUN0770
4 FORMAT(/17X,'FLOW THROUGH ALL THE BOUNDARIES'/5X,'MT',5X,'MASS FLOW
 1E FR.',10X,'FCOL. FLUXIN FLUXES (ENT)')
5 FORMAT(3X,I4,F9.2,F9.4,5X,4P10.4)
25 FORMAT(/5X,'PRESSURE RATIO(INSIDE/ENTRANCE) - EITHER TYPE =',F13.5
 1,' RMP',T76,'I'/6X,'DENSITY RATIO(INSIDE/ENTRANCE) - EITHER TYPE
 2=',F13.5,' RMN',T76,'I'/9X,'FLUX RATIO(INSIDE/ENTRANCE) - EITHER
 3TYPE =',F13.5,' RMP',T76,'I')
26 FORMAT(/10X,'FLUX RATIOS FOR SPECIES RMA =',F7.2/2X,'BOUNDARY',
 1' INFLUX REFLUX NET FLUX NET FLUX/RHO*U'/2X,'ENTRANCE',
 24P10.4/4X,'CAVITY',4F10.4/)
30 FORMAT('1TIME =',F6.3,60X,'RANDOM NUMBER GENERATOR HAS BEEN CALLED
 1 ',I'0,', TIMES') RUN0800
31 FORMAT(' CPU TIME LEFT- ',F8.3) RUN0810
32 FORMAT(7X,'-MOLECULES-/3X,3I6)
33 FORMAT(' TIME = ',F8.3,5X,'COLLISION LOOP=',F8.3,5X,'MOVE LOOP = '
 1,F8.3,5X,'TOTAL TIME = ',F8.3/21X,'2ND MOVE LOOP =',F8.3,5X,
 2'CLEANUP LOOP=',F8.3,4X,'PARTICLE NUMBERS = ',4I6)
34 FORMAT(9X,'-MOLECULAR COLLISIONS-/3(3I14/))
35 FORMAT(2X,'-COLLISIONS WITH SURFACE-/3X,3I8)
36 FORMAT(' MAXIMUM NUMBER OF MOLECULES SO FAR- ',I6//) RUN0880
38 FORMAT(' EXCESS MOLECULES OCCURRED IN ',3A4) RUN0890
40 FORMAT(/' SOMETHING IS WRONG WITH BOX NUMBERING IN RUN ',/9I5,5E14,RUN0900
 15)
44 FORMAT(' NB(',I2,',',',',I4,',') POPULATION EXCEEDED ',I3,', IN MAIN AT TRUN0930
 1IME = ',F7.3) RUN0940
50 FORMAT(/' SNAP SAVED ON TAPE') RUN0950
RUN0960
***** RUN0970
RUN0980

```

CPA=ELTIME(0)

CALL NOUNDF

CALL TRAPS(0,1,1000000,1,1)

LIMIT(4)=NMC

LIMIT(5)=NMP

LIMIT(6)=NPB

LIMIT(7)=NJV

LIMIT(10)=NSP

KAWLS=0

PI=3.141593

PIROOT=SQRT(PI)

MET=0

LARGE=0

NL=1

DUMP=.TRUE.

DEBUG(1)=.FALSE.

DEBUG(2)=.FALSE.

DEBUG(3)=.TRUE.

SAVE=.FALSE.

NEW=.TRUE.

REDO=.FALSE.

PERCNT=.001

ACR=.001

DO 58 I=1,15

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RUN1030

RUN1040

RUN1060

RUN1080

RUN1190

RUN1230

RUN1260

RUN1240

RUN1250

```

58 BODY(I)=0.0
  DO 60 I=1,3
  RNU(I)=0.0
  RMA(I)=0.0
  CHI(I)=0.0
  PLUXIN(I)=0.0
  PCOL(I)=0.0
  LL(I)=0
  DO 59 J=1,18
  ALPHA(I,J)=1.0
59 SIGMA(I,J)=1.0
  DO 60 K=1,3
  ETA(I,K)=0.0
  PHI(I,K)=0.0
60 DIR(I,K)=0.0
  WRITE(6,1)                                     RUN1330
  READ(5,CONTRL)                                RUN1340
  WRITE(6,CONTRL)
  IF(NEW) GO TO 103                             RUN1360
  REWIND 9                                     RUN1370
  READ(9) DENP,U,XREP,TRP,KAWLS,NL,NW,NH,BW,BH,NREG,XLB,XLC,PI,ND,
2   S,SINANG,COSANG,AKN,AKT,NBX,BM,XR,TIME,DTM,TI,ITS,ITP,TST,
3   TLIM,RMA,RNU,DIR,XSTART,MNN,MNB,TR,BZC,CN7,DRP,PCF,FNA,
4   HTF,INM,LLM,NAV,NMAX,NWEDG,PRT,SAMP,AKN1,AKN2,AKT1,AKT2,
5   BTA,C1,C2,C3,C7,C8,DAM,DPA,FL,DELANG,PDN,HTI,HTR,JNT,KNM,
6   NM,WTM,C4,VRM,NCOL,CTI,CTR,CNI,CNR,SN,
7   ST,D1,D2,D3,D4,NRAN,VELR,RMP,RMN,RMF,IFLUX,FLUXIN,
8   XLIM,COEFF,XCB,XS,YCB,TB,ALPHA,SIGMA,NTS,
9   UTL,UTT,VTS,HTS,HTSI,ENT,REM,TMPA,
A   DBA,NB,NBT,TMP,XV,XVA,YV,YVA,ZV,ZVA,T,DB,FNB,XC,YC,ZC,
B   PAU,PAV,PAW,PAX,PAY,PAZ,LCOL,LN,
C   ETA,PHI,CHI,CN,CM,CNG,CMG,CN8,TRP,TRPA,MSP,ANGLE,TP,
D   UTILITY,UTTI,VTSI,ER,RMB,LB,NBM,NBN,VEL,PFV,PCOL,JV
  DTMO=DTM
  READ(5,TIMES)
  WRITE(6,TIMES)
  IF(DTM.EQ.DTMO) GO TO 100
  AIME=TIME*DTMO
  TIME=AIME/DTM+0.1
  DO 99 J=1,NSP
  DO 99 L=1,2
  ENT(J,L)=ENT(J,L)*DTM/DTMO
99 CONTINUE
100 IF(TI.GT.0.0) TST=TI/DTM
  WRITE(6,2)
  WRITE(6,4)
  WRITE(6,5) (MT,RMA(MT),RNU(MT),PCOL(MT),PLUXIN(MT),(ENT(MT,K),K=1,2
1),MT=1,MSP)
  WRITE(6,2)                                     RUN1540
  CALL PRINTA(NWEDG,TITLE,NAME,XCB,YCB,TB,ALPHA,SIGMA,XLIM,
1COEFF,LIMIT,MSP)
  CALL PRINTB(FNA,MSP,FNB,NM,XLIM,XC,YC,ZC,NB,NSP)
  GO TO 280                                     RUN1590
103 READ(5,TIMES)
  WRITE(6,TIMES)

```

```

READ(5,FLOREF)
WRITE(6,FLOREF)
READ(5,MOLEC)
WRITE(6,MOLEC)
IF(MSP.GT.LIMIT(10)) CALL DIAG(10,LIMIT(10),MSP)
CHIM=0.0
RMR=0.0
DMR=0.0
ETT=0.0
DO 105 M=1,MSP
BMR=RME+RMA(M)*RNU(M)
CHIM=CHIM+CHI(M)*RNU(M)
105 CONTINUE
DO 115 K=1,MSP
DO 115 M=1,MSP
ETT=ETT+RNU(M)*RNU(K)*ETA(M,K)
SRMA=SQRT(.5*RMR*(1./RMA(M)+1./RMA(K)))
115 DMR=DMR+RNU(M)*RNU(K)*DIR(M,K)*(TRF/TP)**(ETA(M,K)/2.)*SRMA
XREF=1./(DENP*DMR*1.414214)
VELR=SQRT(16628.64*TP/RMR)
IF(MET.NE.0) VELR=SQRT(99437.92*TP/RMR)
TMR=XREF/VELR
S=U/VELR
NREG=1
ND=1
READ(5,SHAPES)
WRITE(6,SHAPES)
XCB(1)=BODY(1)/XREF
TB(1)=BODY(2)/TP
XLIM(1)=XCB(1)
XSTART=XLIM(1)
RMB=.5*BODY(3)/XREF
TR=TB(1)
104 READ(5,SHAPES)
WRITE(6,SHAPES)
ND=ND+1
IF(ND.GT.LIMIT(3)) CALL DIAG(3,LIMIT(3),ND)
XCB(ND)=BODY(1)/XREF
TB(ND)=BODY(2)/TP
DO 1104 M=1,MSP
ALPHA(M,ND)=BODY(2+2*M)
1104 SIGMA(M,ND)=BODY(3+2*M)
IF(TB(ND).GT.TR) TR=TB(ND)
IF(BODY(3).EQ.0.0) GO TO 104
NSTEP=NREG+1
XLIM(NSTEP)=XCB(ND)
COEFF(1)=0.0
COEFF(2)=1.0
COEFF(3)=0.0
COEFF(4)=-RMB**2
XR=XLIM(NSTEP)-XSTART
AKN=1./XR
AKT=.5/RMB
RMFP1=TB(1)**(.5+ETT/2.)/(2.*PIROOT*S)
RMFP2=TB(ND)**(.5+ETT/2.)/(2.*PIROOT*S)

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RUN1640
RUN1650

RUN1670
RUN1680
RUN1690

RUN1770
RUN1800

```

AKN1=AKN*RMPP1
AKN2=AKN*RMPP2
AKT1=AKT*RMPP1
AKT2=AKT*RMPP2
DO 260 N=1,3
DO 260 M=1,3
NCOL(N,M)=0
260 CONTINUE
READ(5,GEOM)
READ(5,INCULP)
WRITE(6,GEOM)
WRITE(6,INCULP)
DO 116 J=1,42
VARG(J)=0.0
116 CURV(J)=0.0
RMN=RMP*TB(1)/TB(ND)
RMP=SQRT(RMN*RMP)
RM=RMB
BW=XR/NW
BH=RM/NH
IF(NWEDG.GT.LIMIT(1)) CALL DIAG(1,LIMIT(1),NWEDG)
IF(MNM.GT.LIMIT(5)) CALL DIAG(5,LIMIT(5),MNM)
IF(MNB.GT.LIMIT(6)) CALL DIAG(6,LIMIT(6),MNB)
DELANG=180./NWEDG
SINANG=SIN(ANGLE/180.*PI)
COSANG=COS(ANGLE/180.*PI)
VOL=PI*RM*RM*XR
NBI=NW*NH*NWEDG
IF(NBX.GT.LIMIT(4)) CALL DIAG(4,LIMIT(4),NBX)
IF(JV.NE.LIMIT(7)) CALL DIAG(7,LIMIT(7),JV)
BR=SQRT(TR)
SRMX=0.0
DO 916 MT=1,MSP
WTM(MT)=RMA(MT)/RMR
BTA(MT)=SQRT(WTM(MT))
SR(MT)=S*BTA(MT)
SRT=SR(MT)*FLUXIN(MT)/RNU(MT)
IF(SRT.GT.SRMX) SRMX=SRT
916 CONTINUE
INM=LLM*SQRT(TB(1))/PIROOT/SRMX/(1.+RMN)
DDN=INM/VOL
DO 140 MT=1,MSP
FDN(MT)=RNU(MT)*DDN
DFA(MT)=RNU(MT)
SN(MT)=SR(MT)*COSANG
ST(MT)=SR(MT)*SINANG
DO 117 K=1,MSP
DAM(K,MT)=DIR(K,MT)*(TRF/TP)**(ETA(K,MT)/2.)/DMR
CN8(K,MT)=DDN/DAM(K,MT)*1.414214
BT=AMIN1(BTA(K),BTA(MT))
VR1=S+3.* (1.+SQRT(TR))/BT
VR2=3.*SQRT((1.+2.*S**2/(5.+CHIM))* (1./WTM(K)+1./WTM(MT)))
CM(K,MT,1)=AMAX1(VR1,VR2)
CN(K,MT,1)=RAND(0)*CM(K,MT,1)
DF=PHI(K,MT)*(CHI(K)+CHI(MT)+2.0)-1

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RUN2200

RUN2230

RUN2240

RUN2320

RUN2330

RUN2420

RUN2470

RUN2480

RUN2740

RUN2750

```

DS=PHI (K,MT) * (2. - .5*ETA (K,MT)) - 1.0
DO 917 N=2,3
XPM=ACR**AMIN1 (DF, DS)
IF ((DF.GT.0.) .AND. (DS.GT.0.)) XPM=(DF/(DF+DS))**DF* (DS/(DF+DS))**DS
XPN=ACR**AMAX1 (DF, DS)
IF ((DF.LT.0.) .AND. (DS.LT.0.)) XPN=(DF/(DF+DS))**DF* (DS/(DF+DS))**DS
CM (K,MT,N)=XPM-XPN
CN (K,MT,N)=RAND (0)*CM (K,MT,N)
DF=CHI (K)
DS=CHI (MT)
917 CONTINUE
117 CONTINUE
DO 118 K=1,KMX
READ (5,INOUT)
WRITE (6,INOUT)
DO 118 J=1,JV
VEL (J,K,MT)=VARG (J)
PFV (J,K,MT)=CURV (J)
118 CONTINUE
ENT (MT,1)=INM*S*DTM*AKN*FLUXIN (MT)/RNU (MT)
ENT (MT,2)=RMN*ENT (MT,1)*SQRT (TB (ND)/TB (1))
REM (MT,1)=0.0
REM (MT,2)=0.0
LL (MT)=INM*PIROOT*(1.+RMN)*SR (MT)/SQRT (TB (1))*FLUXIN (MT)/RNU (MT)
CHT=CHI (MT)
IF (CHT.GT.0.) CMG (MT)=CHT**CHT*EXP (-CHT)
IF (CHT.EQ.0.) CMG (MT)=1.0
IF (CHT.LT.0.) CMG (MT)=ACR**CHT*EXP (-ACR)
CNG (MT)=RAND (0)*CMG (MT)
140 CONTINUE
XS (1)=0.0
DO 155 N=2,ND
155 XS (N)= (.5* (XCB (N) + XCB (N-1)) - XSTART) *AKN
YCB (1)=0.0
DO 160 N=2,ND
160 YCB (N)=2.* (XCB (N) - XCB (N-1)) /RMB
CALL CELL (BW,BH,NW,NH,XSTART,DELANG,NWEDG,IC,YC,ZC,FNB)
PNA=0.0
DO 210 N=1,NBX
210 PNA=PNA+PNB (N)
NPX=NBX
220 TIME=0
LARGE=0
SAMP=0
PRT=0
NAV=0
AIME=0.
TI=0.0
DT=DTM
NMAX=0
DO 250 MT=1,3
C1 (MT)=RAND (0)
C2 (MT)=RAND (0)
C3 (MT)=RAND (0)
C7 (MT)=RAND (0)

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RUN330

RUN3380

RUN3390

RUN4250

RUN4260

RUN4270

RUN4290

RUN4300

RUN4310

RUN4380

RUN4410

RUN4420

RUN4430

RUN4440

C8(MT)=RAND(0)	RUN4450
D1(MT)=RAND(0)	RUN4460
D2(MT)=RAND(0)	RUN4470
D3(MT)=RAND(0)	RUN4480
D4(MT)=RAND(0)	RUN4490
FL(MT)=0.	RUN4500
HTI(MT)=0.	RUN4510
HTR(MT)=0.	RUN4520
JNT(MT)=0	RUN4530
NM(MT)=0	RUN4540
IPLUX(MT,1)=0	
IPLUX(MT,2)=0	
DO 230 N=1,3	RUN4550
CTI(MT,N)=0.	RUN4560
CTR(MT,N)=0.	RUN4570
CNI(MT,N)=0.	RUN4580
230 CNR(MT,N)=0.	RUN4590
DO 240 N=1,ND	RUN4600
DO 240 K=1,NWEDG	RUN4610
NTS(MT,N,K)=0	RUN4620
HTSI(MT,N,K)=0.	RUN4640
UTLI(MT,N,K)=0.	
UTTI(MT,N,K)=0.	
VTSI(MT,N,K)=0.	
UTL(MT,N,K)=0.	RUN4650
UTT(MT,N,K)=0.	RUN4660
VTS(MT,N,K)=0.	RUN4670
240 HTS(MT,N,K)=0.	RUN4680
DO 245 N=1,NPX	RUN4690
NB(MT,N)=0	RUN4700
NBT(MT,N)=0	RUN4720
DBA(MT,N)=0.	RUN4730
XVA(MT,N)=0.	RUN4740
YVA(MT,N)=0.	RUN4750
ZVA(MT,N)=0.	RUN4760
TMPA(MT,N)=0.	RUN4790
TRPA(MT,N)=0.0	
DO 245 NN=1,NSP	
T(MT,NN,N)=0.0	
245 CONTINUE	RUN4800
250 CONTINUE	
FND=DDN	
DRF=2./(FND*S*S*RMB*RMB*PI)	RUN4930
PCP=1./(FND*S*RMB*RMB*PI)	RUN4940
HTF=.5*DRF/S	
WRITE(6,2)	
WRITE(6,4)	
WRITE(6,5) (MT,RMA(MT),RNU(MT),FCOL(MT),FLUXIN(MT),(ENT(MT,K),K=1,2	
1),MT=1,MSP)	
WRITE(6,2)	RUN5000
CALL PRINTA(NWEDG,TITLE,NAME,XCB,YCB,TB,ALPHA,SIGMA,XLIM,	
1COEFF,LIMIT,MSP)	
CALL GAS(NWEDG,DELANG,ND,BTA,C1,DFA,NM,FNB,DB,NB,NBM,NBN,	
1PAU,PAV,PAW,PAX,PAY,PAZ,XLIM,COEFF,LM,LIMIT(4),LIMIT(6),XCB,TB,	
2LARGE,MNN,MNB,DEBUG(1),LCOL,MSP,ER,CHI,CNG,CMG,NSP,LB)	

CPUTYM=TPIND(0)
 IF(LARGE.NE.0) GO TO 345
 DO 265 I=1,MSP
 265 IF(NM(I).GT.NMAX) NMAX=NM(I)
 CALL PRINTB(FNA,MSP,FNB,NM,XLIM,XC,YC,ZC,NB,NSP)
 IF(DEBUG(2)) WRITE(6,1)
 CALL ACCUM(NMC,NPB,FNB,NB,PAU,PAV,PAW,ER,TMP,TRP,XV,YV,ZV,LH,MSP,
 1NSP,NBM)
 CPA=ELTIME(0)
 CPI=CPA
 GO TO 340
 280 TIME=TIME+1
 IF(TIME.NE.TST+1) GO TO 285
 TI=TST*DTM
 DO 282 MT=1,MSP
 282 IFIUX(MT,1)=0
 282 IFIUX(MT,2)=0
 285 LARGE=0
 CPI=ELTIME(0)
 AIME=TIME*DTM
 DT=AIME-TI
 IF(DEBUG(1)) WRITE(6,33) AIME,CPC,CPM,CPI,CPB,CPA, (NM(I),I=1,3),NMAX
 PRT=PRT+1
 SAMP=SAMP+1
 CALL COLIDE(CN,CM,WTM,DB,DBA,NB,NCOL,LCOL,PAU,PAV,PAW,ER,T,LH,MSP,
 1LIMIT(4),LIMIT(6),ETA,PHI,CHI,CN8,NSP,NBM)
 KNM(1)=0
 KNM(2)=0
 KNM(3)=0
 CPC=ELTIME(0)
 IF(DEBUG(1)) WRITE(6,33) AIME,CPC,CPM,CPI,CPB,CPA, (NM(I),I=1,3),NMAX
 CALL MOVE(0,AKN,NWEDG,XSTART,LIMIT(3),LIMIT(1),LIMIT(8),LIMIT(9),
 1DELANG,BTA,C2,C3,DFA,FL,HTI,HTR,JNT,KNM,NM,XCB,XLIM,CTI,CTR,
 2CNI,CNR,ALPHA,SIGMA,COEFF,HTS,HTSI,NTS,UTL,UTT,VTS,PAU,PAV,PAW,
 3PAX,PAY,PAZ,LCOL,TB,MSP,ER,CHI,CNG,CMG,NSP,UTLI,UTTI,VTSI,IFLUX)
 KNM(1)=NM(1)
 KNM(2)=NM(2)
 KNM(3)=NM(3)
 CPM=ELTIME(0)
 IF(DEBUG(1)) WRITE(6,33) AIME,CPC,CPM,CPI,CPB,CPA, (NM(I),I=1,3),NMAX
 TBI=TB(ND)
 CALL FLOW(NWEDG,MNM,LARGE,BTA,C1,C7,C8,ENT,REM,LCOL,MSP,NM,SN,ST,
 1TBI,PAU,PAV,PAW,PAX,PAY,PAZ,ER,CHI,CNG,CMG,NSP,JV,FCOL,VEL,PFV)
 IF(LARGE.NE.0) GO TO 345
 CPB=ELTIME(0)
 IF(DEBUG(1)) WRITE(6,33) AIME,CPC,CPM,CPI,CPB,CPA, (NM(I),I=1,3),NMAX
 CALL MOVE(1,AKN,NWEDG,XSTART,LIMIT(3),LIMIT(1),LIMIT(8),LIMIT(9),
 1DELANG,BTA,C2,C3,DFA,FL,HTI,HTR,JNT,KNM,NM,XCB,XLIM,CTI,CTR,
 2CNI,CNR,ALPHA,SIGMA,COEFF,HTS,HTSI,NTS,UTL,UTT,VTS,PAU,PAV,PAW,
 3PAX,PAY,PAZ,LCOL,TB,MSP,ER,CHI,CNG,CMG,NSP,UTLI,UTTI,VTSI,IFLUX)
 CPB=CPB+ELTIME(0)
 IF(DEBUG(1)) WRITE(6,33) AIME,CPC,CPM,CPI,CPB,CPA, (NM(I),I=1,3),NMAX
 DO 330 MT=1,MSP
 DO 290 N=1,NBX
 NB(MT,N)=0

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RUN5060

RUN5130

RUN5140

RUN5170

RUN5180

RUN5190

RUN5220

RUN5230

RUN5280

RUN5290

RUN5330

RUN5420

RUN5450

290 CONTINUE RUN5470
 NG=NM(MT)
 N=0 RUN5480
 295 N=N+1 RUN5490
 IF(N.GT.NG) GO TO 310 RUN5500
 X=PAX(MT,N)
 Y=PAY(MT,N)
 Z=PAZ(MT,N)
 R=SQRT(Y*Y+Z*Z)
 TANG=180.*ATAN2(Z,-Y)/PI
 IWDGE=TANG/DELANG+1
 IF(IWDGE.LT.1) IWDGE=1
 IF(IWDGE.GT.NWEDG) IWDGE=NWEDG
 L=X/BW
 IF(L.GE.NW) L=NW-1 RUN5610
 M=R/BH
 IF(M.GE.NH) M=NH-1 RUN5620
 K=NWEDG*(L*NH+M)+IWDGE
 IF(K.LE.NBX) GO TO 305 RUN5630
 WRITE(6,40)L,M,IWDGE,NW,NH,NWEDG,K,MT,N,X,Y,Z,R,TANG
 IF(DUMP) CALL ABEND(4)
 STOP RUN5640
 305 J=NB(MT,K)+1 RUN5660
 IF(J.LE.MNB) GO TO 308
 IF(DEBUG(1)) WRITE(6,44) MT,K,MNB,AIME
 308 NB(MT,K)=J RUN6400
 LB(N)=K
 GO TO 295 RUN6430
 310 CONTINUE
 NBM(MT,1)=0
 DO 320 M=1,NBX
 A=NB(MT,M)
 DB(MT,M)=A*DFA(MT)/PNB(M)
 NBM(MT,M+1)=NBM(MT,M)+NB(MT,M)
 NBN(M)=NBM(MT,M)
 320 CONTINUE RUN6450
 IF(NM(MT).GT.NMAX) NMAX=NM(MT)
 DO 325 N=1,NG
 Q=LB(N)
 NBN(Q)=NBN(Q)+1
 NA=NBN(Q)
 325 LM(MT,NA)=N
 330 CONTINUE RUN6510
 IF(SAMP.LT.ITS) GO TO 335
 CALL ACCUM(NMC,NPB,PNB,NB,PAU,PAV,PAW,ER,TMP,TRP,XV,YV,ZV,LM,MSP,
 1NSP,NBM)
 SAMP=0 RUN6570
 IF(TIME.LE.TST) GO TO 335
 CALL AVRGE(PNB,DB,DBA,NB,NBT,XV,YV,ZV,XVA,YVA,ZVA,TMP,TMPA,TRP,TRP
 1A,MSP,NSP)
 NAV=NAV+1
 335 CPA=ELTIME(0)
 CPI=CPC+CPM+CPB+CPA
 CPJ=2.*CPI+5.
 340 CPUTYM=TFIND(0) RUN6600

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IF(DEBUG(3)) WRITE(6,33) AIME,CPC,CPM,CPI,CPB,CPA, (NM(I),I=1,3),NMAX
IF(TIME.EQ.0) GO TO 355
IF((TIME.GE.TLIM).OR.(CPUTYM.LE.CPJ)) GO TO 345
IF(PRT.LT.ITP) GO TO 280
PRT=0
345 WRITE(6,30) AIME,KAWLS
IF(DEBUG(3)) WRITE(6,31) CPUTYM
WRITE(6,32) (NM(I),I=1,3)
WRITE(6,34) ((NCOL(I,J),J=1,3),I=1,3)
WRITE(6,35) (JNT(I),I=1,3)
IF(LARGE.NE.0) GO TO 360
WRITE(6,36) NMAX
IF(.NOT.SAVE) GO TO 355
IF(PRT.NE.0.AND.CPUTYM.GT.CPJ.AND.TIME.LT.TLIM) GO TO 355
REWIND 9
WRITE(9) DENP,U,XREF,TRP,KAWLS,NL,NW,NH,BW,BH,NREG,XLB,XLC,PI,ND,
2      S,SINANG,COSANG,AKN,AKT,NBX,RM,XR,TIME,DTM,TI,ITS,ITP,TST,
3      TLIM,RMA,RNU,DIR,XSTART,MNM,MNB,TR,BZC,CN7,DRF,FCF,FNA,
4      HTF,INM,LLM,NAV,NMAX,NWEDG,PRT,SAMP,AKN1,AKN2,AKT1,AKT2,
5      BTA,C1,C2,C3,C7,C8,DAM,DFA,FL,DELANG,FDN,HTI,HTR,JNT,KNN,
6      NM,WTM,C4,VRM,NCOL,CTI,CTR,CNI,CNR,SN,
7      ST,D1,D2,D3,D4,NRAN,VELR,RMP,RMN,RMF,IPLUX,FLUXIN,
8      KLIM,COEFF,XCB,XS,YCB,TB,ALPHA,SIGMA,NTS,
9      UTL,UTT,VTS,HTS,HTSI,ENT,REM,TMPA,
A      DBA,NB,NBT,TMP,XV,XVA,YV,YVA,ZV,ZVA,T,DB,FNB,XC,YC,ZC,
B      PAU,PAV,PAW,PAX,PAY,PAZ,LCOL,LH,
C      ETA,PHI,CHI,CN,CM,CNG,CMG,CN8,TRP,TRPA,MSP,ANGLE,TF,
D      UTLI,UTTI,VTSI,ER,RMB,LB,NBM,NBN,VEL,PPV,FCOL,JV
      WRITE(6,50) RUN7050
355 CONTINUE
      WRITE(6,25) RMP,RMN,RMF
      DO 356 MT=1,MSP
      FIN1=ENT(MT,1)
      FIN2=ENT(MT,2)
      RF1=IPLUX(MT,1)*DTM/DT
      RF2=IPLUX(MT,2)*DTM/DT
      RNF1=1.-RF1/FIN1
      RNF2=(FIN2-RF2)/FIN1
      RPS1=RNF1*FLUXIN(MT)/RNU(MT)
      RPS2=RNF2*FLUXIN(MT)/RNU(MT)
356 WRITE(6,26) RMA(MT),FIN1,RF1,RNF1,RPS1,FIN2,RF2,RNF2,RPS2
      IF(TIME.EQ.0) GO TO 280
      IF(TIME.LE.TST) GO TO 350
      CALL PRINT1(DT,COSANG,SINANG,RMA,RNU,DRF,FCF,HTF,FL,HTI,HTR,CTI,
1CTR,CNI,CNR)
      CALL PRINT2(AKN,XSTART,DT,RNU,RMA,DRF,FCF,HTF,UTLI,UTTI,VTSI,HTSI,
1DELANG,NWEDG,XS,XCB,YCB,HTS,NTS,UTL,UTT,VTS,LIMIT(3),LIMIT(1),MSP)
      CALL PRINT4(MSP,CHI,RNU,NSP,TRPA,FDN,WTM,DB,NBT,TMPA,XVA,
1YVA,ZVA,1,NBT,XC,YC,ZC)
      GO TO 353
350 CONTINUE
      CALL PRINT4(MSP,CHI,RNU,NSP,TRP,FDN,WTM,DB,NB,TMP,XV,YV,ZV,
10,NB,XC,YC,ZC)
353 IF(DEBUG(2)) WRITE(6,1)
      IF((TIME.LT.TLIM).AND.(CPUTYM.GT.CPJ)) GO TO 280

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IP(IC.EQ.1COPY) RETURN
IC=IC+1
WRITE(6,2)
WRITE(6,4)
WRITE(6,5) (MT,RMA(MT),RNU(MT),FCOL(MT),FLUXIN(MT),(ENT(MT,K),K=1,2
1),MT=1,MSP)
WRITE(6,2)
WRITE(6,3) IC
CALL PRINTA(NWEDG,TITLE,NAME,XCB,YCB,TB,ALPHA,SIGMA,XLIM,
1COEFF,LIMIT,MSP)
CALL PRINTB(FNA,MSP,PNB,NM,XLIM,XC,YC,ZC,NB,NSP)
SAVE=.FALSE.
GO TO 345
360 WRITE(6,38) (DBG1(I,LARGE),I=1,3)
IF((REDO).AND.(TIME.LE.TST)) GO TO 364
IF(DUMP) CALL ABEND(9)
STOP
364 CONTINUE
IP(NEW) GO TO 365
REWIND 9
READ(9) DENP,U,XREP,TRP,KAWLS,NL,NW,NH,BW,BH,NREG,XLB,XLC,PI,ND,
2 S,SINANG,COSANG,AKN,AKT,NBX,RM,XR,TIME,DTM,TI,ITS,ITP,TST,
3 TLIM,RMA,RNU,DIR,XSTART,MNM,MNB,TR,BZC,CN7,DRP,FCP,FNA,
4 HTF,INM,LLM,NAV,NMAX,NWEDG,PRT,SAMP,AKN1,AKN2,AKT1,AKT2,
5 BTA,C1,C2,C3,C7,C8,DAM,DPA,PL,DELANG,FDN,HTI,HTR,JNT,KNM,
6 NM,WTM,C4,VRM,NCOL,CTI,CTR,CNI,CNR,SN,
7 ST,D1,D2,D3,D4,NRAN,VELR,RMP,RMN,RMP,IPLUX,FLUXIN,
8 XLIM,COEFF,XCB,XS,YCB,TB,ALPHA,SIGMA,NTS,
9 UTL,UTT,VTS,HTS,HTSI,ENT,REM,TMPA,
A DBA,NB,NBT,TMP,XV,XVA,YV,YVA,ZV,ZVA,T,DB,PNB,XC,YC,ZC,
B PAU,PAV,PAW,PAY,PAZ,LCOL,LM,
C ETA,PHI,CHI,CN,CM,CNG,CNG,CN8,TRP,TRPA,MSP,ANGLE,TF,
D UTLI,UTTI,VTSI,ER,RMB,LB,NBM,NBN,VEL,PPV,FCOL,JV
365 ANM=INM
INM=9*ANM/10
DDN=.9*DDN
DRP=DRP/.9
FCP=FCP/.9
HTF=HTF/.9
DO 370 MM=1,MSP
FDN(MM)=FDN(MM)*INM/ANM
LL(MM)=9*LL(MM)/10
DO 366 KK=1,MSP
366 CN8(KK,MM)=CN8(KK,MM)*.9
DO 370 NK=1,2
ENT(MM,NK)=ENT(MM,NK)*INM/ANM
370 REM(MM,NK)=0.0
IF(NEW) GO TO 220
TST=TIME+TST
TI=0.
PRT=ITP
WRITE(6,2)
WRITE(6,4)
WRITE(6,5) (MT,RMA(MT),RNU(MT),FCOL(MT),FLUXIN(MT),(ENT(MT,K),K=1,2
1),MT=1,MSP)

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      WRITE(6,2)
      IF((LARGE.EQ.2).OR.(LARGE.EQ.3)) GO TO 280
      REDO=.FALSE.
      GO TO 360
      END

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RUN7410

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SUBROUTINE DIAG(N, ITEST, NUM)
REAL*8 PARAM(10)/' NWEDGE', ' NREG', ' ND', ' NPX', ',
1 MNN', ' MNB', ' NBX', ' NS', ' MJ', ' MSP'/

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DIAG010

DIAG040

DIAG050

DIAG060

DIAG070

DIAG080

FORMATS

DIAG090

DIAG100

DIAG120

DIAG130

DIAG140

DIAG150

DIAG160

DIAG170

DIAG180

DIAG190

DIAG210

DIAG220

DIAG230

DIAG240

DIAG250

DIAG260

DIAG270

DIAG280

DIAG290

DIAG300

DIAG310

DIAG320

DIAG340

DIAG350

DIAG370

DIAG390

DIAG400

DIAG420

DIAG440

DIAG460

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      WRITE(6,42)
      WRITE(6,44) ITEST, NUM, PARAM(N)
      WRITE(6,56)
      GO TO (1,2,3,4,5,6,7,8,9,10), N
1      WRITE(6,62)
      WRITE(6,32)
      GO TO 11
2      WRITE(6,64)
      GO TO 11
3      WRITE(6,66)
      WRITE(6,62)
      GO TO 11
4      WRITE(6,68)
      WRITE(6,76)
      GO TO 11
5      WRITE(6,70)
      GO TO 11
6      WRITE(6,72)
      GO TO 11

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7 WRITE(6,80)                                     DIAG480
  GO TO 11
8 WRITE(6,78)                                     DIAG500
  GO TO 11
9 WRITE(6,80)                                     DIAG520
  GO TO 11
10 WRITE(6,75)
11 WRITE(6,74)
  STOP
  END

```

```

SUBROUTINE PRINTA(NWEDG,TITLE,NAME,XCB,YCB,TB,ALPHA,SIGMA,XLIM,
1COEFF,LIMIT,MSP)
  INTEGER TST,TLIM,TIME
  LOGICAL SAVE,NEW
  DIMENSION LIMIT(1),TITLE(6),NAME(2),XCB(1),YCB(1),TB(1)
  DIMENSION ALPHA(3,1),SIGMA(3,1),XLIM(1),COEFF(4)
  DIMENSION RNU(3),RMA(3),CHI(3),DIR(3,3),PHI(3,3),ETA(3,3)
  DIMENSION WTM(3),DAM(3,3),VELS(3),XSP(3)
  COMMON /FIRST/NL,NW,NH
  COMMON /SECND/BW,BH,RMP,RMN,RMF
  COMMON /THIRD/PI,NREG,S,SINANG,COSANG,AKN,AKT,AKN1,AKN2,AKT1,AKT2
  COMMON /FOURTH/NBX,RM,XR,DUMP,C9,LL(3),LLM
  COMMON /FIFTH/ND,TIME,DTM,TI,ITS,ITP,TST,TLIM,RMA,RNU,DIR
  COMMON /SIXTH/RMB,XSTART,INM,MNM,MNB,NEW,SAVE,PERCNT,NSR,TR
  COMMON/EIGHT/DENP,U,TF,ANGLE,TRF,CHI,PHI,ETA,WTM,DAM,VELR,XREF
  DATA NOT/'NOT'/

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PRA0030
PRA0040

PRA0100
PRA0120
PRA0130
PRA0140
PRA0150
PRA0160
PRA0170

FORMATS

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1 FORMAT(16X,40(''),T74,'I'//9X,'3-D',I2,'-FLUID PROGRAM - ')
2 FORMAT('+',31X,A4)                           PRA0190
3 FORMAT('+',35X,      'A RESTART OF A PREVIOUS RUN',T74,'I'/12X,2PRA0200
  1A4,' - ',6A4,' - ',I2,' REGIONS',T74,'I',16(/T74,'I'))      PRA0210
4 FORMAT(7X,'FRONT OF BODY =',E12.4,' XSTAET MAX HEIGHT =',E12.4,
  1RMB',T74,'I'/7X,'BODY TEMPERATURES =',F12.2,' T PR.STRM.',F12.2,
  2' T ENTR.',F12.2,' T CAVITY'/7X,'X-LIMIT',T37,'BODY COEFFICIENTS',
  3T74,'I')
6 FORMAT(5F14.6,3X,'I')                         PRA0240
10 FORMAT(1X,72(''))                           PRA0250
12 FORMAT(//14X,'PARAMETERS OF SEGMENTS FOR BODY COLLISIONS',T96,'I'/
  18X,'X-COORD. TEMP. ALPHA1 ALPHA2 ALPHA3 SIGMA1 SIGMA2
  2 SIGMA3 AREAS',T96,'I')
14 FORMAT(4X,E12.4,7F9.4,E12.4,T96,'I')
17 FORMAT(//25X,'ARRAY STORAGE USED'/5X,I6,'      *',10I6,T96,'I')
18 FORMAT(1H1/17X,'LENGTH OF CELL IN MEAN-FREE-PATHS = ',F12.4,' BW'
  A,T76,'I'
  1/17X,'HEIGHT OF CELL IN MEAN-FREE-PATHS = ',F12.4,' BH',T76,'I'
  2/16X,'NUMBER OF L1 CELLS ALONG FLOW AXIS =',I13,' NW',T76,'I'
  3/17X,'NUMBER OF L1 CELLS IN RADIAL DIR. =',I13,' NH',T76,'I'
  4/21X,'NUMBER OF LEVELS OF CELL SIZE =',I13,' NL',T76,'I')
23 FORMAT(3X,'NUMBER OF AZIMUTHAL WEDGES WITHIN ',I3,' DEGREES =',I,PRA0470
  113,' NWEDG I'//)

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24 FORMAT(16X,'BASIC TIME INTERVAL FOR COLLISIONS =',E13.4,' DTM
1 I'/8X,'TIME INTERVAL FOR SAMPLING FLOW FIELD INFO =',E13.4,' DTS
2 I'/24X,'TIME INTERVAL FOR PRINTING =',E13.4,' DTP I'/9X,
3' TIME TO STEADY-STATE CONDITIONS (ASSUMED) =',E13.4,' TST I'/
419X,'TIME AT WHICH RUN IS TERMINATED =',E13.4,' TLIM I'/)
25 FORMAT(/5X,'PRESSURE RATIO(INSIDE/ENTRANCE) - EITHER TYPE =',F13.5
1,' RMP',T76,'I'/6X,'DENSITY RATIO(INSIDE/ENTRANCE) - EITHER TYPE
2=',F13.5,' RMN',T76,'I'/9X,'FLUX RATIO(INSIDE/ENTRANCE) - EITHER
3TYPE =',F13.5,' RMP',T76,'I')
26 FORMAT(5X,'FREE STREAM NUMBER OF MOLECULES - EITHER TYPE =',I13,
1 A INM I'/9X,'INITIAL NUMBER OF MOLECULES - MAXIMUM =',I13,
1' LLM I'/9X,'MAXIMUM NUMBER OF MOLECULES - EITHER TYPE =',I13
2,' MNB I'/1X,'MAX NUMBER OF MOLECULES IN ANY CELL - EITHER TY
3PE =',I13,' MNB I')
27 FORMAT( /22X,'VELOCITY OF FREE STREAM FLOW =',E13.4,' U',T76,'I'/1
19X,'SPEED RATIO OF FREE STREAM FLOW =',E13.4,' S',T76,'I'/19X,'MAC
AH NUMBER OF FREE STREAM FLOW =',E13.4,' M',T76,'I'/19X,'SPECIFIC H
BEAT RATIO (CALCULATED) =',E13.4,' GAMMA',T76,'I'/ 35X,'ANG
2LE OF ATTACK =',F13.4,' ANGLE I'/16X,'NUMBER DENSITY OF FREE ST
3REAM FLOW =',E13.4,' N',T76,'I'/19X,'TEMPERATURE OF FREE STREAM PL
4OW =',F13.4,' TF',T76,'I'/16X,'MOLE FRACTIONS OF FREE STREAM FLOW
5=',3E13.4,' RNU I'/16X,'MOLECULAR WEIGHTS OF SPECIES ABOVE =',3F13
6.4,' RMA I'/18X,'INITIAL NUMBERS OF SPECIES ABOVE =',3I13,' LL I'
7)
28 FORMAT( /10X,'REFERENCE TEMPERATURE FOR MOLECULAR DATA =',F13.4,
1' TRF',T90,'I'/14X,'CROSS-SECTION',26X,'TEMP EXPONENT',T90,'I'/3(3X,
23E12.4,3X,3F12.6,T90,'I')/ 5X,'CHI/2-1',11X,'ROTATIONAL PARAMETER
3 PHI',T90,'I'/3(F12.4,5X,3F12.6,T90,'I'))
29 FORMAT(/9X,'DATA SAVED ON TAPE 9')
30 FORMAT( 31,'REF MOLECULAR SPEED =',E13.4,' VELR',T76,'I'/20X,'SP
1ECIES FREE STREAM MOLECULAR SPEEDS',T76,'I'/14X,3E16.6,T76,'I'/26X
2,'REFERENCE MEAN FREE PATH =',E13.4,' XREF',T76,'I'/26X,'SPECIES M
3EAN FREE PATHS',T76,'I'/14X,3E16.6,T76,'I'/11X,'LONGITUDINAL KNUDS
4EN NUMBER (FREE STRM.) =',E13.4,' AKN',T76,'I'/13X,'TRANSVERSE KNUD
5SEN NUMBER (FREE STRM.) =',E13.4,' AKT',T76,'I'/11X,'LONGITUDINAL K
6NUDSEN NUMBER ( ENTRANCE ) =',E13.4,' AKN1',T76,'I'/13X,'TRANSVERSE
7 KNUDSEN NUMBER ( ENTRANCE ) =',E13.4,' AKT1',T76,'I'/11X,'LONGITUD
8INAL KNUDSEN NUMBER ( CAVITY ) =',E13.4,' AKN2',T76,'I'/13X,'TRAN
9SVERSE KNUDSEN NUMBER ( CAVITY ) =',E13.4,' AKT2',T76,'I')

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PRA0680

PRA0690

PRA0700

```

IARRAY=708+LIMIT(3)*(32+56*LIMIT(1))+20*LIMIT(2)+LIMIT(4)*(120+4*LPRA0710
LIMIT(6))+56*LIMIT(5)+LIMIT(8)*(68+96*LIMIT(9))+20*LIMIT(7)+224*LIMPRA0720
2IT(1)

```

PRA0730

```

WRITE(6,1) MSP
IF(NEW) WRITE(6,2) NOT
WRITE(6,3) NAME,TITLE,NREG
WRITE(6,4) XSTART,RMB
DO 100 I=1,NREG
100 WRITE(6,6) XLIM(I+1),(COEFF(J),J=1,4)
WRITE(6,10)
WRITE(6,12)
DO 110 I=1,ND
110 WRITE(6,14) XCB(I),TB(I),(ALPHA(J,I),J=1,3),(SIGMA(J,I),J=1,3),YCB

```

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PRA0750

PRA0760

PRA0770

PRA0780

PFL0800

PRA0810

PRA0820

```

1(I)
  WRITE(6,10)                                            PRA0850
  WRITE(6,17) IARRAY, (LIMIT(I), I=1,10)
  WRITE(6,18) BW,BH,NW,NH,NL                            PRA0880
  IETAZ=180.
  WRITE(6,23) IETAZ,NWEDG
  DTS=DTM*ITS
  DTP=DTM*ITP
  AST=DTM*TST
  ALIM=DTM*TLIM
  CHT=0.0
  DO 120 J=1,MSP
120 CHT=CHT+CHI(J)*RNU(J)
  GAMMA=(7.+2.*CHT)/(5.+2.*CHT)
  AM=S*SQRT(2./GAMMA)
  WRITE(6,24) DTM,DTS,DTP,AST,ALIM
  WRITE(6,26) INM,LLM,MNM,MNB
  WRITE(6,25) RMP,RMN,RMP
  WRITE(6,27) U,S,AM,GAMMA,ANGLE,DENP,TF, (RNU(I), I=1,3), (RMA(I), I=1,3
1), (LL(I), I=1,3)
  WRITE(6,28) TRP, ((DIR(I,K), K=1,3), (ETA(I,K), K=1,3), I=1,3), (CHI(I),
1 (PHI(I,K), K=1,3), I=1,3)
  DO 210 I=1,3
  VELS(I)=0.0
210 XSP(I)=0.0
  DO 220 J=1,MSP
  VELS(J)=VELR/SQRT(WTM(J))
  XT=0.0
  DO 215 M=1,MSP
215 XT=XT+RNU(M)*DAM(J,M)*SQRT(1.+WTM(J)/WTM(M))
220 XSP(J)=1.414214*XREP/XT
  WRITE(6,30) VELR, (VELS(I), I=1,3), XREP, (XSP(I), I=1,3), AKN, AKT,
1AKN1, AKT1, AKN2, AKT2
  IP(SAVE) WRITE(6,29)
  RETURN
  END

```

```

SUBROUTINE PRINTB(PNA,MSP,FNB,NM,XLIM,XC,YC,ZC,NB,N)
INTEGER*2 NB
DIMENSION FNB(1),NM(1),XLIM(1),XC(1)
DIMENSION YC(1),ZC(1),NB(N,1)
COMMON /FIRST/NL,NW,NH
COMMON /THIRD/PI,NREG,S,SINANG,COSANG,AKN,AKT
COMMON /FORTH/NBX
1 FORMAT(1H1)                                            PRB0070
2 FORMAT(2X,-----CELL GEOMETRY-----)                  PRB0080
1-----'/2X, 'BOX LEVEL POSITION OF CENTER VOLUME PRB0100
2INITIAL POPULATION'/2X,'NUM.',12X,'X',7X,'Y THETA',12X,' TOT
3AL ',', EACH SPECIES ',', CELL#)
3 FORMAT(1X,I4,I5,3X,2F8.3,F7.1,E12.3,2X,I4,2X,3I5,3X,I4)
4 FORMAT(2X,-----TOTALS-----',E12.4,8X,3I5)
  WRITE(6,1)
  WRITE(6,2)
  DO 200 I=1,NBX
  X=(YC(I)-XLIM(1))*AKN

```

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PRA0940
PRA0950
PRA0960
PRA0970

PRA0980

PRA1040
PRA1050
PRA1060

PRB0070
PRB0080

PRB0090

PRB0100
PRB0105

PRB0110

PRB0120

PRB0130

PRB0140

PRB0150

PRB0170

PRB0180

```

Y=YC(I)*AKT*2.0
IX=IX*NW+1
IY=IY*NH+1
M1=NB(1,I)
M2=0
IF(MSP.GE.2) M2=NB(2,I)
M3=0
IF(MSP.GE.3) M3=NB(3,I)
MM=M1+M2+M3
* 140 WRITE(6,3) IX,IY,X,Y,ZC(I),PNB(I),MM,M1,M2,M3,I
200 CONTINUE
NM2=0
IF(MSP.GE.2) NM2=NM(2)
NM3=0
IF(MSP.GE.3) NM3=NM(3)
* WRITE(6,4) FNA,MM(1),NM2,NM3
RETURN
END

```

PRB0370

PRB0390
PRB0400

```

SUBROUTINE CELL(A,B,KW,KH,XO,DELANG,NWEDG,XC,YC,ZC,PNB)
DIMENSION XC(1),YC(1),ZC(1),PNB(1)
COMMON /THIRD/PI

```

CELL030
CELL040
CELL050
CELL060
CELL070
CELL080
CELL090
CELL100
CELL110
CELL120
CELL130

THE PURPOSE OF THIS SUBROUTINE IS TO

1. COMPUTE THE VOLUME OF EACH CELL (ALL 3 POSSIBLE LEVELS) AND STORE THE RESULT IN THE ARRAY CALLED 'PNB'.
2. COMPUTE THE X, R, AND THETA COORDINATES OF THE CENTER OF EACH CELL (ALL 3 POSSIBLE LEVELS) AND STORE THE RESULTS IN CELLS CALLED 'XC', 'YC', AND 'ZC'.

```

I=0
X=XO-0.5*A
FACTOR=DELANG*PI*B*B*A/180.
DO 110 K=1,KW
X=X+A
Y=-.5*B
DO 110 L=1,KH
Y=Y+B
Z=-.5*DELANG
DO 110 M=1,NWEDG
Z=Z+DELANG
I=I+1
XC(I)=X
YC(I)=Y
ZC(I)=Z
110 PNB(I)=FACTOR*(2*L-1)
RETURN
END

```

CELL360
CELL370

```

SUBROUTINE IMPACT(RM,G1,G2,G3,ET,EI,PHI,CHI,ETA,XM,CIM)
COMMON/THIRD/PI
IF(PHI.EQ.0.) GO TO 20

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```

IF (CHI.EQ.0.) GO TO 20
DF=PHI*CHI-1.
DS=PHI*(2.-.5*ETA)-1.
E=ET+EI
10 X=RAND(0)
IF (X.EQ.0.0) GO TO 10
XT=X**DF*(1.-X)**DS
IF (XT.GT.XM) GO TO 15
CIM=CIM+XT
IF (CIM.LT.XM) GO TO 10
CIM=CIM-XM
15 ET=(1.-PHI)*ET+(1.-X)*PHI*E
EI=(1.-PHI)*EI+X*PHI*E
20 GP=SQRT(ET/RM)
EP=2.*PI*RAND(0)
CSX=2.*RAND(0)-1.
SSX=SQRT(1.-CSX**2)
G1=GP*CSX
G2=GP*SSX*COS(EP)
G3=GP*SSX*SIN(EP)
RETURN
END

```

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```

SUBROUTINE GAS(NWEDG, DELANG, ND, BTA, C1, DFA, NM, FNB, DB, NB, NBM, NBN,
1 PAU, PAV, PAX, PAY, PAZ, XLIM, COEFF, LM, I2, I3, XCB, TB, LARGE,
2 NM, MNB, DEBUG1, LCOL, IP, ER, CHI, CNG, CMG, I, LB)
INTEGER*2 LM(I,1), LCOL(I,1), LB(1), NBM(I,1), NBN(1)
INTEGER*2 NB
LOGICAL DUMP, DEBUG1
DIMENSION BTA(1), C1(1), DFA(1), NM(1), FNB(1), CHI(1)
DIMENSION DB(I,1), NB(I,1), PAU(I,1), PAV(I,1), PAW(I,1)
DIMENSION PAX(I,1), PAY(I,1), PAZ(I,1), ER(I,1), COEFF(4), XLIM(1)
DIMENSION CNG(1), CMG(1), XCB(1), TB(1)
COMMON /FIRST/NL, NW, NH
COMMON /SECND/BW, BH, RMP, RMN, RMP
COMMON /THIRD/PI, NREG, S, SINANG, COSANG, AKN, AKT
COMMON /FORTH/NBX, RM, XR, DUMP, C9, LL(3)
-----
```

GAS0060

GAS0150

GAS0160

GAS0170

1. COMPUTE THE INITIAL VELOCITY OF EACH MOLECULE. GAS0180
2. COMPUTE THE INITIAL POSITION OF EACH MOLECULE. GAS0190
3. CREATE AN ARRAY WHICH STORES THE CELL POPULATIONS- GAS0200
4. CREATE A CROSS-REFERENCING ARRAY (WHOSE CONSTRUCTION IS GAS0210
5. COMPUTE AN ARRAY WHICH STORES THE NUMBER DENSITY IN EACH GAS0220
6. COMPUTE THE NUMBER OF MOLECULES IN EACH CELL GAS0230
7. COMPUTE THE NUMBER OF MOLECULES IN EACH CELL GAS0240
8. COMPUTE THE NUMBER OF MOLECULES IN EACH CELL GAS0250
9. COMPUTE THE NUMBER OF MOLECULES IN EACH CELL GAS0260
10. COMPUTE THE NUMBER OF MOLECULES IN EACH CELL GAS0270
11. COMPUTE THE NUMBER OF MOLECULES IN EACH CELL GAS0280
12. COMPUTE THE NUMBER OF MOLECULES IN EACH CELL GAS0290

GAS0300

```

2 FORMAT(/' SOMETHING IS WRONG WITH BOX NUMBERING IN GAS'/9I5,5E14.5
1//)
3 FORMAT(/' SOMETHING WRONG IN CELL VOLUMES IN GAS'/5X,5I5,2E14.5)

```

```

4 FORMAT(' NB(',I2,',',',I4,',') POPULATION EXCEEDED ',I3,' IN GAS')      GAS0340
BP=(1.-RMN)
CP=1.-RMN**2
DO 180 MT=1,IP
N=0
110 N=N+1
IF(N.GT.LL(MT)) GO TO 180
IF(N.GT.MNN) GO TO 190
120 P=.001+.998*RAND(0)
IF(BP.NE.0.0) P=(1.-SQRT(1.-CP*P))/BP
X=XLIM(1)+*XR
R=RM*SQRT(RAND(0))
D=PI*RAND(0)
PA(X,MT,N)=X
PA(Y,MT,N)=R*COS(D)
PA(Z,MT,N)=R*SIN(D)
DO 126 J=2,ND
IF(X.LE.XCB(J)) GO TO 128
126 CONTINUE
WRITE(6,3) J,ND,MT,LL(MT),N,X,XCB(J)
IF(DUMP) CALL ABEND(14)
STOP
128 TL=TB(J-1)+(TB(J)-TB(J-1))*(X-XCB(J-1))/(XCB(J)-XCB(J-1))
130 V=4.*RAND(0)
VV=V*V
C1(MT)=C1(MT)+VV*EXP(1.-VV)
IF(C1(MT).LT.1.) GO TO 130
C1(MT)=C1(MT)-1.
A=1.-2.*RAND(0)
B=SQRT(1.-A*A)
C=2.*PI*RAND(0)
V=V/BTA(MT)*SQRT(TL)
PA(U,MT,N)=V*A
PA(V,MT,N)=V*B*COS(C)
PA(W,MT,N)=V*B*SIN(C)
EX=0.0
IF(CHI(MT).LE.-1.) GO TO 136
135 EX=9.*RAND(0)
IF(EX.EQ.0.0) GO TO 135
XT=EX**CHI(MT)*EXP(-EX)
IF(XT.GE.CMG(MT)) GO TO 136
CNG(MT)=CNG(MT)+XT
IF(CNG(MT).LT.CMG(MT)) GO TO 135
CNG(MT)=CNG(MT)-CMG(MT)
136 ER(MT,N)=EX*TL
TANG=180.*(1.-D/PI)
IWDGE=TANG/DELANG+1
IF(IWDGE.GT.NWEDG) IWDGE=NWEDG
L=X/BW
IF(L.GE.NW) L=NW-1
M=R/BH
IF(M.GE.NH) M=NH-1
K=NWEDG*(L*NH+M)+IWDGE
IF(K.LE.NBX) GO TO 165
WRITE(6,2) L,M,IWDGE,NW,NH,NWEDG,K,MT,N,X,Y,Z,R,TANG

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GAS0370
GAS0380
GAS0390

GAS0400
GAS0410
GAS0420
GAS0430
GAS0440
GAS0450
GAS0460
GAS0470

GAS0510

GAS0720

GAS0780
GAS0790

IF (DUMP) CALL ABEND(11)	GAS0830
STOP	GAS0840
165 J=NB(MT,K)+1	GAS1190
LCOL(MT,N)=0	
IF (J.LE.MNB) GO TO 166	GAS1240
IF (DEBUG1) WRITE(6,4) MT,K,MNB	GAS1250
166 NB(MT,K)=J	GAS1270
LB(N)=K	
167 IF (N.LT.LL(MT)) GO TO 110	
NM(MT)=N	GAS1310
NBM(MT,1)=0	
DO 170 N=1,NBX	GAS1330
A=NB(MT,N)	GAS1360
DB(MT,N)=A*DFA(MT)/PNB(N)	GAS1370
NBM(MT,N+1)=NBM(MT,N)+NB(MT,N)	
NBN(N)=NBM(MT,N)	
170 CONTINUE	GAS1380
NG=NM(MT)	
DO 175 N=1,NG	
NQ=LB(N)	
NBN(NQ)=NBN(NQ)+1	
NA=NBN(NQ)	
175 LM(MT,NA)=N	
180 CONTINUE	GAS1390
RETURN	GAS1400
190 LARGE=1	GAS1410
RETURN	GAS1420
END	GAS1430

SUBROUTINE FLOW(NWEDG,MNM,LARGE,BTA,C1,C7,C8,ENT,REM,LCOL,IP,NM,
 1SN,ST,TBI,PAU,PAV,PAW,PAX,PAZ,ER,CHI,CNG,CMG,I,JV,FCOL,VEL,
 2PFV)

INTEGER*2 LCOL	
DIMENSION BTA(1),NM(1),SN(1),ST(1)	
DIMENSION C1(1),C7(1),C8(1),CMG(1),FCOL(1),VELK(4)	
DIMENSION PAU(I,1),PAV(I,1),PAW(I,1),PAX(I,1),PAY(I,1),PAZ(I,1)	
DIMENSION ENT(3,1),REM(3,1),LCOL(I,1),ER(I,1),CHI(1),CNG(1)	
DIMENSION VEL(JV,4,1),PFV(JV,4,1)	
COMMON /THIRD/PI	FLO0100
COMMON /FORTH/NBX,RM,XR	FLO0110

 THE PURPOSE OF THIS SUBROUTINE IS TO ADD A NEW BATCH OF MOLECULES
 TO THE SAMPLE THROUGH THE UPSTREAM BOUNDARY.

-----	FLO0130
-----	FLO0140
-----	FLO0150
DO 370 MT=1,IP	
XGO=0.	FLO0180
E=1.	FLO0190
FRAC=FCOL(MT)	
ARG=SN(MT)	
STT=ST(MT)	
TV=1./BTA(MT)	
TR=1.	
DO 180 NT=1,2	FLO0200
VM=(SQRT(ARG**2+2.0)+ARG)/2.	
SM=AMAX1(0.0,VM-4.)	

```

SMM=VM+4.-SM
AM=ENT(MT,NT)+REM(MT,NT)
M=AM
REM(MT,NT)=AM-M
IF(M.EQ.0) GO TO 170
DO 160 N=1,M
IF(NM(MT).GE.NNM) GO TO 380
NM(MT)=NM(MT)+1
NMX=NM(MT)
R=RM*SQRT(RAND(0))
D=PI*RAND(0)
PAZ(MT,NMX)=R*COS(D)
PAW(MT,NMX)=R*SIN(D)
LCOL(MT,NMX)=0
IF(FRAC.EQ.0.0) GO TO 130
PF=RAND(0)
IF(PF.GT.FRAC) GO TO 130
KMX=3
VELK(4)=0.0
IF(CHI(MT).GT.-1) KMX=4
DO 110 K=1,KMX
P=RAND(0)
DO 102 J=2,JV
IF(PFV(J,K,MT).GT.P) GO TO 105
102 CONTINUE
VELK(K)=VEL(JV,K,MT)
GO TO 110
105 VELK(K)=VEL(J-1,K,MT)+(P-PPV(J-1,K,MT))*(VEL(J,K,MT)-VEL(J-1,K,MT))/(PFV(J,K,MT)-PFV(J-1,K,MT))
110 CONTINUE
PAU(MT,NMX)=VELK(1)
PAV(MT,NMX)=VELK(2)
PAW(MT,NMX)=VELK(3)
ER(MT,NMX)=VELK(4)
GO TO 160
130 V=SMM*RAND(0)*SMM
C1(MT)=C1(MT)+V*EXP(VM**2-V**2+2.*ARG*(V-VM))/VM
IF(C1(MT).LT.1.) GO TO 130
C1(MT)=C1(MT)-1.
PAU(MT,NMX)=E*V*TV
140 V=8.*RAND(0)-4.
C7(MT)=C7(MT)+EXP(-V*V)
IF(C7(MT).LT.1.) GO TO 140
C7(MT)=C7(MT)-1.
PAV(MT,NMX)=STT+V*TV
150 V=8.*RAND(0)-4.
C8(MT)=C8(MT)+EXP(-V*V)
IF(C8(MT).LT.1.) GO TO 150
C8(MT)=C8(MT)-1.
PAW(MT,NMX)=V*TV
X=0.0
IF(CHI(MT).LE.-1.) GO TO 156
155 X=9.*RAND(0)
IF(X.LE.0.0) GO TO 155
XT=X**CHI(MT)*EXP(-X)

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PL00340
PL00360
PL00370
PL00390
PL00420
PL00430

```

IF (XT.GE.CMG(MT)) GO TO 156
CNG(MT)=CNG(MT)+XT
IF (CNG(MT).LT.CMG(MT)) GO TO 155
CNG(MT)=CNG(MT)-CMG(MT)
156 ER(MT,NMX)=X*TR
160 PAX(MT,NMX)=XGO
170 CONTINUE
  ARG=0.0
  XGO=XR
  E=-1.
  FRAC=0.0
  STT=0.0
  TV=SQRT(TBI)/BTA(MT)
  TR=TBI
180 CONTINUE
370 CONTINUE
  RETURN
380 LARGE=2
  RETURN
  END

```

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FLO0610
FLO0620
FLO0640
FLO0650
FLO0660
FLO1140
FLO1150
FLO1160
FLO1170
FLO1180

```

SUBROUTINE COLIDE(CN,CM,WTM,DB,DBA,NB,NCOL,LCOL,PAU,PAV,PAW,ER,T,
1LM,MT,I2,I3,ETA,PHI,CHI,CN8,NP,NBM)
  INTEGER TIME
  INTEGER*2 LM(NP,1), LCOL(NP,1)
  INTEGER*2 NBM,NB
  DIMENSION CN(3,3,1),CM(3,3,1),WTM(1),DB(NP,1),DBA(NP,1),NB(NP,1)
  DIMENSION NBM(NP,1),NCOL(3,1),T(NP,NP,1),ETA(3,1),PHI(3,1),CHI(1)
  DIMENSION PAU(NP,1),PAV(NP,1),PAW(NP,1),ER(NP,1),CN8(3,1),WA(2)
  COMMON /FOURTH/NBX
  COMMON /FIFTH/ND,TIME,DTM

```

COL0030
COL0080
COL0090
COL0100

THE PURPOSE OF THIS SUBROUTINE IS TO ADVANCE THE ELAPSED TIMES INCOL0110
CELLS BY AN AMOUNT APPROXIMATELY EQUAL TO THE PRE-SELECTED COLLISCOL0120
TIME. THERE ARE FOUR TIMES FOR EACH CELL, SAVED IN AN ARRAY CALLECOL0130
'T', CORRESPONDING TO THE FOUR TYPES OF MOLECULAR COLLISIONS WHICCOL0140
CAN OCCUR. TO ADVANCE THE VARIOUS TIMES, AN APPROPRIATE NUMBER OFCOL0150
THE CORRESPONDING MOLECULAR COLLISIONS IS COMPUTED. THE ACTUAL COL0160
MOLECULES TO COLLIDE ARE SELECTED AT RANDOM, AND THEIR VELOCITY VCOL0170
DIRECTIONS AFTER COLLISION ARE SELECTED AT RANDOM. COL0180
-----COL0190

```

AIME=DTM*TIME
DO 240 MTA=1,MT
DO 230 MTB=1,MTA
  D = WTM(MTA) + WTM(MTB)
  WA(MTA)=WTM(MTA)/D
  WA(MTB)=WTM(MTB)/D
  RM=WTM(MTA)*WTM(MTB)/D
  CHT=CHI(MTA)+CHI(MTB)+2.0
  PHT=PHI(MTA,MTB)
  ETT=ETA(MTA,MTB)
  DO 220 N=1,NBX
  IF (T(MTA,MTB,N).LT.AIME) GO TO 100
  IF (T(MTB,MTA,N).GE.AIME) GO TO 220
100 NA=NB(MTA,N)*NB(MTB,N)

```

```

IF(MTA.EQ.MTB) NA=(NA-NB(MTA,N))/2.
IF(NA.LT.1) GO TO 220
KS=0
*20 KC=0
CPUT=ELTIME(0)
KS=KS+1
IF(KS.GT.NA) GO TO 220
130 KC=KC+1
IF(KC.GT.NA) GO TO 220
*35 I=NB(MTA,N)*RAND(0)+1+NBM(MTA,N)
IF(I.GT.NBM(MTA,N+1)) I=NBM(MTA,N+1)
J=LM(MTA,I)
140 K=NB(MTB,N)*RAND(0)+1+NBM(MTB,N)
IF(K.GT.NBM(MTB,N+1)) K=NBM(MTB,N+1)
IF(MTA.EQ.MTB.AND.I.EQ.K) GO TO 140
* L=LM(MTB,K)
GM1=WA(MTA)*PAU(MTA,J)+WA(MTB)*PAU(MTB,L)
GM2=WA(MTA)*PAV(MTA,J)+WA(MTB)*PAV(MTB,L)
GM3=WA(MTA)*PAW(MTA,J)+WA(MTB)*PAW(MTB,L)
G1=PAU(MTA,J)-PAU(MTB,L)
G2=PAV(MTA,J)-PAV(MTB,L)
* G3=PAW(MTA,J)-PAW(MTB,L)
GS=G1**2+G2**2+G3**2
IF(GS.LT.1.0E-8) GO TO 130
ET=RM*GS
EI=ER(MTA,J)+ER(MTB,L)
VR=GS**(.5-ETT/2.)
* IF(VR.GE.CN(MTA,MTB,1)) GO TO 160
CN(MTA,MTB,1)=CN(MTA,MTB,1)+VR
IF(CN(MTA,MTB,1).LT.CM(MTA,MTB,1)) GO TO 130
CN(MTA,MTB,1)=CN(MTA,MTB,1)-CM(MTA,MTB,1)
160 CONTINUE
CPUT=ELTIME(0)
* CALL IMPACT(RM,G1,G2,G3,ET,EI,PHT,CHT,ETT,CN(MTA,MTB,2),CN(MTA,MTB
1,2))
165 CONTINUE
CPUT=ELTIME(0)
IF(PHT.EQ.0.) GO TO 175
X1=0.0
* IF(CHI(MTA).EQ.-1.) GO TO 175
X1=1.0
IF(CHI(MTB).EQ.-1.) GO TO 175
170 X1=RAND(0)
IF((CHI(MTA).EQ.0.).AND.(CHI(MTB).EQ.0.)) GO TO 175
XT=X1**CHI(MTA)*(1.-X1)**CHI(MTB)
IF(XT.GT.CN(MTA,MTB,3)) GO TO 175
CN(MTA,MTB,3)=CN(MTA,MTB,3)+XT
IF(CN(MTA,MTB,3).LT.CM(MTA,MTB,3)) GO TO 170
CN(MTA,MTB,3)=CN(MTA,MTB,3)-CM(MTA,MTB,3)
175 CONTINUE
C=DBA(MTA,N)
D=DBA(MTB,N)
IF(C.EQ.0.0) C=DB(MTA,N)
IF(D.EQ.0.0) D=DB(MTB,N)
IF(T(MTA,MTB,N).GE.AIME) GO TO 180

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PAU(MTA,J)=GM1+WA(MTB)*G1
 PAV(MTA,J)=GM2+WA(MTB)*G2
 PAW(MTA,J)=GM3+WA(MTB)*G3
 IF(PHT.GT.0.) ER(MTA,J)=X1*EI
 LCOL(MTA,J)=1+LCOL(MTA,J)
 NCOL(MTA,MTB)=NCOL(MTA,MTB)+1
 T(MTA,MTB,N)=T(MTA,MTB,N)+CN8(MTA,MTB)/NB(MTA,N)/D/VR
 IF(MTA.EQ.MTB) GO TO 190
 180 IF(T(MTB,MTA,N).GE.AIME) GO TO 210
 190 PAU(MTB,L)=GM1-WA(MTA)*G1
 PAV(MTB,L)=GM2-WA(MTA)*G2
 PAW(MTB,L)=GM3-WA(MTA)*G3
 IF(PHT.GT.0.) ER(MTB,L)=(1.-X1)*EI
 LCOL(MTB,L)=1+LCOL(MTB,L)
 NCOL(MTB,MTA)=NCOL(MTB,MTA)+1
 T(MTB,MTA,N)=T(MTB,MTA,N)+CN8(MTB,MTA)/NB(MTB,N)/C/VR
 210 CONTINUE
 IF(T(MTA,MTB,N).LT.AIME.OR.T(MTB,MTA,N).LT.AIME) GO TO 120
 220 CONTINUE
 230 CONTINUE
 240 CONTINUE
 RETURN
 END

COL0810

COL0920

SUBROUTINE MOVE(KSWCH,AKN,NWEDG,XSTART,I2,I3,I4,I5,DELANG,
 1BTA,C2,C3,DFA,PL,HTI,HTR,JNT,KNM,NM,XCB,XLIM,CTI,CTR,CNI,
 2CNR,ALPHA,SIGMA,COEFF,HTS,HTSI,NTS,UTL,UTT,VTS,PAU,PAV,PAW,PAX,
 3PAY,PAZ,LCOL,TB,IP,ER,CHI,CNG,CMG,I,UTLI,UTTI,VTSI,IPLUX)
 INTEGER*2 LCOL(I,1)
 INTEGER TST,TIME
 LOGICAL DUMP
 REAL LAM,MU,NU

MOV0070

MOV0080

DIMENSION BTA(1),C2(1),C3(1),PL(1),HTI(1)
 DIMENSION HTR(1),TB(1),XCB(1),ALPHA(3,1),SIGMA(3,1),COEFF(4)
 DIMENSION PAU(I,1),PAV(I,1),PAW(I,1),CTI(3,1),CTR(3,1)
 DIMENSION CNI(3,1),CNR(3,1),DFA(1),JNT(1),XLIM(1),KNM(1),NM(1)
 DIMENSION HTS(3,I2,I3),HTSI(3,I2,I3),NTS(3,I2,I3)
 DIMENSION UTLI(3,I2,I3),UTTI(3,I2,I3),VTSI(3,I2,I3)
 DIMENSION UTL(3,I2,I3),UTT(3,I2,I3),VTS(3,I2,I3)
 DIMENSION PAX(I,1),PAY(I,1),PAZ(I,1)
 DIMENSION ER(I,1),CNG(1),CMG(1),CHI(1),IPLUX(3,2)
 COMMON /THIRD/PI,NREG
 COMMON /FORTH/NBX,RM,XR,DUMP
 COMMON /FIFTH/ND,TIME,DTM,TI,ITS,ITP,TST
 COMMON /SVNTH/LAM,MU,NU,MT,N,J,XI,YI,ZI,TUSE
 NAMELIST/CHECK/TIME,X,Y,Z,DX,DY,DZ,TLEFT,RADS,RMS,XR

MOV0170

MOV0180

MOV0200

 THE PURPOSE OF THIS SUBROUTINE IS TO ADVANCE THE SPATIAL POSITION
 OF ALL THE MOLECULES BY AN AMOUNT APPROPRIATE TO THEIR CURRENT VE
 LOCITIES AND THE PRE-SELECTED COLLISION TIME.

210

220

230

240

250

NAREA=NREG+1
 RMS=RM**2
 DO 150 MT=1,IP
 N=KNM(MT)

MOV0270

MOV0290

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```

10 N=N+1                                MOV0300
  TLEFT=DTM
  IF (KSWCH.EQ.1) TLEFT=TLEFT*RAND(0)
  IF (N.GT.NM(MT)) GO TO 150
15 LAM=PAU(MT,N)
  MU=PAV(MT,N)
  NU=PAW(MT,N)
  XI=PAX(MT,N)
  YI=PAY(MT,N)
  ZI=PAZ(MT,N)
  DX=TLEFT*LAM
  DY=TLEFT*MU
  DZ=TLEFT*NU
  X=XI+DX
  Y=YI+DY
  Z=ZI+DZ
  RADS=Y**2+Z**2
  TUSE=TLEFT
  KS=0
  IF (RADS.GT.RMS) CALL INTERS(X,Y,Z,RMS,KS)
  IF (X.GT.XLIM(1)) GO TO 18
  IPLUX(MT,1)=IPLUX(MT,1)+1
  GO TO 100
18 IF (X.LT.XLIM(NAREA)) GO TO 19
  IPLUX(MT,2)=IPLUX(MT,2)+1
  GO TO 100
19 CONTINUE
  IF (KS.GT.0) CALL DRAG(AKN,NWEDG,XSTART,I2,I3,I4,I5,DELANG,JNT,BTA,
  1C2,C3,DFA,PL,HTI,HTR,TB,XCB,CTI,CTR,CNI,CNR,ALPHA,SIGMA,RM,HTS,
  2HTSI,NTS,UTL,UTT,VTS,LCOL,IP,ER,CHI,CNG,CMG,I,UTLI,UTTI,VTSI)
  IF (Z.GT.0.0) GO TO 20
  Z=-Z
  NU=-NU
20 CONTINUE
  PAX(MT,N)=X
  PAY(MT,N)=Y
  PAZ(MT,N)=Z
  PAU(MT,N)=LAM
  PAV(MT,N)=MU
  PAW(MT,N)=NU
  IF ((X.GT.0.0).AND.(X.LT.XR).AND.((Y**2+Z**2).LT.RMS)) GO TO 25
  WRITE(6,CHECK)
  GO TO 100
25 CONTINUE
  TLEFT=TLEFT-TUSE
  IF (TLEFT.GT.0.0) GO TO 15
  GO TO 10
100 NZ=NM(MT)                                MOV0820
  PAX(MT,N)=PAX(MT,NZ)
  PAY(MT,N)=PAY(MT,NZ)
  PAZ(MT,N)=PAZ(MT,NZ)
  PAU(MT,N)=PAU(MT,NZ)
  PAV(MT,N)=PAV(MT,NZ)
  PAW(MT,N)=PAW(MT,NZ)
  ER(MT,N)=ER(MT,NZ)
  ORIGINAL PAGE IS
  OF POOR QUALITY
  MOV0830
  MOV0840
  MOV0850
  MOV0860
  MOV0870
  MOV0880

```

```

LCOL(MT,N)=LCOL(MT,NZ)          MOV0900
N=N-1                           MOV0910
NM(MT)=NM(MT)-1                MOV0920
GO TO 10                         MOV0930
150 CONTINUE                      MOV0940
RETURN                           MOV0950
END                               MOV0960

SUBROUTINE ACCUM(I2,I3,FNB,NB,PAU,PAV,PAW,ER,TMP,TRP,XV,YV,ZV,LM,
1IP,I,NBM)
INTEGER*2 LM(I,1),NBM(I,1)
INTEGER*2 NB
DIMENSION FNB(1),NB(I,1),PAU(I,1),PAV(I,1),PAW(I,1),TMP(I,1)
DIMENSION XV(I,1),YV(I,1),ZV(I,1),ER(I,1),TRP(I,1)
COMMON /FORTH/NBX               ACUM050
-----ACUM060
THE PURPOSE OF THIS SUBROUTINE IS TO ACCUMULATE TEMPERATURES,      ACUM070
VELOCITIES, AND DENSITIES IN VARIOUS ARRAYS FOR DETERMINING THE      ACUM080
AVERAGE FLOW FIELD PROPERTIES AFTER STEADY-STATE HAS BEEN REACHEDACUM090
-----ACUM100
DO 180 N=1,NBX                 ACUM120
DO 110 MT=1,IP                  ACUM140
XV(MT,N)=0.0                     ACUM160
YV(MT,N)=0.0                     ACUM170
ZV(MT,N)=0.0                     ACUM180
TMP(MT,N)=0.                      ACUM190
TRP(MT,N)=0.0                     ACUM200
TTX=0.                           ACUM210
TTY=0.                           ACUM220
TTZ=0.                           ACUM230
TTR=0.0                          ACUM240
M=NB(MT,N)                       ACUM250
IF(M.LT.1) GO TO 110             ACUM260
U=0.                             ACUM270
V=0.                             ACUM280
W=0.                             ACUM290
DO 100 L=1,M                     ACUM300
NA=NBM(MT,N)+L                  ACUM310
J=LM(MT,NA)                     ACUM320
PU=PAU(MT,J)
PV=PAV(MT,J)
PW=PAW(MT,J)
U=U+PU
V=V+PV
W=W+PW
TTR=TTR+ER(MT,J)
TTX=TTX+PU*PU
TTY=TTY+PV*PV
100 TTZ=TTZ+PW*PW
M=NB(MT,N)
XV(MT,N)=U/M                     ACUM390
YV(MT,N)=V/M                     ACUM400
ZV(MT,N)=W/M                     ACUM410
TMP(MT,N)=(TTX+TTY+TTZ)/M       ACUM420
TRP(MT,N)=TTR/M

```

```
110 CONTINUE
180 CONTINUE
  RETURN
END
```

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ACUM430
ACUM440
ACUM450
ACUM460

```
SUBROUTINE AVRGE(FNB,DB,DBA,NB,NBT,XV,YV,ZV,XVA,YVA,ZVA,TMP,TMPA,
1TRP,TRPA,IP,I)
  INTEGER*2 NB,NBT
  DIMENSION FNB(1),DB(I,1),DBA(I,1),NB(I,1),NBT(I,1),TMP(I,1)
  DIMENSION TMPA(I,1),XV(I,1),XVA(I,1),YV(I,1),YVA(I,1),ZV(I,1)
  DIMENSION ZVA(I,1),TRP(I,1),TRPA(I,1)
  COMMON /FORTH/NBX
```

AVG0050
AVG0060

THE PURPOSE OF THIS SUBROUTINE IS TO COMPUTE THE AVERAGE FLOW
FIELD PROPERTIES.

AVG0070
AVG0080

```
DO 110 N=1,NBX
DO 100 MT=1,IP
```

AVG0090
AVG0110

```
A=NBT(MT,N)
B=NB(MT,N)
```

AVG0150

```
C=A+B
```

AVG0170

```
NBT(MT,N)=C
```

AVG0180

```
IF(C.LT.1.) GO TO 100
```

AVG0190

```
DBA(MT,N)=(DBA(MT,N)*A+DB(MT,N)*B)/C
```

AVG0200

```
XVA(MT,N)=(XVA(MT,N)*A+XV(MT,N)*B)/C
```

AVG0210

```
YVA(MT,N)=(YVA(MT,N)*A+YV(MT,N)*B)/C
```

AVG0220

```
ZVA(MT,N)=(ZVA(MT,N)*A+ZV(MT,N)*B)/C
```

AVG0230

```
TMPA(MT,N)=(TMPA(MT,N)*A+TMP(MT,N)*B)/C
```

AVG0240

```
TRPA(MT,N)=(TRPA(MT,N)*A+TRP(MT,N)*B)/C
```

AVG0250

```
100 CONTINUE
```

AVG0260

```
110 CONTINUE
```

AVG0270

```
  RETURN
```

AVG0280

```
SUBROUTINE DRAG(AKN,NWEDG,XSTART,I2,I3,I4,I5,DELANG,JNT,BTA,
1C2,C3,DFA,FL,HTI,HTR,TB,XCB,CTI,CTR,CNI,CNR,ALPHA,SIGMA,RM,HTS,
2HTSI,NTS,UTL,UTT,VTS,LCOL,IP,ER,CHI,CNG,CMG,I,UTLI,UTTI,VTSI)
```

DRG0040

```
  INTEGER*2 LCOL
```

DRG0050

```
  INTEGER TIME,TST
```

DRG0060

```
  REAL LAM,MU,NU,JAY,KAY
```

```
  DIMENSION BTA(1),C2(1),C3(1),FL(1),HTI(1),CTI(3,1),CTR(3,1)
```

DRG0140

```
  DIMENSION HTR(1),TB(1),XCB(1),ALPHA(3,1),SIGMA(3,1)
```

DRG0150

```
  DIMENSION HTS(3,I2,I3),HTSI(3,I2,I3),NTS(3,I2,I3)
```

DRG0160

```
  DIMENSION UTL(3,I2,I3),UTT(3,I2,I3),VTS(3,I2,I3),LCOL(I,1)
```

170

```
  DIMENSION UTLI(3,I2,I3),UTTI(3,I2,I3),VTSI(3,I2,I3)
```

180

```
  DIMENSION CNI(3,1),CNR(3,1),DFA(1),JNT(1)
```

190

```
  DIMENSION ER(I,1),CNG(1),CMG(1),CHI(1)
```

200

```
  COMMON /THIRD/PI
```

210

```
  COMMON /PIFTH/ND,TIME,DTM,TI,ITS,ITP,TST
```

THE PURPOSE OF THIS SUBROUTINE IS TO ACCUMULATE THE DRAG AND HEAT
TRANSFER INCREMENTS ON THE BODY CONTRIBUTED BY EACH MOLECULE WHICH
COLLIDES WITH THE BODY. IN ADDITION, EACH MOLECULE WHICH COLLIDES
WITH THE BODY IS ASSIGNED AN APPROPRIATE NEW VELOCITY (OF REFLECTION)

WHICH IS USED TO CONTINUE ITS SPATIAL TRANSLATION (IN SUBROUTINE DRAG)

230

```

CALL NORMAL(JAY,KAY,RM)
JNT(MT)=JNT(MT)+1
TANG=180.*ATAN2(ZCL,-YCL)/PI
IWDG=TANG/DELANG+1
IF(IWDG.GT.NWEDG) IWDG=NWEDG
D=(LAM*LAM+MU*MU+NU*NU)
G=ER(MT,N)
H=G
DO 100 M=2,ND
IF(XCL.LT.XCB(M)) GO TO 110
100 CONTINUE
110 TBX=TB(M-1)+(TB(M)-TB(M-1))*(XCL-XCB(M-1))/(XCB(M)-XCB(M-1))
WI=(NU*JAY-MU*KAY)
VID=MU*JAY+NU*KAY
UID=LAM
E=RAND(0)
IF(E.LT.SIGMA(MT,M)) GO TO 115
VRD=-VID
URD=UID
WR=WI
GO TO 125
115 V=4.*RAND(0)
C2(MT)=C2(MT)+.544331*V*V*V*EXP(1.5-V*V)
IF(C2(MT).LT.1.) GO TO 115
C2(MT)=C2(MT)-1.
IF(NTS(MT,M,IWDG).NE.0) GO TO 117
ATR=ALPHA(MT,M)*TBX/SIGMA(MT,M)
GO TO 118
117 ATR=ALPHA(MT,M)*TBX/SIGMA(MT,M)+(1.-ALPHA(MT,M)/SIGMA(MT,M))*HTS DRG0530
118 I(MT,M,IWDG)/NTS(MT,M,IWDG)/(3.+CHI(MT))
ABR=SQRT(ATR)
V=V*ABR/BTA(MT)
AA=RAND(0)
A=SQRT(AA)
B=SQRT(1.-AA)
C=2.*PI*RAND(0)
VRD=V*B*COS(C)
WR=V*B*SIN(C)
IF(CHI(MT).EQ.-1.) GO TO 125
122 X=9.*RAND(0)
IF(X.EQ.0.0) GO TO 122
XTEMP=1.0
IF(CHI(MT).NE.0.0) XTEMP=X**CHI(MT)
CNG(MT)=CNG(MT)+XTEMP*EXP(-X)
IF(CNG(MT).LT.CMG(MT)) GO TO 122
124 CONTINUE
CNG(MT)=CNG(MT)-CMG(MT)
IF(CNG(MT).GE.CMG(MT)) GO TO 124
ER(MT,N)=X*ATR
H=ER(MT,N)
125 UR=URD
LAM=UR

```

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DRG0320DRG0330
DRG0340

DRG0350

DRG0370

DRG0400
DRG0410

DRG0420

DRG0430

DRG0440

DRG0450

DRG0460

DRG0470

DRG0480

DRG0490

DRG0500

DRG0510

DRG0520

DRG0530

DRG0550

DRG0560

DRG0570

DRG0610

DRG0620

DRG0630

DRG0640

DRG0650

LE: GKBINT AUG82 A

PRINCETON UNIVERSITY TIME-SHARING SYSTEM

MU=JAY*VRD-KAY*WR
 NU=KAY*VRD+JAY*WR
 IF (TIME.LE.TST) RETURN
 XMZ=(XCL-XSTART)*AKN
 YMZ= RM*AKN
 B=(URD*URD+VRD*VRD+WR*WR)
 UTI=UID*UID+WI*WI
 UYI=-WI*KAY
 UYR=-WR*KAY
 PL(MT)=PL(MT)+DFA(MT)
 HTI(MT)=HTI(MT)+D+G
 HTR(MT)=HTR(MT)-B-H
 CTI(MT,1)=CTI(MT,1)+UID
 CTI(MT,2)=CTI(MT,2)+UYI
 CTI(MT,3)=CTI(MT,3)+(XMZ*UYI-YMZ*UID)
 CNI(MT,2)=CNI(MT,2)+VID*JAY
 CNI(MT,3)=CNI(MT,3)+XMZ*JAY*VID
 CTR(MT,1)=CTR(MT,1)-URD
 CTR(MT,2)=CTR(MT,2)-UYR
 CTR(MT,3)=CTR(MT,3)-(XMZ*UYR-YMZ*URD)
 CNR(MT,2)=CNR(MT,2)-VRD*JAY
 CNR(MT,3)=CNR(MT,3)-XMZ*JAY*VRD
 NTS(MT,M,IWDG)=NTS(MT,M,IWDG)+1
 UTLLI(MT,M,IWDG)=UTLLI(MT,M,IWDG)+UID
 UTL(MT,M,IWDG)=UTL(MT,M,IWDG)+(UID-URD)
 UTTI(MT,M,IWDG)=UTTI(MT,M,IWDG)+WI
 UTT(MT,M,IWDG)=UTT(MT,M,IWDG)+(WI-WR)
 VTSI(MT,M,IWDG)=VTSI(MT,M,IWDG)-VID
 VTS(MT,M,IWDG)=VTS(MT,M,IWDG)+(VRD-VID)
 HTSI(MT,M,IWDG)=HTSI(MT,M,IWDG)+D+G
 HTS(MT,M,IWDG)=HTS(MT,M,IWDG)+D-B+G-H
 RETURN
 END

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DRG0730
DRG0740
DRG0750
DRG0760
DRG0790
DRG0820
DRG0830
DRG0840
DRG0860
DRG0880
DRG0890
DRG0900
DRG0920
DRG0940
DRG0960
DRG0970
DRG0980
DRG1170
DRG1180

SUBROUTINE INTERS(X,Y,Z,RMS,KS)
 REAL LAM,MU,NU
 COMMON /SVNTH/LAM,MU,NU,MT,N,J,XI,YI,ZI, TUSE
 1 FORMAT(//'* SOMETHING IS WRONG IN INTERS*/6E15.6/4E15.6)
 2 FORMAT(//'* TROUBLE IN INTERS - Tyme =*,E15.6/6E15.6/4E15.6)
 A=MU**2+NU**2
 B=(YI*MU+ZI*NU)/A
 C=(RMS-YI**2-ZI**2)/A
 IF (C.LT.0.0) C=0.0
 DISCR=B**2+C
 IF (DISCR.GE.0.0) GO TO 10
 WRITE(6,1) XI,YI,ZI,MU,NU,RMS,A,B,C,DISCR
 STOP
 10 Tyme=SQRT(DISCR)-B
 IF (C.EQ.0.0) Tyme=0.0
 IF ((Tyme.LE.Tuse).AND.(Tyme.GE.0.0)) GO TO 11
 WRITE(6,2) Tyme,XI,YI,ZI,MU,NU,RMS,A,B,C,DISCR
 IF (Tyme.GT.Tuse) Tyme=Tuse
 IF (Tyme.LT.0.0) Tyme=0.0
 11 KS=1
 Tuse=Tyme

```

XI=XI+LAM*TYME
YI=YI+MU*TYME
ZI=ZI+NU*TYME
IF(ZI.GT.0.0) GO TO 20
ZI=-ZI
NU=-NU
20 CONTINUE
YI=.9999*YI
ZI=.9999*ZI
X=XI
Y=YI
Z=ZI
RETURN
END

```

ORIGINAL PAGE 19
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```

SUBROUTINE NORMAL(JAY,KAY,RM)
REAL LAM,MU,NU,JAY,KAY
COMMON /SVNTH/LAM,MU,NU,MT,N,J,XCL,YCL,ZCL
JAY=-YCL/RM
KAY=-ZCL/RM
RETURN
END

```

NORM020
NORM040
NORM130
NORM140

```

SUBROUTINE PRINT1(DT,COSANG,SINANG,RMA,RNU,DRF,FCF,HTF,PL,HTI,HTR,
1CTI,CTR,CNI,CNR)
DIMENSION DD(3),WD(2,5),PP(4,4),QQ(4,4),RR(4,4),SS(4,4),TT(4,4)
DIMENSION UU(4,4),P1(4,4),Q1(4,4),R1(4,4),PA(4),PB(4),PC(4)
DIMENSION PL(1),HTI(1),HTR(1),CTI(3,1),CTR(3,1),CNI(3,1),CNR(3,1)
DIMENSION RMA(1),RNU(1)
DATA WD/'X-PO','RCE ','Y-PO','RCE ','Z-MO','MENT','DRAG',' ',' ','LPT10070
1IPT',' '/
----- PT10080
----- PT10090
----- PT10100
----- PT10110
----- PT10120
----- PT10130
----- PT10140
----- PT10150
----- PT10160
----- PT10170
----- PT10180

```

THE PURPOSE OF THIS SUBROUTINE IS TO PRINT OUT THE GROSS SURFACE COEFFICIENTS OF THE BODY.

```

FORMAT(//1X,50('*'),' GROSS SURFACE COEFFICIENTS ',50('*')/' MOLEC
1ULAR WEIGHT',12X,F8.3,3(19X,F8.3)/25X,
2 'INC.      REF.      TOT.      INC.      REF.      TOT.      INC.
3REF.      TOT.      INC.      REF.      TOT.')
10 FORMAT(' NUMBER FLUX      ',4(F9.3,18X))
12 FORMAT(1X,2A4,2X,'SHEAR      ',4(3F8.3,3X))
14 FORMAT(11X,'PRESSURE ',4(3F8.3,3X))
16 FORMAT(11X,'TOTAL      ',4(3F8.3,3X)/)
18 FORMAT(' HEAT TRANSFER',7X,4(3F8.3,3X)/)

```

PT10280
***** PT10290
PT10300
PT10310

RMR=0.0

```

DO 50 MT=1,3
DD(MT)=RMA(MT)*RNU(MT)*DRF/DT
50 RMR=RMR+RMA(MT)*RNU(MT)
WRITE(6,1) (RMA(MT),MT=1,3),RMR
PF=PL(1)*FCF/DT
QF=PL(2)*FCF/DT
RF=PL(3)*FCF/DT
SP=PF+QF+RF
WRITE(6,10) PF,QF,RF,SP
DO 200 I=1,3
PP(4,I)=0.0
QQ(4,I)=0.0
RR(4,I)=0.0
SS(4,I)=0.0
TT(4,I)=0.0
UU(4,I)=0.0
P1(4,I)=0.0
Q1(4,I)=0.0
R1(4,I)=0.0
DO 150 MT=1,3
PP(MT,I)=CTI(MT,I)*DD(MT)/RMR
QQ(MT,I)=CTR(MT,I)*DD(MT)/RMR
SS(MT,I)=CNI(MT,I)*DD(MT)/RMR
TT(MT,I)=CNE(MT,I)*DD(MT)/RMR
P1(MT,I)=PP(MT,I)+SS(MT,I)
PT10400
QQ(MT,I)=QQ(MT,I)+TT(MT,I)
PT10460
RR(MT,I)=PP(MT,I)+QQ(MT,I)
PT10470
UU(MT,I)=SS(MT,I)+TT(MT,I)
PT10480
R1(MT,I)=P1(MT,I)+Q1(MT,I)
PT10490
PP(4,I)=PP(4,I)+PP(MT,I)
QQ(4,I)=QQ(4,I)+QQ(MT,I)
RR(4,I)=RR(4,I)+RR(MT,I)
SS(4,I)=SS(4,I)+SS(MT,I)
TT(4,I)=TT(4,I)+TT(MT,I)
UU(4,I)=UU(4,I)+UU(MT,I)
P1(4,I)=P1(4,I)+P1(MT,I)
Q1(4,I)=Q1(4,I)+Q1(MT,I)
R1(4,I)=R1(4,I)+R1(MT,I)
150 CONTINUE
WRITE(6,12) (WD(J,I),J=1,2), (PP(K,I),QQ(K,I),RR(K,I),K=1,4)
WRITE(6,14) (SS(K,I),TT(K,I),UU(K,I),K=1,4)
WRITE(6,16) (P1(K,I),Q1(K,I),R1(K,I),K=1,4)
200 CONTINUE
AA=COSANG
PT10630
BB=SINANG
PT10640
DO 300 I=4,5
PT10650
DO 250 K=1,4
PT10660
PP(K,4)=AA*PP(K,1)+BB*PP(K,2)
PT10680
QQ(K,4)=AA*QQ(K,1)+BB*QQ(K,2)
PT10690
RR(K,4)=AA*RR(K,1)+BB*RR(K,2)
PT10700
SS(K,4)=AA*SS(K,1)+BB*SS(K,2)
PT10710
TT(K,4)=AA*TT(K,1)+BB*TT(K,2)
PT10720
UU(K,4)=AA*UU(K,1)+BB*UU(K,2)
PT10730
P1(K,4)=AA*P1(K,1)+BB*P1(K,2)
PT10740
Q1(K,4)=AA*Q1(K,1)+BB*Q1(K,2)
PT10750

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250 R1(K,4)=AA*R1(K,1)+BB*R1(K,2) PT10760
 WRITE(6,12)(WD(J,I),J=1,2),(PP(K,4),QQ(K,4),RR(K,4),K=1,4)
 WRITE(6,14)(SS(K,4),TT(K,4),UU(K,4),K=1,4)
 WRITE(6,16)(P1(K,4),Q1(K,4),R1(K,4),K=1,4)
 AA=-SINANG PT10800
 BB=COSANG PT10810
 300 CONTINUE PT10820
 HD=HTP/DT PT10830
 PA(4)=0.0
 PB(4)=0.0
 PC(4)=0.0
 DO 400 MT=1,3
 PA(MT)=HTI(MT)*RMA(MT)*RNU(MT)*HD/RMR
 PB(MT)=HTR(MT)*RMA(MT)*RNU(MT)*HD/RMR
 PC(MT)=PA(MT)+PB(MT)
 PA(4)=PA(4)+PA(MT)
 PB(4)=PB(4)+PB(MT)
 PC(4)=PC(4)+PC(MT)
 400 CONTINUE
 WRITE(6,18)(PA(I),PB(I),PC(I),I=1,4)
 RETURN PT10950
 END PT10960

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SUBROUTINE PRINT2(AKN,XSTART,DT,RNU,RMA,DRP,PCP,HTP,UTLI,UTTI,VTSI
 1,HTSI,DELANG,NWEDG,XS,XCB,YCB,HTS,NTS,UTL,UTT,VTS,I2,I3,IP)
 DIMENSION RMA(1),RNU(1),XS(1),XCB(1),YCB(1)
 DIMENSION HTS(3,I2,I3),NTS(3,I2,I3),UTL(3,I2,I3)
 DIMENSION UTT(3,I2,I3),VTS(3,I2,I3),UTLI(3,I2,I3),UTTI(3,I2,I3)
 DIMENSION VTSI(3,I2,I3),HTSI(3,I2,I3)
 COMMON /FIFTH/ND

PT20060

----- PT20070

THE PURPOSE OF THIS SUBROUTINE IS TO PRINT OUT THE DISTRIBUTION PT20080
 ON SURFACE OF THE SURFACE COEFFICIENTS PT20090

----- PT20100

PT20110

PT20120

PT20130

PT20140

PT20150

PT20160

FORMATS

8 FORMAT(//1X,45('''),' DISTRIBUTION ON SURFACE ',45(''')/71X,'INC.
 1 TOT. INC. TOT. INC. TOT.'/11X,'SEGMENT GEOMETRY',
 214X,'MOL. MOLE SAMP NUM. SKIN SKIN PRES- PRES-
 3 HEAT HEAT'/' NO. CENTER DELX CENTER DELANG',4X,'WGHT. P
 4RACT.',10X,
 5 'FLUX FRCTN FRCTN SURE SURE TRNSF TRNSF')
 10 FORMAT(1X,I3,F8.3,F7.3,F9.3,F8.3,1X,2F8.4,I6,7F8.4)
 11 FORMAT(37X,F8.4,' 1.0000',I6,7F8.4)

PT20230

***** PT20240

PT20250

PT20260

RMR=0.0
 DO 50 MT=1,IP
 50 RMR=RMR+RMA(MT)*RNU(MT)

```

      WRITE(6,8)
      I=0
      DO 110 N=2,ND
      DTY=DT*YCB(N)/180.
      P=XS(N)
      Q=2.*((XCB(N)-XSTART)*AKN-XS(N))
      J=0
      R=-.5*DELANG
      PMLT=FCP/(DTY*DELANG)
      QMLT=DRP/(DTY*DELANG)
      SMLT=HTF/(DTY*DELANG)
      DO 100 K=1,NWEDG
      R=R+DELANG
      I=I+1
      J=J+1
      M3=0
      P3=0.0
      Q3=0.0
      Q4=0.0
      R3=0.0
      R4=0.0
      S3=0.0
      S4=0.0
      DO 90 MT=1,IP
      M1=NTS(MT,N,J)
      M3=M3+M1
      P1=NTS(MT,N,J)*PMLT*RNU(MT)
      P3=P3+P1
      Q1=SQRT(UTLI(MT,N,J)**2+UTTI(MT,N,J)**2)*RNU(MT)*RMA(MT)*QMLT/RMR
      Q2=SQRT(UTL(MT,N,J)**2+UTT(MT,N,J)**2)*RNU(MT)*RMA(MT)*QMLT/RMR
      Q3=Q3+Q1
      Q4=Q4+Q2
      R1=VTSI(MT,N,J)*RNU(MT)*RMA(MT)*QMLT/RMR
      R2=VTS(MT,N,J)*RNU(MT)*RMA(MT)*QMLT/RMR
      S1=HTSI(MT,N,J)*RNU(MT)*RMA(MT)*SMLT/RMR
      S2=HTS(MT,N,J)*RNU(MT)*RMA(MT)*SMLT/RMR
      R3=R3+R1
      R4=R4+R2
      S3=S3+S1
      S4=S4+S2
  90  WRITE(6,10)I,P,Q,R,DELANG,RMA(MT),RNU(MT),M1,P1,Q1,Q2,R1,R2,S1,S2
      WRITE(6,11)RMR,M3,P3,Q3,Q4,R3,R4,S3,S4
  100 CONTINUE
  110 CONTINUE
      RETURN
      END

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PT20280
PT20320
PT20330
PT20340
PT20350
PT20380
PT20490
PT20500
PT20660
PT20670
PT20680
PT20690

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      SUBROUTINE PRINT4(MSP,CHI,RNU,I,TRP,FDN,WTM,DB,NS,TMP,XV,
      1YV,ZV,KS,NB,XC,YC,ZC)
      INTEGER*2 NB,NS
      DIMENSION FDN(1),RNU(1),CHI(1),WTM(1),TMP(I,1),TRP(I,1)
      DIMENSION DB(I,1),NB(I,1),XV(I,1),YV(I,1),ZV(I,1),DBT(3),NS(I,1)
      DIMENSION XC(1),YC(1),ZC(1)
      COMMON /SECND/BW,BH,RMP,RMN,RMF
      COMMON /FORTH/NBX,RM,XR

```

 THE PURPOSE OF THIS SUBROUTINE IS TO PRINT OUT THE INSTANTANEOUS PT50060
 FLOW-FIELD PROPERTIES. PT50070
 PT50080
 PT50090
 PT50100
 PT50110
 PT50120
 PT50130
 PT50140
 PT50150

FORMATS

1 FORMAT(//1X,45('*'),' INSTANTANEOUS FLOW FIELD INFORMATION ',45('*' PT50160
 1')) PT50170
 2 FORMAT(/2X,'HORIZONTAL NUMBER=',I3,3X,'VERTICAL NUMBER=',I3,3X,'HO
 1RIZONTAL POSITION=',F8.5,3X,' RADIAL POSITION =',F8.5/2X,'BOX#
 2ANGLE Samp DENSITY MACH NO X VEL. Y VEL. Z VEL. T(KIN) T(RO
 3T) TEMP.',14X,'MOLE FRACTIONS')
 3 FORMAT(//1X,46('*'),' ACCUMULATED FLOW FIELD INFORMATION ',46('*'))
 1)
 4 FORMAT(1X,I4,E11.3,I6,8F8.3,3E11.3)

PT50220
 ***** PT50230
 PT50240
 PT50250

IP(KS.EQ.0) WRITE(6,1)
 IP(KS.NE.0) WRITE(6,3)
 DO 40 MT=1,3
 40 DBT(MT)=0.0
 FDA=0.
 CHT=0.
 DO 50 MT=1,MSP
 CHT=CHT+CHI(MT)*RNU(MT)
 50 FDA=FDA+FDN(MT)*WTM(MT)
 YCT=0.0
 XCT=0.0
 DO 110 N=1,NBX
 IF ((XC(N).EQ.XCT).AND.(YC(N).EQ.YCT)) GO TO 52 PT50290
 XCT=XC(N)
 YCT=YC(N)
 XXT=XCT/XR
 YYT=YCT/RM
 IXC=XC/BW + 1
 IYC=YCT/BH + 1
 WRITE(6,2) IXC,IYC,XXT,YYT
 52 ZCT=ZC(N)
 NSAMP=0
 DBA=0.
 XVM=0.
 YVM=0.
 ZVM=0.
 TMPM=0.
 TRPM=0.
 E=0.
 F=0.
 DO 100 MT=1,MSP
 NSAMP=NSAMP+NS(MT,N)

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XVM=XVM+XV (MT,N)*RNU (MT)*WTM (MT)*NB (MT,N)
YVM=YVM+YV (MT,N)*RNU (MT)*WTM (MT)*NB (MT,N)
ZVM=ZVM+ZV (MT,N)*RNU (MT)*WTM (MT)*NB (MT,N)
DBA=DBA+DB (MT,N)*WTM (MT)
TMPM=TMPM+WTM (MT)*RNU (MT)*TMP (MT,N)*NB (MT,N)
TRPM=TRPM+RNU (MT)*NB (MT,N)*TRP (MT,N)
E=E+WTM (MT)*RNU (MT)*NB (MT,N)
100 F=F+RNU (MT)*NB (MT,N)
DBA=DBA/FDA
IF (E.EQ.0.0) GO TO 55
XVM=XVM/E
YVM=YVM/E
ZVM=ZVM/E
VS=XVM**2+YVM**2+ZVM**2
TMPM=TMPM/E-VS
TRPM=TRPM/F
55 CONTINUE
TTM=(TMPM+TRPM)/(2.5+CHT)
TMPM=TMPM/1.5
IF (CHT.NE.-1.) TRPM=TRPM/(1.+CHT)
AMS=SQRT(VS)
IF (TTM.GT.0.) AMS=SQRT((5.+2.*CHT)*VS/(TTM*(3.5+CHT)))
CCZ=COS(ZCT/57.29578)
SCZ=SQRT(1.-CCZ**2)
RVM=ZVM*SCZ-YVM*CCZ
TVM=YVM*SCZ+ZVM*CCZ
DO 60 MT=1,MSP
DBT (MT)=RNU (MT)*NB (MT,N)
IF (F.NE.0.) DBT (MT)=DBT (MT)/F
60 CONTINUE
WRITE(6,4) N,ZCT,NSAMP,DBA,AMS,XVM,RVM,TVM,TMPM,TRPM,TTM,(DBT (J),
1J=1,3)
110 CONTINUE
RETURN
END

```

PT50470
PT50480
PT50490

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:INTRY
:CONTRL NAME='INTE','RNAL',TITLE=' PAR','ABOL','A AT',' 95K','M M','ON. ',
:DEBUG=.F.,.F.,.T.,NEW=.F.,SAVE=.T.,ICOPY=0,REDO=.T. &END
:TIMES DTM=.005,ITS=6,ITP=6,TST=2,TLIM=12 &END
:FLOREF LLM=2000,MNM=5000,MNB=150,MSP=1,MET=0,U=7485.9,ANGLE=28.,RNU=1.,2*0.,
:RMA=28.94,0.,0.,TF=195.51,DENF=2.52E+19 &END
:MOLEC TRF=1000,DIR=3.5E-19,ETA=.104,PHI=0.0,CHI=-1.,ACR=.001 &END
:SHAPES BODY=0.0,1000.,.00235 &END
:SHAPES BODY=.0025,555.,0.0,2*1.0 &END
:SHAPES BODY=.0050,345.,0.0,2*1.0 &END
:SHAPES BODY=.0100,300.,0.0,2*1.0 &END
:SHAPES BODY=.0200,300.,0.0,2*1.0 &END
:SHAPES BODY=.0300,300.,0.0,2*1.0 &END
:SHAPES BODY=.0400,300.,0.0,2*1.0 &END
:SHAPES BODY=.0500,300.,0.0,2*1.0 &END
:SHAPES BODY=.0600,300.,0.0,2*1.0 &END
:SHAPES BODY=.0700,300.,0.0,2*1.0 &END
:SHAPES BODY=.0800,300.,0.0,2*1.0 &END
:SHAPES BODY=.0870,300.,1.0,2*1.0 &END
:GEOM NWEDG=2,NW=20,NH=3, &END

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;INCUP1 FLUXIN=2.1429,FCOL=1.0,RMP=0.3,JV=22,KMX=3 &END
;INOUT VARG=0.,1.,2.,3.,4.,5.,6.,7.,8.,9.,10.,11.,12.,13.,14.,15.,16.,17.,18.,
19.,20.,21.,CURV=0.0,.070,.170,.282,.369,.459,.537,.599,.656,.710,.750,.785,
.815,.845,.872,.900,.922,.951,.975,.988,.996,1.00, &END
;INOUT VARG=-20.,-19.,-17.,-15.,-13.,-11.,-9.,-7.,-5.,-3.,-1.,1.,3.,5.,7.,9.,
11.,13.,15.,17.,19.,CURV=4*0.,.003,.013,.036,.084,.149,.250,.406,.611,.762,
.871,.932,.962,.984,.995,.999,3*1.0, &END
;INOUT VARG=-20.,-19.,-17.,-15.,-13.,-11.,-9.,-7.,-5.,-3.,-1.,1.,3.,5.,7.,9.,
11.,13.,15.,17.,19.,CURV=4*0.0,.003,.013,.036,.084,.149,.250,.406,.611,.762,
.871,.932,.962,.984,.995,.999,3*1.0, &END
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